# ANALYSIS \& PDE 

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Nils Dencker


#### Abstract

In this paper we consider the solvability of pseudodifferential operators when the principal symbol vanishes of at least second order at a nonradial involutive manifold $\Sigma_{2}$. We shall assume that the subprincipal symbol is of principal type with Hamilton vector field tangent to $\Sigma_{2}$ at the characteristics, but transversal to the symplectic leaves of $\Sigma_{2}$. We shall also assume that the subprincipal symbol is essentially constant on the leaves of $\Sigma_{2}$ and does not satisfying the Nirenberg-Trèves condition $(\Psi)$ on $\Sigma_{2}$. In the case when the sign change is of infinite order, we also need a condition on the rate of vanishing of both the Hessian of the principal symbol and the complex part of the gradient of the subprincipal symbol compared with the subprincipal symbol. Under these conditions, we prove that $P$ is not solvable.


## 1. Introduction

We will consider the solvability for a classical pseudodifferential operator $P \in \Psi_{\mathrm{cl}}^{m}(M)$ on a $C^{\infty}$ manifold $M$. This means that $P$ has an expansion $p_{m}+p_{m-1}+\cdots$, where $p_{k} \in S_{\text {hom }}^{k}$ is homogeneous of degree $k$ for all $k$, and $p_{m}=\sigma(P)$ is the principal symbol of the operator. A pseudodifferential operator is said to be of principal type if the Hamilton vector field $H_{p_{m}}$ of the principal symbol does not have the radial direction $\xi \cdot \partial_{\xi}$ on $p_{m}^{-1}(0)$, in particular $H_{p_{m}} \neq 0$. We shall consider the case when the principal symbol vanishes of at least second order at an involutive manifold $\Sigma_{2}$; then $P$ is not of principal type.
$P$ is locally solvable at a compact set $K \subseteq M$ if the equation

$$
\begin{equation*}
P u=v \tag{1-1}
\end{equation*}
$$

has a local solution $u \in \mathcal{D}^{\prime}(M)$ in a neighborhood of $K$ for any $v \in C^{\infty}(M)$ in a set of finite codimension. We can also define microlocal solvability of $P$ at any compactly based cone $K \subset T^{*} M$; see Definition 2.6.

For pseudodifferential operators of principal type, it is known [Dencker 2006; Hörmander 1981] that local solvability is equivalent to condition $(\Psi)$ on the principal symbol, which means that

$$
\begin{equation*}
\text { Im } a p_{m} \text { does not change sign from }- \text { to }+ \text { along the oriented bicharacteristics of } \operatorname{Re} a p_{m} \tag{1-2}
\end{equation*}
$$

for any $0 \neq a \in C^{\infty}\left(T^{*} M\right)$. The oriented bicharacteristics are the positive flow-out of the Hamilton vector field $H_{\operatorname{Re} a p_{m}} \neq 0$ on which $\operatorname{Re} a p_{m}=0$; these are also called semibicharacteristics of $p_{m}$. Condition (1-2) is invariant under multiplication of $p_{m}$ with nonvanishing factors, and symplectic changes of variables; thus it is invariant under conjugation of $P$ with elliptic Fourier integral operators. Observe that the sign changes in (1-2) are reversed when taking adjoints, and that it suffices to check (1-2) for some $a \neq 0$ for which $H_{\operatorname{Re} a p} \neq 0$, according to [Hörmander 1985b, Theorem 26.4.12].

[^0]For operators which are not of principal type, the situation is more complicated and the solvability may depend on the lower-order terms. When the set $\Sigma_{2}$, where the principal symbol vanishes of second order, is involutive, the subprincipal symbol $\sigma_{\text {sub }}(P)=p_{m-1}$ is invariantly defined at $\Sigma_{2}$. In fact, on $\Sigma_{2}$ it is equal to the refined principal symbol; see [Hörmander 1985a, Theorem 18.1.33].

In the case where the principal symbol is real and vanishes of at least second order at the involutive manifold, there are several results, mostly in the case when the principal symbol is a product of real symbols of principal type. Then the operator is not solvable if the imaginary part of the subprincipal symbol has a sign change of finite order on a bicharacteristic of one the factors of the principal symbol; see [Egorov 1977; Popivanov 1974; Wenston 1977; 1978].

This necessary condition for solvability has been extended to some cases when the principal symbol is real and vanishes of second order at the involutive manifold. The conditions for solvability then involve the sign changes of the imaginary part of the subprincipal symbol on the limits of bicharacteristics from outside the manifold, thus on the leaves of the symplectic foliation of the manifold; see [Mendoza and Uhlmann 1983; 1984; Mendoza 1984; Yamasaki 1983]. This has been extended to more general limit bicharacteristics of real principal symbols in [Dencker 2016].

When $\Sigma_{2}$ is not involutive, there are examples where the operator is solvable for any lower-order terms. For example when $P$ is effectively hyperbolic, then even the Cauchy problem is solvable for any lower-order term; see [Hörmander 1977; Nishitani 2004]. There are also results in the cases when the principal symbol is a product of principal-type symbols not satisfying condition ( $\Psi$ ); see [Cardoso and Trèves 1974; Gilioli and Trèves 1974; Goldman 1975; Trèves 1973; Yamasaki 1980].

In the present paper, we shall consider the case when the principal symbol (not necessarily realvalued) vanishes of at least second order at a nonradial involutive manifold $\Sigma_{2}$. We shall assume that the subprincipal symbol is of principal type with Hamilton vector field tangent to $\Sigma_{2}$ at the characteristics, but transversal to the symplectic leaves of $\Sigma_{2}$. We shall also assume that the subprincipal symbol is essentially constant on the symplectic leaves of $\Sigma_{2}$ by (2-8), and does not satisfy condition $(\Psi)$; see Definition 2.4. In the case when the sign change is of infinite order, we will need a condition on the rate of vanishing of both the Hessian of the principal symbol and the complex part of the gradient of the subprincipal symbol on the semibicharacteristic of the subprincipal symbol; see condition (2-11). Under these conditions, $P$ is not solvable in a neighborhood of the semibicharacteristic; see Theorem 2.7, which is the main result of the paper. In this case $P$ is an evolution operator; see [Colombini et al. 2003; 2010] for some earlier results on the solvability of evolution operators.

## 2. Statement of results

Let $\sigma(P)=p \in S_{\text {hom }}^{m}$ be the homogeneous principal symbol. We shall assume that

$$
\begin{equation*}
\sigma(P) \text { vanishes of at least second order at } \Sigma_{2} \tag{2-1}
\end{equation*}
$$

where

$$
\begin{equation*}
\Sigma_{2} \text { is a nonradial involutive manifold. } \tag{2-2}
\end{equation*}
$$

Here nonradial means that the radial direction $\left\langle\xi, \partial_{\xi}\right\rangle$ is not in the span of the Hamilton vector fields of the manifold, i.e., not equal to $H_{f}$ on $\Sigma_{2}$ for some $f \in C^{1}$ vanishing at $\Sigma_{2}$. Then by a change of homogeneous symplectic coordinates we may assume that

$$
\begin{equation*}
\Sigma_{2}=\left\{\xi^{\prime}=0\right\}, \quad \xi=\left(\xi^{\prime}, \xi^{\prime \prime}\right) \in \mathbb{R}^{k} \times \mathbb{R}^{n-k} \tag{2-3}
\end{equation*}
$$

for some $k>0$; this can be achieved by conjugation by elliptic Fourier integral operators. If $P$ is of principal type near $\Sigma_{2}$ then, since solvability is an open property, we find that a necessary condition for $P$ to be solvable at $\Sigma_{2}$ is that condition ( $\Psi$ ) for the principal symbol is satisfied in some neighborhood of $\Sigma_{2}$. Naturally, this condition is empty on $\Sigma_{2}$, where we instead need some conditions on the subprincipal symbol

$$
\begin{equation*}
p_{s}=p_{m-1}+\frac{i}{2} \sum_{j} \partial_{x_{j}} \partial_{\xi_{j}} p \tag{2-4}
\end{equation*}
$$

which is equal to $p_{m-1}$ on $\Sigma_{2}$ and invariantly defined as a function on $\Sigma_{2}$ under symplectic changes of coordinates and conjugation with elliptic pseudodifferential operators. (In the Weyl quantization, the subprincipal symbol is equal to $p_{m-1}$.) When composing $P$ with an elliptic pseudodifferential operator $C$, the value of the subprincipal symbol of $C P$ is equal to $c p_{s}+\frac{1}{2} i H_{p} c=c p_{s}$ at $\Sigma_{2}$, where $c=\sigma(C)$. Observe that the subprincipal symbol is complexly conjugated when taking the adjoint of the operator.

Let $T \Sigma_{2}^{\sigma}$ be the symplectic polar to $T \Sigma_{2}$, which spans the symplectic leaves of $\Sigma_{2}$. If $\Sigma_{2}=\left\{\xi^{\prime}=0\right\}$ and $x=\left(x^{\prime}, x^{\prime \prime}\right) \in \mathbb{R}^{k} \times \mathbb{R}^{n-k}$, then the leaves are spanned by $\partial_{x^{\prime}}$. Let

$$
\begin{equation*}
T^{\sigma} \Sigma_{2}=T \Sigma_{2} / T \Sigma_{2}^{\sigma} \tag{2-5}
\end{equation*}
$$

which is a symplectic space over $\Sigma_{2}$, which in these coordinates is given by

$$
\begin{equation*}
T^{\sigma} \Sigma_{2}=\left\{\left(\left(x_{0}, 0, \xi_{0}^{\prime \prime}\right) ;\left(0, y^{\prime \prime}, 0, \eta^{\prime \prime}\right)\right):\left(y^{\prime \prime}, \eta^{\prime \prime}\right) \in T^{*} \mathbb{R}^{n-k}\right\} \tag{2-6}
\end{equation*}
$$

Next, we are going to study the Hamilton vector field $H_{p_{m-1}}$ at $\Sigma_{2}$. If $H_{p_{m-1}} \subseteq T \Sigma_{2}$ at $\Sigma_{2}$ then we find that $d p_{s}$ vanishes on $T \Sigma_{2}^{\sigma}$ so $d p_{s}$ is well defined on $T^{\sigma} \Sigma_{2}$. In fact, $p_{s}=p_{m-1}$ on $\Sigma_{2}$ so if we choose coordinates so that (2-3) holds, then $H_{p_{m-1}} \subseteq T \Sigma_{2}$ is equivalent to

$$
\begin{equation*}
H_{p_{m-1}} \xi^{\prime}=-\partial_{x^{\prime}} p_{m-1}=-\partial_{x^{\prime}} p_{s}=0 \quad \text { when } \xi^{\prime}=0 \tag{2-7}
\end{equation*}
$$

which is invariant under multiplication with nonvanishing factors when $p_{s}=0$. Let $H_{p_{s}}$ be the Hamilton vector field of $p_{s}$ with respect the symplectic structure on the symplectic manifold $T^{\sigma} \Sigma_{2}$. In the chosen coordinates we have

$$
H_{p_{s}}=\partial_{\xi^{\prime \prime}} p_{s} \partial_{x^{\prime \prime}}-\partial_{x^{\prime \prime}} p_{s} \partial_{\xi^{\prime \prime}}
$$

modulo $\partial_{x^{\prime}}$, which is nonvanishing if $\partial_{x^{\prime \prime} \xi^{\prime \prime}} p_{s} \neq 0$. Since $p_{s}=p_{m-1}$ on $\Sigma_{2}$, the difference between $H_{p_{m-1}}$ and $H_{p_{s}}$ is tangent to the leaves of $\Sigma_{2}$. Actually, since the subprincipal symbol is only well defined on $\Sigma_{2}$, the vector field $H_{p_{s}}$ is only well defined up to terms tangent to the leaves.

Because of that, we would need that the subprincipal symbol $p_{s}$ is constant on the leaves of $\Sigma_{2}$, but that condition is not invariant under multiplication with nonvanishing factors when $p_{s} \neq 0$. Instead we
shall use the invariant condition

$$
\begin{equation*}
\left|d p_{s}\right|_{T L}\left|\leq C_{0}\right| p_{s} \mid \tag{2-8}
\end{equation*}
$$

for any leaf $L$ of $\Sigma_{2}$. Then $p_{s}$ is constant on the leaves modulo nonvanishing factors, according to the following lemma.

Lemma 2.1. If $\left.d p_{s}\right|_{T^{\sigma} \Sigma_{2}} \neq 0$, then condition (2-8) is equivalent to the fact that $p_{s}$ is constant on the leaves of $\Sigma_{2}$ after multiplication with a smooth nonvanishing factor. Thus, if $\Sigma_{2}=\left\{\xi^{\prime}=0\right\}$ then (2-8) gives $p_{s}\left(x, 0, \xi^{\prime \prime}\right)=c\left(x, \xi^{\prime \prime}\right) q\left(x^{\prime \prime}, \xi^{\prime \prime}\right)$ with $0 \neq c \in C^{\infty}$.
Proof. Choose coordinates so that $\Sigma_{2}=\left\{\xi^{\prime}=0\right\}$. If $p_{s} \neq 0$ at a point $w_{0} \in \Sigma_{2}$ then (2-8) gives that $\partial_{x^{\prime}} \log p_{s}$ is uniformly bounded near $w_{0}$, where $\log p_{s}$ is a branch of the complex logarithm. Thus, by integrating with respect to $x^{\prime}$ in a simply connected neighborhood starting at $x^{\prime}=x_{0}^{\prime}$, we find that

$$
\begin{equation*}
p_{s}\left(x, 0, \xi^{\prime \prime}\right)=c\left(x, \xi^{\prime \prime}\right) q\left(x^{\prime \prime}, \xi^{\prime \prime}\right) \tag{2-9}
\end{equation*}
$$

where $q\left(x^{\prime \prime}, \xi^{\prime \prime}\right)=p_{s}\left(x_{0}^{\prime}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right) \in C^{\infty}$, so $0 \neq c \in C^{\infty}$ satisfies $c\left(x_{0}^{\prime}, x^{\prime \prime}, \xi^{\prime \prime}\right)=1$. When $p_{s}=0$ we find that $\left.d \operatorname{Re} z p_{s}\right|_{T^{\sigma} \Sigma_{2}} \neq 0$ for some $z \in \mathbb{C} \backslash\{0\}$ by assumption. Thus we obtain locally that

$$
p_{s}\left(x, 0, \xi^{\prime \prime}\right)=c_{ \pm}\left(x, \xi^{\prime \prime}\right) q_{ \pm}\left(x^{\prime \prime}, \xi^{\prime \prime}\right) \quad \text { on } S_{ \pm}=\left\{ \pm \operatorname{Re} z p_{s}\left(x, 0, \xi^{\prime \prime}\right)>0\right\}
$$

where $q_{ \pm}\left(x^{\prime \prime}, \xi^{\prime \prime}\right)=p_{s}\left(x_{0}^{\prime}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right), 0 \neq c_{ \pm} \in C^{\infty}$ and $c_{ \pm}\left(x_{0}^{\prime}, x^{\prime \prime}, \xi^{\prime \prime}\right)=1$ on $S_{ \pm}$. Then we find that $p_{s}^{-1}(0)$ is independent of $x^{\prime}$ and

$$
\partial_{x^{\prime \prime}}^{\alpha} \partial_{\xi^{\prime \prime}}^{\beta} q_{ \pm}\left(x^{\prime \prime}, \xi^{\prime \prime}\right)=\partial_{x^{\prime \prime}}^{\alpha} \partial_{\xi^{\prime \prime}}^{\beta} p_{s}\left(x_{0}^{\prime}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right) \quad \forall \alpha, \beta, \quad \text { on } S_{ \pm},
$$

so by taking the limit at $S=\left\{\operatorname{Re} z p_{s}=0\right\}$, we find that the functions $q_{ \pm}$extend to $q \in C^{\infty}$. Since $c_{+} q=c_{-} q=p_{s}$ at $S$, we find that $c_{+}=c_{-}$at $S$ when $q \neq 0$. When $q=0$ at $S$, we may differentiate in the normal direction of $S$ to obtain $c_{-} \partial_{\nu} q=c_{+} \partial_{\nu} q$, and since $\partial_{\nu} q \neq 0$, the functions $c_{ \pm}$extend to a continuous function $c$. By differentiating and taking the limit, we find that

$$
\nabla c_{-} q+c \nabla q=\nabla p_{s}=\nabla c_{+} q+c \nabla q \quad \text { at } S
$$

which similarly gives that $\nabla c_{-}=\nabla c_{+}$at $S$, so $c \in C^{1}$. By repeatedly differentiating $c_{ \pm} q$, we find by induction that $c \in C^{\infty}$, so we get smooth quotients $c$ and $q$ in (2-9).

Now, a semibicharacteristic of $p_{s}$ will be a bicharacteristic of $\operatorname{Re} a p_{s}$ on $T^{\sigma} \Sigma_{2}$, where $C^{\infty} \ni a \neq 0$, with the natural orientation. Observe that condition (2-7) is only invariant under multiplication with nonvanishing factors when $p_{s}=0$.
Definition 2.2. We say that the operator $P$ is of subprincipal type if, when $p_{s}=0$ on $\Sigma_{2}$, the following hold: $\left.H_{p_{m-1}}\right|_{\Sigma_{2}} \subseteq T \Sigma_{2}$,

$$
\begin{equation*}
\left.d p_{s}\right|_{T^{\sigma} \Sigma_{2}} \neq 0 \tag{2-10}
\end{equation*}
$$

and the corresponding Hamilton vector field $H_{p_{s}}$ of (2-10) does not have the radial direction. We call $H_{p_{s}}$ the subprincipal Hamilton vector field and the (semi)bicharacteristics are called the subprincipal (semi)bicharacteristics on $\Sigma_{2}$.

Clearly, if (2-3) holds, then the condition that the Hamilton vector field does not have the radial direction means that $\partial_{\xi^{\prime \prime}} p_{s} \neq 0$ or $\partial_{x^{\prime \prime}} p_{s} \nVdash \xi^{\prime \prime}$ when $p_{s}=0$ on $\Sigma_{2}=\left\{\xi^{\prime}=0\right\}$.

In the case when the principal symbol $p$ is real, a necessary condition for solvability of the operator is that the imaginary part of the subprincipal symbol does not change sign from - to + when going in the positive direction on a $C^{\infty}$ limit of normalized bicharacteristics of the principal symbol $p$ at $\Sigma_{2}$; see [Dencker 2016]. When $p$ vanishes of exactly second order on $\Sigma_{2}=\left\{\xi^{\prime}=0\right\}$, such limit bicharacteristics are tangent to the leaves of $\Sigma_{2}$. In fact, then Taylor's formula gives $H_{p}=\left\langle B \xi^{\prime}, \partial_{x^{\prime}}\right\rangle+\mathcal{O}\left(\left|\xi^{\prime}\right|^{2}\right)$, where $B \neq 0$, so the normalized Hamilton vector fields have limits that are tangent to the leaves. When the principal symbol is proportional to a real-valued symbol, this gives examples of nonsolvability when the subprincipal symbol is not constant on the leaves of $\Sigma_{2}$. Thus condition (2-8) is natural if there are no other conditions on the principal symbol.

Remark 2.3. If $p_{s}$ is real-valued, then by the proof of Lemma 2.1 it follows from (2-8) that $p_{s}$ has constant sign on the leaves of $\Sigma_{2}$, since then $c>0$ in (2-9).

Definition 2.4. We say that $P$ satisfies condition $\operatorname{Sub}(\Psi)$ if $\operatorname{Im} a p_{s}$ does not change sign from - to + when going in the positive direction on the subprincipal bicharacteristics of $\operatorname{Re} a p_{s}$ for any $0 \neq a \in C^{\infty}$.

Thus, condition $\operatorname{Sub}(\Psi)$ is condition ( $\Psi$ ) given by (1-2) on the subprincipal symbol $p_{s}$. Observe that since $p_{s}$ is only defined on $\Sigma_{2}$, the Hamilton vector field $H_{p_{s}}$ is only well defined in $T^{\sigma} \Sigma_{2}=T \Sigma_{2} / T \Sigma_{2}^{\sigma}$; thus it is well defined modulo $\partial_{x^{\prime}}$. But if $(2-8)$ holds then we find that $\operatorname{Sub}(\Psi)$ is a condition on $p_{s}$ with respect to the symplectic structure of $T^{\sigma} \Sigma_{2}$. In fact, by the invariance of condition ( $\Psi$ ) given by [Hörmander 1985b, Lemma 26.4.10], condition $\operatorname{Sub}(\Psi)$ holds for any $a \neq 0$ such that $H_{\operatorname{Re} a p_{s}} \neq 0$, so we may assume by Lemma 2.1 that $p_{s}$ is constant on the leaves of $\Sigma_{2}$.

Since condition $\operatorname{Sub}(\Psi)$ is invariant under symplectic changes of variables and multiplication with nonvanishing functions, it is invariant under conjugation of the operator by elliptic Fourier integral operators. Observe that the sign change is reversed when taking the adjoint of the operator.

Recall that the Hessian of the principal symbol Hess $p$ is the quadratic form given by $\partial^{2} p$ at $\Sigma_{2}$, which is defined on the normal bundle $N \Sigma_{2}$ since it vanishes on $T \Sigma_{2}$. By the calculus, Hess $p$ is invariant, modulo nonvanishing smooth factors, under symplectic changes of variables and multiplication of $P$ with elliptic pseudodifferential operators.

Next, we assume that condition $\operatorname{Sub}(\Psi)$ is not satisfied on a semibicharacteristic $\Gamma$ of $p_{s}$; that is, $\operatorname{Im} a p_{s}$ changes sign from - to + on the positive flow-out of $H_{\operatorname{Re} a p_{s}} \neq 0$ for some $0 \neq a \in C^{\infty}$. Now if the sign change is not of finite order, we shall also need an extra condition on the rate of vanishing of both the Hessian of the principal symbol and the complex part of the gradient of the subprincipal symbol on the subprincipal semibicharacteristic. Then, we shall assume that there exists $C>0, \varepsilon>0$ and $0 \neq a \in C^{\infty}$ so that $\left.d \operatorname{Re} a p_{s}\right|_{T \Sigma_{2}} \neq 0$ and

$$
\begin{equation*}
\| \text { Hess } p \|+\left|d p_{s} \wedge d \bar{p}_{s}\right| \leq C\left|p_{s}\right|^{\varepsilon} \quad \text { when } \operatorname{Re} a p_{s}=0 \text { on } \Sigma_{2} \tag{2-11}
\end{equation*}
$$

near $\Gamma$. Since (2-11) also holds for smaller $\varepsilon$ and larger $C$, it is no restriction to assume $\varepsilon \leq 1$. The motivation for (2-11) is to prevent the transport equation (6-1) from dispersing the support of the solution
before the sign change of the imaginary part of the subprincipal symbol localizes it; see Remark 3.1. We also find that $\nabla p_{s}$ is proportional to a real vector when $p_{s}=0$ since then $d p_{s} \wedge d \overline{p_{s}}=0$.

Remark 2.5. Condition (2-11) is invariant under multiplication of $P$ with elliptic pseudodifferential operators, and symplectic changes of coordinates. If (2-8) also holds, then we obtain

$$
\begin{equation*}
\| d \text { Hess }\left.p\right|_{T L} \| \leq C_{1}\left|p_{s}\right|^{\varepsilon / 2} \tag{2-12}
\end{equation*}
$$

for any leaf $L$ of $\Sigma_{2}$ when $\operatorname{Re} a p_{s}=0$ near $\Gamma$.
In fact, multiplication with an elliptic pseudodifferential operator with principal symbol $c$ changes the principal symbol into $c p$, the Hessian of the principal symbol into $c$ Hess $p$ and the subprincipal symbol into

$$
c p_{s}+\frac{i}{2} H_{p} c \quad \text { at } \Sigma_{2}
$$

where the last term vanishes at $\Sigma_{2}$ and contains the factor Hess $p$, modulo terms vanishing of second order at $\Sigma_{2}$. Now we have

$$
\left|d c p_{s} \wedge d \overline{c p_{s}}\right| \leq|c|^{2}\left|d p_{s} \wedge d \overline{p_{s}}\right|+C\left|p_{s}\right|
$$

Thus we find that (2-11) holds with $p$ replaced by $c p, p_{s}$ replaced by $c p_{s}$ and $a$ replaced with $a / c$. If (2-8) also holds and we choose coordinates so that $\Sigma_{2}=\left\{\xi^{\prime}=0\right\}$, then we obtain from Lemma 2.1 that $\left|p_{s}\left(x^{\prime}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right)\right| \cong\left|p_{s}\left(x_{0}^{\prime}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right)\right|$ when $\left|x^{\prime}-x_{0}^{\prime}\right| \leq c$. Thus (2-11) gives

$$
\left\|\operatorname{Hess} p\left(x^{\prime}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right)\right\| \leq C_{2}\left|p_{s}\left(x_{0}^{\prime}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right)\right|^{\varepsilon} \quad \text { when }\left|x^{\prime}-x_{0}^{\prime}\right| \leq c
$$

To show (2-12) it suffices to consider an element $b_{j k}\left(x^{\prime}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right)$ of Hess $p$. Clearly $\left|b_{j k}\right| \leq \|$ Hess $p \|$, so by adding $C_{2}\left|p_{s}\left(x_{0}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right)\right|^{\varepsilon}$, we obtain

$$
0 \leq b_{j k}\left(x^{\prime}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right) \leq 2 C_{2}\left|p_{s}\left(x_{0}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right)\right|^{\varepsilon} \quad \text { when }\left|x^{\prime}-x_{0}^{\prime}\right| \leq c
$$

Then we find that

$$
\left|\partial_{x^{\prime}} b_{j k}\left(x_{0}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right)\right| \leq C \sqrt{b_{j k}\left(x_{0}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right)} \leq C^{\prime}\left|p_{s}\left(x_{0}, x^{\prime \prime}, 0, \xi^{\prime \prime}\right)\right|^{\varepsilon / 2}
$$

by [Hörmander 1983, Lemma 7.7.2].
We shall study the microlocal solvability of the operator, which is given by the following definition. Recall that $H_{(s)}^{\text {loc }}(X)$ is the set of distributions that are locally in the $L^{2}$ Sobolev space $H_{(s)}(X)$.

Definition 2.6. If $K \subset S^{*} X$ is a compact set, then we say that $P$ is microlocally solvable at $K$ if there exists an integer $N$ so that for every $f \in H_{(N)}^{\text {loc }}(X)$ there exists $u \in \mathcal{D}^{\prime}(X)$ such that $K \cap \mathrm{WF}(P u-f)=\varnothing$.

Observe that solvability at a compact set $M \subset X$ is equivalent to solvability at $\left.S^{*} X\right|_{M}$ by [Hörmander 1985b, Theorem 26.4.2], and that solvability at a set implies solvability at a subset. Also, by [Hörmander 1985b, Proposition 26.4.4] the microlocal solvability is invariant under conjugation by elliptic Fourier integral operators and multiplication by elliptic pseudodifferential operators. We can now state the main result of the paper.

Theorem 2.7. Assume that $P \in \Psi_{\mathrm{cl}}^{m}(X)$ has principal symbol that vanishes of at least second order at a nonradial involutive manifold $\Sigma_{2}$, is of subprincipal type, does not satisfy condition $\operatorname{Sub}(\Psi)$ on the subprincipal semibicharacteristic $\Gamma \subset \Sigma_{2}$, and satisfies (2-8) near $\Gamma$. In the case the sign change in $\operatorname{Sub}(\Psi)$ is of infinite order, we also assume condition (2-11) near $\Gamma$. Then $P$ is not locally solvable at $\Gamma$.

Example 2.8. Let

$$
\begin{equation*}
P=D_{1} D_{2}+B\left(x, D_{x}\right) \tag{2-13}
\end{equation*}
$$

with $B \in \Psi_{\mathrm{cl}}^{1}$. Then $\sigma(B)$ is the subprincipal symbol on $\Sigma_{2}=\left\{\xi_{1}=\xi_{2}=0\right\}$. Mendoza and Uhlmann [1983] proved that $P$ is not solvable if $\operatorname{Im} \sigma(B)$ changes sign as $x_{1}$ or $x_{2}$ increases on $\Sigma_{2}$, and they proved in [Mendoza and Uhlmann 1984] that $P$ is solvable if $\operatorname{Im} \sigma(B) \neq 0$ on $\Sigma_{2}$. From this it is natural to conjecture that the condition for solvability of $P$ is that $\operatorname{Im} \sigma(B)$ does not change sign on the leaves of $\Sigma$, which are foliated by $\partial_{x_{1}}$ and $\partial_{x_{2}}$. But the following is a counterexample to that conjecture. Let

$$
\begin{equation*}
P=D_{1} D_{2}+D_{t}+i f\left(t, x, D_{x}\right) \tag{2-14}
\end{equation*}
$$

with real and homogeneous $f(t, x, \xi) \in S_{\text {hom }}^{1}$ satisfying $\partial_{x_{j}} f=\mathcal{O}(|f|)$ for $j=1,2$. This operator is of subprincipal type and satisfies (2-8). Then Theorem 2.7 gives that $P$ is not solvable if $t \mapsto f(t, x, \xi)$ changes sign of finite order from - to + , but observe that $f$ has constant sign on the leaves of $\Sigma_{2}$ by Remark 2.3. Thus the solvability of the operator $P$ in (2-13) also depends on the real part of the subprincipal symbol at $\Sigma_{2}$. In fact, with the above conditions one can prove that $D_{1} D_{2}+i f\left(t, x, D_{x}\right)$ is solvable.

Example 2.9. The linearized Navier-Stokes equation

$$
\begin{equation*}
\partial_{t} u+\sum_{j} a_{j}(t, x) \partial_{x_{j}} u+\Delta_{x} u=f, \quad a_{j}(x) \in C^{\infty} \tag{2-15}
\end{equation*}
$$

is of subprincipal type. The symbol is

$$
\begin{equation*}
i \tau+i \sum_{j} a_{j}(t, x) \xi_{j}-|\xi|^{2} \tag{2-16}
\end{equation*}
$$

so the subprincipal symbol is proportional to a real symbol on $\Sigma_{2}=\{\xi=0\}$. Thus condition $\operatorname{Sub}(\Psi)$ is satisfied.

Now let $S^{*} M \subset T^{*} M$ be the cosphere bundle where $|\xi|=1$, and let $\|u\|_{(k)}$ be the $L^{2}$ Sobolev norm of order $k, u \in C_{0}^{\infty}$. In the following, $P^{*}$ will be the $L^{2}$ adjoint of $P$. To prove Theorem 2.7 we shall use the following result.
Remark 2.10. If $P$ is microlocally solvable at $\Gamma \subset S^{*} \mathbb{R}^{n}$, then [Hörmander 1985b, Lemma 26.4.5] gives that for any $Y \Subset \mathbb{R}^{n}$ such that $\Gamma \subset S^{*} Y$, there exists an integer $v$ and a pseudodifferential operator $A$ so that $\mathrm{WF}(A) \cap \Gamma=\varnothing$ and

$$
\begin{equation*}
\|u\|_{(-N)} \leq C\left(\left\|P^{*} u\right\|_{(v)}+\|u\|_{(-N-n)}+\|A u\|_{(0)}\right), \quad u \in C_{0}^{\infty}(Y) \tag{2-17}
\end{equation*}
$$

where $N$ is given by Definition 2.6.

We shall prove Theorem 2.7 in Section 8 by constructing localized approximate solutions to $P^{*} u \cong 0$ and use (2-17) to show that $P$ is not microlocally solvable at $\Gamma$. We shall first find a normal form for the adjoint operator.

## 3. The normal form

Assume that $P^{*}$ has the symbol expansion $p_{m}+p_{m-1}+\cdots$, where $p_{j} \in S_{\text {hom }}^{j}$ is homogeneous of degree $j$. By multiplying $P^{*}$ with an elliptic pseudodifferential operator, we may assume that $m=2$. Choose local symplectic coordinates $(t, x, y, \tau, \xi, \eta)$ so that $\Sigma_{2}=\{\eta=0\}$, which is foliated by leaves spanned by $\partial_{y}$. Since $p_{2}$ vanishes of at least second order at $\Sigma_{2}$, we find that

$$
p_{2}(t, x, y, \tau, \xi, \eta)=\sum_{j k} B_{j k}(t, x, y, \tau, \xi, \eta) \eta_{j} \eta_{k}
$$

where $B_{j k}$ is homogeneous of degree 0 for all $j, k$.
The differential inequality (2-8) in these coordinates means that $\left|\partial_{y} p_{1}\right| \leq C\left|p_{1}\right|$ when $\eta=0$, which by Lemma 2.1 gives that

$$
p_{1}(t, x, y, \tau, \xi, 0)=q(t, x, y, \tau, \xi) r_{1}(t, x, \tau, \xi)
$$

near $\Gamma$, where $q$ is a nonvanishing smooth homogeneous function. By multiplying with pseudodifferential operators with principal symbol equal to $q^{-1}$ on $\Sigma_{2}$, we may assume that $q \equiv 1$ and that $p_{1}$ is constant on the leaves of $\Sigma_{2}$. The Hamilton vector field of $p_{1}$ is then tangent to $\Sigma_{2}$ by (2-7).

We have assumed that $P$ does not satisfy condition $\operatorname{Sub}(\Psi)$ on a semibicharacteristic $\Gamma$ of $p_{1}$ on $\Sigma_{2}$. Since we are now considering the adjoint $P^{*}$ this means that $\operatorname{Im} a p_{1}$ changes sign from + to - on the flow-out $\Gamma$ of $H_{\operatorname{Re} a p_{1}}$ on $\operatorname{Re} a p_{1}^{-1}(0)$ for some $0 \neq a \in C^{\infty}$. By the invariance of condition $\operatorname{Sub}(\Psi)$ given by [Hörmander 1985b, Lemma 26.4.10], it is no restriction to assume that $a$ is homogeneous and constant in $y$. By multiplication with an elliptic pseudodifferential operator having principal symbol $a^{-1}$, we may assume that $a \equiv 1$. Since $\operatorname{Im} p_{1}$ changes sign on $\Gamma$, there is a maximal semibicharacteristic $\Gamma^{\prime} \subset \Gamma$ on which $\operatorname{Im} p_{1}=0$. Here $\Gamma^{\prime}$ could be a point, which is always the case if the sign change is of finite order.

Since $P$ is of subprincipal type, we find that $\partial_{t, x, \tau, \xi} \operatorname{Re} p_{1} \neq 0$ on $\Gamma^{\prime}$ by (2-10), so $\Gamma^{\prime}$ is transversal to the leaves of $\Sigma_{2}$. Since $\left.\operatorname{Im} p_{1}\right|_{\Gamma}$ has opposite signs near the boundary of $\Gamma^{\prime}$, we may shrink $\Gamma$ so that it is not a closed curve. Since $H_{\operatorname{Re} p_{1}}$ is tangent to $\Sigma_{2}$, we can complete $\tau=\operatorname{Re} p_{1}$ to a symplectic coordinate system in a convex neighborhood of $\Gamma^{\prime}$ so that $\eta=0$ on $\Sigma_{2}$, which preserves the leaves. In fact, this is obtained by solving the equation $H_{\tau} \eta=0$ with initial value on a submanifold transversal to $H_{\tau}$. The change of variables can be then done by conjugation with suitable elliptic Fourier integral operators.

Now, using Malgrange's preparation theorem in a neighborhood of $\Gamma^{\prime}$ in $\Sigma_{2}$, we find that

$$
p_{1}(t, x, y, \tau, \xi, 0)=q(t, x, \tau, \xi)(\tau+r(t, x, \xi)), \quad q \neq 0
$$

near $\Gamma$, since $p_{1}$ is constant on the leaves of $\Sigma_{2}$. In fact, on $\Gamma^{\prime}$ we have $p_{1}=0$ and $d p_{1} \neq 0$, so the division can be done locally and by a partition of unity globally near $\Gamma$ after possibly shrinking $\Gamma$. Then
using Taylor's formula on $p_{1}$, we find since $q \neq 0$ that

$$
\begin{equation*}
p_{1}(t, x, y, \tau, \xi, \eta)=q(t, x, \tau, \xi)(\tau+r(t, x, \xi)+A(t, x, y, \tau, \xi, \eta) \cdot \eta) \tag{3-1}
\end{equation*}
$$

By multiplying $P$ with an elliptic pseudodifferential operator, we may again assume $q \equiv 1$. Since $p_{2}$ vanishes of second order at $\Sigma_{2}$, this only changes $A$ with terms which have Hess $p_{2}$ as a factor and terms that vanish at $\Sigma_{2}$.

We can write $r=r_{1}+i r_{2}$ and $A=A_{1}+i A_{2}$ with real-valued $r_{j}$ and $A_{j}, j=1,2$. Now we may complete

$$
\operatorname{Re} p_{1}=\tau+r_{1}(t, x, \xi)+A_{1}(t, x, y, \tau, \xi, \eta) \cdot \eta
$$

to a symplectic coordinate system in a convex neighborhood of $\Gamma^{\prime}$. Since $H_{\operatorname{Re} p_{1}} \in T \Sigma_{2}$ at $\Sigma_{2}$, we may keep $\Sigma_{2}=\{\eta=0\}$, which preserves the leaves of $\Sigma_{2}$ on which $p_{1}$ is constant. Thus, we find that

$$
\begin{equation*}
p_{1}=\tau+i f(t, x, \xi)+i A(t, x, y, \tau, \xi, \eta) \cdot \eta \tag{3-2}
\end{equation*}
$$

where $f=r_{2}$ and $A=A_{2}$ are real-valued. We also find that

$$
\begin{equation*}
\Gamma=\left\{\left(t, x_{0}, y_{0}, 0, \xi_{0}, 0\right)\right\}, \quad t \in I \tag{3-3}
\end{equation*}
$$

where $I$ is an interval in $\mathbb{R}$. The symplectic change of coordinates can be made by conjugation with elliptic Fourier integral operators, which only changes $A$ with terms having Hess $p_{2}$ as a factor and terms that vanish at $\Sigma_{2}$. Observe that $A$ need not be real-valued after these changes.

We have assumed that condition $\operatorname{Sub}(\Psi)$ is not satisfied for $P$ on the subprincipal semibicharacteristic $\Gamma$. Thus the imaginary part of the subprincipal symbol of $P^{*}$ on $\Sigma_{2}$

$$
\begin{equation*}
t \mapsto f\left(t, x_{0}, \xi_{0}\right) \tag{3-4}
\end{equation*}
$$

changes sign from + to - as $t$ increases on $I \subset \mathbb{R}$. Similarly, we have $f=0$ on $\Gamma^{\prime}$, where $\Gamma^{\prime}$ is given by (3-3) with $I$ replaced by $I^{\prime} \subset I$. By reducing to minimal bicharacteristics on which $t \mapsto f(t, x, \xi)$ changes sign as in [Hörmander 1981, p. 75], we may assume that $f$ vanishes of infinite order on a bicharacteristic $\Gamma^{\prime}$ arbitrarily close to the original bicharacteristic, if $\Gamma^{\prime}$ is not a point (see [Wittsten 2012, Section 2] for a more refined analysis). If $\Gamma^{\prime}$ is not a point then it is a one-dimensional bicharacteristic by [Hörmander 1981, Definition 3.5], which means that the Hamilton vector field on $\Gamma^{\prime}$ is proportional to a real vector.

In fact, if $f\left(a, x_{0}, \xi_{0}\right)>0>f\left(b, x_{0}, \xi_{0}\right)$ for some $a<b$, then we can define

$$
L(x, \xi)=\inf \{t-s: a<s<t<b \text { such that } f(s, x, \xi)>0>f(t, x, \xi)\}
$$

when $(x, \xi)$ is close to $\left(x_{0}, \xi_{0}\right)$, and we put $L_{0}=\liminf _{(x, \xi) \rightarrow\left(x_{0}, \xi_{0}\right)} L(x, \xi)$. Then for every $\varepsilon>0$ there exists an open neighborhood $V_{\varepsilon}$ of $\left(x_{0}, \xi_{0}\right)$ such that the diameter of $V_{\varepsilon}$ is less than $\varepsilon$ and $L(x, \xi)>L_{0}-\varepsilon / 2$ when $(x, \xi) \in V_{\varepsilon}$. By definition, there exists $\left(x_{\varepsilon}, \xi_{\varepsilon}\right) \in V_{\varepsilon}$ and $a<s_{\varepsilon}<t_{\varepsilon}<b$ so that $t_{\varepsilon}-s_{\varepsilon}<L_{0}+\varepsilon / 2$ and $f\left(s_{\varepsilon}, x_{\varepsilon}, \xi_{\varepsilon}\right)>0>f\left(t_{\varepsilon}, x_{\varepsilon}, \xi_{\varepsilon}\right)$. Then it is easy to see that

$$
\begin{equation*}
\partial_{x}^{\alpha} \partial_{\xi}^{\beta} f\left(t, x_{\varepsilon}, \xi_{\varepsilon}\right)=0 \quad \forall \alpha, \beta \text { when } s_{\varepsilon}+\varepsilon<t<t_{\varepsilon}-\varepsilon \tag{3-5}
\end{equation*}
$$

since else we would have a sign change in an interval with length less than $L_{0}-\varepsilon / 2$ in $V_{\varepsilon}$. We may choose a sequence $\varepsilon_{j} \rightarrow 0$ so that $s_{\varepsilon_{j}} \rightarrow s_{0}$ and $t_{\varepsilon_{j}} \rightarrow t_{0}$. Then $L_{0}=t_{0}-s_{0}$ and (3-5) holds at ( $x_{0}, \xi_{0}$ ) for $s_{0}<t<t_{0}$.

We also obtain the following condition from (2-11).
Remark 3.1. If the sign change of $t \mapsto f(t, x, \xi)$ is of infinite order on $\Gamma$, then we find from assumption (2-11) that

$$
\begin{equation*}
\left\|\left\{B_{j k}\right\}_{j k}\right\|+|A|+|d f| \lesssim|f|^{\varepsilon} \quad \text { near } \Gamma \text { on } \Sigma_{2} \tag{3-6}
\end{equation*}
$$

for some $\varepsilon>0$. Here $a \lesssim b$ (and $b \gtrsim a$ ) means that $a \leq C b$ for some $C>0$.
In fact, terms having Hess $\left.p_{2}\right|_{\Sigma_{2}}=\left\{B_{j k}\right\}_{j k}$ as a factor can be estimated by (2-11), so we may assume that (3-2) holds with real $A$. The subprincipal symbol is equal to $p_{s}=p_{1}+i \sum_{j k} \partial_{y_{j}} B_{j k} \eta_{k}$ modulo terms that are $\mathcal{O}\left(|\eta|^{2}\right)$, so $p_{s}=p_{1}$ on $\Sigma_{2}$. By Remark 2.5 and (2-8), we can estimate the terms $\partial_{y_{j}} B_{j k} d \eta_{k}$ in $d p_{s}$ by replacing $\varepsilon$ with $\varepsilon / 2$ in (2-11), so we may replace $p_{s}$ by $p_{1}$ in the estimate. Let $0 \neq a=a_{1}+i a_{2}$ with real-valued $a_{j}$ in (2-11) so that $\left.d \operatorname{Re} a p_{1}\right|_{T \Sigma_{2}} \neq 0$. We have $d p_{1}=d \tau+i(d f+A d \eta)$ on $\Sigma_{2}$, so

$$
\left|d p_{1} \wedge d \overline{p_{1}}\right| \cong|d f|+|A| \quad \text { on } \Sigma_{2}
$$

Thus we find from (2-11) that $|d f|+|A|=0$ on $\Gamma^{\prime}$. Since $\left.d \operatorname{Re} a p_{1}\right|_{T \Sigma_{2}} \neq 0$, we find that $a_{1} \neq 0$ on $\Gamma^{\prime}$. On $\Sigma_{2}$ we have $\operatorname{Re} a p_{1}=a_{1} \tau-a_{2} f=0$ when $\tau=a_{2} f / a_{1}$. We obtain

$$
\operatorname{Im} a p_{1}=a_{2} \tau+a_{1} f=|a|^{2} f / a_{1} \quad \text { when } \operatorname{Re} a p_{1}=0 \text { on } \Sigma_{2} \text { near } \Gamma^{\prime}
$$

which gives (3-6) from (2-11).
We obtain the following normal form for these operators of subprincipal type:

$$
\begin{equation*}
P^{*}=D_{t}+F\left(t, x, y, D_{t}, D_{x}, D_{y}\right) \tag{3-7}
\end{equation*}
$$

where $F \sim F_{2}+F_{1}+\cdots$ with homogeneous $F_{j} \in C^{\infty}\left(\mathbb{R}, S_{\text {hom }}^{j}\right)$. Here $F_{2}$ vanishes of at least second order on $\Sigma_{2}=\{\eta=0\}$, so we find by Taylor's formula that

$$
\begin{equation*}
F_{2}(t, x, y, \tau, \xi, \eta)=B(t, x, y, \tau, \xi, \eta)=\sum_{j k} B_{j k}(t, x, y, \tau, \xi, \eta) \eta_{j} \eta_{k} \tag{3-8}
\end{equation*}
$$

with homogeneous $B_{j k}$. Then $\left.\left\{B_{j k}\right\}_{j k}\right|_{\Sigma_{2}}=\operatorname{Hess} F_{2}(t, x, y, \tau, \xi, 0)$. Also we have that $F_{1}$ vanishes on the semibicharacteristic $\Gamma^{\prime}$ and

$$
\begin{equation*}
F_{1}(t, x, y, \tau, \xi, \eta)=i f(t, x, \xi)+A(t, x, y, \tau, \xi, \eta) \cdot \eta \tag{3-9}
\end{equation*}
$$

Here $f$ is real and homogeneous of degree 1 and $\left.A\right|_{\Sigma_{2}}=\left.\partial_{\eta} F_{1}\right|_{\Sigma_{2}}$. We have that the principal symbol $\sigma\left(P^{*}\right)$ is equal to $F_{2}$, and the subprincipal symbol $\sigma_{\text {sub }}\left(P^{*}\right)$ is equal to $\tau+i f$ on $\Sigma_{2}$. Thus we obtain the following result.

Proposition 3.2. Assume that P satisfies the conditions in Theorem 2.7. Then by conjugation with elliptic Fourier integral operators and multiplication with an elliptic pseudodifferential operator, we may assume

$$
\begin{equation*}
P^{*}=D_{t}+F\left(t, x, y, D_{t}, D_{x}, D_{y}\right) \tag{3-10}
\end{equation*}
$$

microlocally near $\Gamma=\left\{\left(t, x_{0}, y_{0}, 0, \xi_{0}, 0\right): t \in I\right\} \subset \Sigma_{2}$, where $S_{\mathrm{cl}}^{2} \ni F \cong F_{2}+F_{1}+\cdots$ with $F_{j} \in S_{\text {hom }}^{j}$ is homogeneous of degree $j$ and

$$
F_{2}(t, x, y, \tau, \xi, \eta)=\sum_{j k} B_{j k}(t, x, y, \tau, \xi, \eta) \eta_{j} \eta_{k} \in S_{\mathrm{hom}}^{2}
$$

vanishes of second order on $\Sigma_{2}$. We may also assume

$$
F_{1}(t, x, y, \tau, \xi, \eta)=i f(t, x, \xi)+A(t, x, y, \tau, \xi, \eta) \cdot \eta
$$

is homogeneous of degree 1 and $f$ is real-valued such that $t \mapsto f\left(t, x_{0}, \xi_{0}\right)$ changes sign from + to - as $t$ increases on $I \subset \mathbb{R}$. If $f\left(t, x_{0}, \xi_{0}\right)=0$ on a subinterval $I^{\prime} \subseteq I$ such that $\left|I^{\prime}\right| \neq 0$, then we may assume that $\partial_{t}^{k} \partial_{x}^{\alpha} \partial_{\xi}^{\beta} f\left(t, x_{0}, \xi_{0}\right)=0$ for all $k, \alpha, \beta$, for $t \in I^{\prime}$. If the sign change of $f$ is of infinite order then (3-6) is satisfied near $\Gamma$.

For the proof of Theorem 2.7, we shall modify the Moyer-Hörmander construction of approximate solutions of the type

$$
\begin{equation*}
u_{\lambda}(t, x, y)=e^{i \lambda \omega(t, x, y)} \sum_{j \geq 0} \phi_{j}(t, x, y) \lambda^{-j / N}, \quad \lambda \geq 1 \tag{3-11}
\end{equation*}
$$

with $N$ to be determined later. Here the phase function $\omega(t, x)$ will be complex-valued, but $\operatorname{Im} \omega \geq 0$ and $\partial \operatorname{Re} \omega \neq 0$ when $\operatorname{Im} \omega=0$. Letting $z=(t, x, y)$, we therefore have the formal expansion

$$
\begin{equation*}
p(z, D)(\exp (i \lambda \omega) \phi) \sim \exp (i \lambda \omega) \sum_{\alpha} \partial_{\zeta}^{\alpha} p\left(z, \lambda \partial_{z} \omega(z)\right) \mathcal{R}_{\alpha}(\omega, \lambda, D) \phi(z) / \alpha! \tag{3-12}
\end{equation*}
$$

where $\mathcal{R}_{\alpha}(\omega, \lambda, D) \phi(z)=\left.D_{w}^{\alpha}(\exp (i \lambda \tilde{\omega}(z, w)) \phi(w))\right|_{w=z}$ and

$$
\tilde{\omega}(z, w)=\omega(w)-\omega(z)+(z-w) \partial \omega(z)
$$

Observe that the values of the symbol are given by an almost analytic extension; see Theorem 3.1 in Chapter VI and Chapter X:4 in [Trèves 1980]. This gives

$$
\begin{align*}
e^{-i \lambda \omega} P^{*} e^{i \lambda \omega} \phi= & \left(\lambda \partial_{t} \omega+\lambda^{2} B\left(t, x, y, \partial_{t, x, y} \omega\right)+i \lambda f\left(t, x, \partial_{x} \omega\right)-\lambda \partial_{\eta}^{2} B\left(t, x, y, \partial_{t, x, y} \omega\right) \partial_{y}^{2} \omega / 2\right) \phi \\
+ & D_{t} \phi+\lambda \partial_{\eta} B\left(t, x, y, \partial_{t, x, y} \omega\right) D_{y} \phi+\partial_{\eta}^{2} B\left(t, x, y, \partial_{t, x, y} \omega\right) D_{y}^{2} \phi / 2 \\
& +i \partial_{\xi} f\left(t, x, \partial_{x} \omega\right) D_{x} \phi+A\left(t, x, y, \partial_{t, x, y} \omega\right) D_{y} \phi+\sum_{j \geq 0} \lambda^{-j} R_{j}\left(t, x, y, D_{t, x, y}\right) \phi \tag{3-13}
\end{align*}
$$

where $R_{0}(t, x, y)=F_{0}\left(t, x, y, \partial_{t, x, y} \omega\right)$. Here the values of the symbols at $\left(t, x, y, \partial_{t, x, y} \omega\right)$ will be replaced by finite Taylor expansions at $\left(t, x, y, \partial_{t, x, y} \operatorname{Re} \omega\right)$. In fact, the almost analytic extensions are determined by these Taylor expansions.

Because of the inhomogeneity coming from the terms of $B$, we shall use a phase function $\omega(t, x)$ which is constant in $y$ so that

$$
\begin{equation*}
u_{\lambda}(t, x, y)=e^{i \lambda \omega(t, x)} \sum_{j \geq 0} \phi_{j}(t, x, y) \lambda^{-j / N}, \quad \lambda \geq 1 \tag{3-14}
\end{equation*}
$$

When $\partial_{y} \omega \equiv 0$ the expansion (3-13) becomes

$$
\begin{align*}
e^{-i \lambda \omega} P^{*} e^{i \lambda \omega} \phi= & \lambda\left(\partial_{t} \omega+i f\left(t, x, \partial_{x} \omega\right)\right) \phi+D_{t} \phi+\partial_{\eta}^{2} B\left(t, x, y, \partial_{t, x} \omega, 0\right) D_{y}^{2} \phi / 2 \\
& +A\left(t, x, y, \partial_{t, x} \omega, 0\right) D_{y} \phi+i \partial_{\xi} f\left(t, x, \partial_{x} \omega\right) D_{x} \phi+\sum_{j \geq 0} \lambda^{-j} R_{j}\left(t, x, y, D_{t, x, y}\right) \phi \tag{3-15}
\end{align*}
$$

where $R_{0}(t, x, y)=F_{0}\left(t, x, y, \partial_{t, x} \omega, 0\right)$, and $R_{m}\left(t, x, y, D_{t, x, y}\right)$ are differential operators of order $j$ in $t$, order $k$ in $x$ and order $\ell$ in $y$, where $j+k+\ell \leq m+2$ for $m>0$. In fact, this follows since $\partial_{\tau}^{j} \partial_{\xi}^{\alpha} \partial_{\eta}^{\beta} F_{k} \in S^{k-j-|\alpha|-|\beta|}$ by homogeneity.

## 4. The eikonal equation

We shall first solve the eikonal equation approximately, which is given by the highest-order term of (3-15)

$$
\begin{equation*}
\partial_{t} \omega+i f\left(t, x, \partial_{x} \omega\right)=0 \tag{4-1}
\end{equation*}
$$

where $t \mapsto f(t, x, \xi)$ changes sign from + to - for some $(x, \xi)$ as $t$ increases in a neighborhood of $\Gamma=\left\{\left(t, x_{0}, \xi_{0}\right): t \in I\right\}$ on which $f(t, x, \xi)$ vanishes. If $|I| \neq 0$ then by reducing to minimal bicharacteristics as in Section 3, we may assume that $f$ vanishes of infinite order at $\Gamma$. We shall choose the phase function so that $\operatorname{Im} \omega \geq 0$ and $\partial_{x}^{2} \operatorname{Im} \omega>0$ near the interval. By changing coordinates, it is no restriction to assume $0 \in I$. We shall use the approach by Hörmander [1981] in the principal-type case and use the phase function to localize in $t$ and $x$. Observe that since $\omega$ does not depend on $y$, the localization in the $y$-variables will be done in the amplitude $\phi$.

We shall take the Taylor expansion of $\omega$ in $x$ :

$$
\begin{equation*}
\omega(t, x)=w_{0}(t)+\left\langle x-x_{0}(t), \xi_{0}(t)\right\rangle+\sum_{2 \leq|\alpha| \leq K} w_{\alpha}(t)\left(x-x_{0}(t)\right)^{\alpha} / \alpha!. \tag{4-2}
\end{equation*}
$$

Here $\alpha=\left(\alpha_{1}, \alpha_{2}, \ldots\right)$, with $\alpha_{j} \in \mathbb{N}, \alpha!=\prod_{j} \alpha_{j}!$ and $|\alpha|=\alpha_{1}+\alpha_{2}+\cdots$. Then we find that

$$
\begin{align*}
\partial_{t} \omega(t, x)=w_{0}^{\prime}(t)- & \left\langle x_{0}^{\prime}(t), \xi_{0}(t)\right\rangle+\left\langle x-x_{0}(t), \xi_{0}^{\prime}(t)\right\rangle \\
& +\sum_{2 \leq|\alpha| \leq K} w_{\alpha}^{\prime}(t)\left(x-x_{0}(t)\right)^{\alpha} / \alpha!-\sum_{1 \leq|\alpha| \leq K-1}^{k} w_{\alpha+e_{k}}(t)\left(x-x_{0}(t)\right)^{\alpha} x_{0, k}^{\prime}(t) / \alpha!, \tag{4-3}
\end{align*}
$$

where $e_{k}=(0, \ldots, 0,1,0, \ldots, 0)$ is the $k$-th unit vector. We also find

$$
\begin{equation*}
\partial_{x_{j}} \omega(t, x)=\xi_{0 j}(t)+\sum_{1 \leq|\alpha| \leq K-1} w_{\alpha+e_{j}}(t)\left(x-x_{0}(t)\right)^{\alpha} / \alpha!=\xi_{0 j}(t)+\sigma_{j}(t, x) \tag{4-4}
\end{equation*}
$$

Here $\xi_{0}(t)=\left(\xi_{0,1}(t), \ldots\right)$ and $\sigma=\left\{\sigma_{j}\right\}_{j}$ is a finite expansion in powers of $\Delta x=x-x_{0}$. We define the value of $f\left(t, x, \partial_{x} \omega\right)$ by the Taylor expansion

$$
\begin{align*}
f\left(t, x, \partial_{x} \omega\right)=f(t, x, & \left.\xi_{0}+\sigma\right) \\
& =f\left(t, x, \xi_{0}\right)+\sum_{j} \partial_{\xi_{j}} f\left(t, x, \xi_{0}\right) \sigma_{j}+\sum_{j k} \partial_{\xi_{j}} \partial_{\xi_{k}} f\left(t, x, \xi_{0}\right) \sigma_{j} \sigma_{k} / 2+\cdots \tag{4-5}
\end{align*}
$$

Now the value at $x=x_{0}$ of (4-1) is equal to $w_{0}^{\prime}(t)-\left\langle x_{0}^{\prime}(t), \xi_{0}(t)\right\rangle+i f\left(t, x_{0}(t), \xi_{0}(t)\right)$. This vanishes if

$$
\left\{\begin{array}{l}
\operatorname{Re} w_{0}^{\prime}(t)=\left\langle x_{0}^{\prime}(t), \xi_{0}(t)\right\rangle  \tag{4-6}\\
\operatorname{Im} w_{0}^{\prime}(t)=-f\left(t, x_{0}(t), \xi_{0}(t)\right)
\end{array}\right.
$$

so by putting $w_{0}(0)=0$, this will determine $w_{0}$ once we have $\left(x_{0}(t), \xi_{0}(t)\right)$.
We shall simplify the notation and put $w_{k}=\left\{w_{\alpha} k!/ \alpha!\right\}_{|\alpha|=k}$ so that $w_{k}$ is a multilinear form. The first-order terms in $x-x_{0}$ of (4-1) vanish if

$$
\xi_{0}^{\prime}(t)-w_{2}(t) x_{0}^{\prime}(t)+i\left(\partial_{x} f\left(t, x_{0}(t), \xi_{0}(t)\right)+\partial_{\xi} f\left(t, x_{0}(t), \xi_{0}(t)\right) w_{2}(t)\right)=0
$$

We find by taking real and imaginary parts that

$$
\left\{\begin{array}{l}
\xi_{0}^{\prime}=\operatorname{Re} w_{2} x_{0}^{\prime}+\partial_{\xi} f\left(t, x_{0}, \xi_{0}\right) \operatorname{Im} w_{2}  \tag{4-7}\\
x_{0}^{\prime}=\left(\operatorname{Im} w_{2}\right)^{-1}\left(\partial_{x} f\left(t, x_{0}, \xi_{0}\right)+\partial_{\xi} f\left(t, x_{0}, \xi_{0}\right) \operatorname{Re} w_{2}\right)
\end{array}\right.
$$

with $\left(x_{0}(0), \xi_{0}(0)\right)=\left(x_{0}, \xi_{0}\right)$, which will determine $x_{0}(t)$ and $\xi_{0}(t)$ if $\left|\operatorname{Im} w_{2}\right| \neq 0$.
The second-order terms in $x-x_{0}$ vanish if

$$
w_{2}^{\prime} / 2-w_{3} x_{0}^{\prime} / 2+i\left(\partial_{\xi} f w_{3} / 2+\partial_{x}^{2} f / 2+\partial_{x} \partial_{\xi} f w_{2}+w_{2} \partial_{\xi}^{2} f w_{2} / 2\right)=0
$$

which gives

$$
\begin{equation*}
w_{2}^{\prime}=w_{3} x_{0}^{\prime}-i\left(\partial_{\xi} f w_{3}+\partial_{x}^{2} f+2 \partial_{x} \partial_{\xi} f w_{2}+w_{2} \partial_{\xi}^{2} f w_{2}\right) \tag{4-8}
\end{equation*}
$$

with initial data $w_{2}(0)$ such that $\operatorname{Im} w_{2}(0)>0$.
We find that the terms of order $k>2$ vanish if

$$
\begin{equation*}
w_{k}^{\prime}-w_{k+1} x_{0}^{\prime}=F_{k}\left(t, x_{0}, \xi_{0},\left\{w_{j}\right\}\right) \tag{4-9}
\end{equation*}
$$

where we may choose $w_{k}(0)=0$. Here $F_{k}$ is a linear combination of the derivatives of $f$ of order $\leq k$ multiplied by polynomials in $w_{j}$ with $2 \leq j \leq k+1$. When $k=K$ we get $w_{K}^{\prime}=F_{K}\left(t, x_{0}, \xi_{0},\left\{w_{j}\right\}\right)$, where $j \leq K$. The equations (4-7)-(4-9) form a quasilinear system of differential equations, which can be solved in a convex neighborhood of 0 . In the case when $|I| \neq 0$, we have assumed that $\partial_{t, x, \xi}^{\alpha} f\left(t, x_{0}, \xi_{0}\right) \equiv 0$ for all $\alpha$, for $t \in I$. Then we find from (4-7)-(4-9) that $x_{0}, \xi_{0}$ and $w_{k}$ are constant in $t \in I$, so we may solve (4-7)-(4-9) in a convex neighborhood of $I$. Observe that the higher-order terms cannot change the condition that $\operatorname{Im} \partial_{x}^{2} \omega \geq c>0$ and $\operatorname{Im} \omega(t, x) \geq 0$ if $\left|x-x_{0}(t)\right| \ll 1$. Summing up, we have proved the following result.
Proposition 4.1. Let $\Gamma=\left\{\left(t, x_{0}, \xi_{0}\right): t \in I\right\}$ and assume that $\partial_{t}^{k} \partial_{x}^{\alpha} \partial_{\xi}^{\beta} f\left(t, x_{0}, \xi_{0}\right)=0$ for all $t \in I$ in the case $|I| \neq 0$. Then we may solve (4-1) with $\omega(t, x)$ given by (4-2) in a convex neighborhood $\Omega$ of $\Gamma$ modulo $\mathcal{O}\left(\left|x-x_{0}(t)\right|^{M}\right)$ for any $M$ such that $\left(x_{0}(t), \xi_{0}(t)\right)=\left(x_{0}, \xi_{0}\right)$ when $t \in I$ and $w_{k}(t) \in C^{\infty}$ such that $w_{0}(t)=0, \operatorname{Im} w_{2}(t)>0$ and $w_{k}(t)=0, k>2$, when $t \in I$.

Then we obtain $\operatorname{Im} \omega(t, x) \geq c\left|x-x_{0}(t)\right|^{2}$ near $\Gamma$, with $c>0$, so the errors that are $\mathcal{O}\left(\left|x-x_{0}\right|^{M}\right)$ in the eikonal equation will give terms that are bounded by $C_{M} \lambda^{-M / 2}$. But we have to show that $t \mapsto f\left(t, x_{0}(t), \xi_{0}(t)\right)$ also changes sign from + to - as $t$ increases for some choice of $\left(x_{0}, \xi_{0}\right)$. This
problem will be studied in the next section, with a special emphasis on the finite vanishing case. By (4-6) we then obtain that $t \mapsto \operatorname{Im} w_{0}(t)$ has a local minimum on $I$ which can be equal to 0 by subtracting a constant.

## 5. The change of sign

We have assumed that condition $\operatorname{Sub}(\Psi)$ for $P$ is not satisfied near the subprincipal semibicharacteristic $\Gamma=\left\{\left(t, x_{0}, \xi_{0}\right): t \in I\right\}$, so that $t \mapsto f(t, x, \xi)$ changes sign from + to - for some $(x, \xi)$ as $t$ increases near $\Gamma$. But after solving the eikonal equation, we have to know that $t \mapsto f\left(t, x_{0}(t), \xi_{0}(t)\right)$ has the same sign change, possibly after changing the starting point $\left(x_{0}, \xi_{0}\right)$. In order to do so, we shall use the invariance of condition $\operatorname{Sub}(\Psi)$, but note that condition (3-6) is only assumed when the change of sign is of infinite order. Therefore we shall first consider the case when the sign change is of finite order and show that this condition is preserved after solving the eikonal equation. Thus assume that

$$
\begin{equation*}
\partial_{t}^{k} f\left(t_{0}, x_{0}, \xi_{0}\right)<0 \quad \text { and } \quad \partial_{t}^{j} f\left(t_{0}, x_{0}, \xi_{0}\right)=0 \quad \text { for } j<k \tag{5-1}
\end{equation*}
$$

for some odd integer $k$, where we may assume $t_{0}=0$. Now, if the order of the zero is not constant in a neighborhood of ( $x_{0}, \xi_{0}$ ) then in any neighborhood the mapping $t \mapsto f(t, x, \xi)$ must have a zero of odd order with sign change from + to - , and the order of vanishing is constant almost everywhere on $f^{-1}(0)$. We obtain this because $\partial_{t}^{k} f \neq 0, t \mapsto f(t, x, \xi)$ goes from + to - and the set where the order of the zero changes is nowhere dense in $f^{-1}(0)$ since it is the union of boundaries of closed sets in the relative topology. By possibly changing ( $t_{0}, x_{0}, \xi_{0}$ ), we may assume that (5-1) holds with $t_{0}=0$, and that the order of the zero is odd and constant near $\left(x_{0}, \xi_{0}\right)$. Then the zeros form a smooth manifold by the implicit function theorem. Using Taylor's formula, we find that $f(t, w)=a(t, w)\left(t-t_{0}(w)\right)^{k}$, where $k \geq 1$ is odd, $w=(x, \xi), t_{0}\left(w_{0}\right)=0$ and $a<0$ in a neighborhood of $\left(0, w_{0}\right)=\left(0, x_{0}, \xi_{0}\right)$. Then we find

$$
\begin{equation*}
\partial_{w} f=\partial_{w} a\left(t-t_{0}\right)^{k}-a k\left(t-t_{0}\right)^{k-1} \partial_{w} t_{0} \tag{5-2}
\end{equation*}
$$

which vanishes of at least order $k-1$ in $t$ at $f^{-1}(0)$. Let $w(t)=\left(x_{0}(t), \xi_{0}(t)\right)$. Then

$$
f(t, w(t))=f\left(t, w_{0}\right)+\partial_{w} f\left(t, w_{0}\right) \Delta w(t)+\mathcal{O}\left(|\Delta w(t)|^{2}\right)
$$

where $\Delta w(t)=w(t)-w_{0}$. Now $t \mapsto f\left(t, w_{0}\right)$ vanishes of order $k$ in $t$ at 0 and $t \mapsto \partial_{w} f\left(t, w_{0}\right)$ vanishes of at least order $k-1$, so if $t \mapsto \Delta w(t)$ vanishes of at least order $k>1$ then by (5-2) we find that $t \mapsto f(t, w(t))$ vanishes of order $k$. Since $(d / d t) \Delta w(t)=w^{\prime}(t)$, we will need the following result.

Lemma 5.1. Let $\left(x_{0}(t), \xi_{0}(t)\right)$ be the solution to equation (4-7) with $\operatorname{Im} w_{2}(0) \neq 0$ and assume that $t \mapsto \partial_{w} f\left(t, x_{0}, \xi_{0}\right)$ vanishes of order $r \geq 1$ at $t=0$. Then $\left(x_{0}^{\prime}(t), \xi_{0}^{\prime}(t)\right)$ vanishes of order $r$ and $\Delta w(t)$ vanishes of order $r+1$ at $t=0$.

Proof. By (4-7) we have

$$
\begin{equation*}
w^{\prime}(t)=\left(x_{0}^{\prime}(t), \xi_{0}^{\prime}(t)\right)=A(t) \partial_{w} f(t, w(t)), \quad w(0)=w_{0} \tag{5-3}
\end{equation*}
$$

Here we have $|A(0)| \neq 0$ if $\operatorname{Im} w_{2}(0) \neq 0$; in fact $w^{\prime}(0)=0$ then gives $\partial_{\xi} f\left(0, w_{0}\right)=0$ and $\partial_{x} f\left(0, w_{0}\right)=0$ by (4-7).

Now we define $\phi_{0}(t)=\partial_{w} f\left(t, w_{0}\right)$ and $\phi_{1}(t)=\partial_{w} f(t, w(t))$. Then we have $w^{\prime}(t)=A(t) \phi_{1}(t)$ and the condition is that $\phi_{0}(t)$ vanishes of order $r \geq 1$ at 0 . We shall proceed by induction, and first assume that $r=1$. Since $w(0)=w_{0}$ we find $\phi_{1}(0)=\phi_{0}(0)=0$ and thus $w^{\prime}(0)=0$.

Next, for $r>1$ we assume by induction that $w^{\prime}(t)$ vanishes of order $r-1$ at 0 so $w^{(k)}(0)=0$ for $k<r$, and then we shall show that $w^{(r)}(0)=0$ so that $w^{\prime}$ vanishes of order $r$. Using the chain rule we obtain

$$
\begin{equation*}
\partial_{t}^{r}(g(t, w(t)))=\sum_{\substack{0 \leq j \leq r \\ \sum_{i} r_{i}+j=r}} c_{j, \alpha} \partial_{t}^{j} \partial_{w}^{\alpha} g(t, w(t)) \prod_{i=1}^{|\alpha|} w^{\left(r_{i}\right)}(t) \tag{5-4}
\end{equation*}
$$

for any $g(t, w) \in C^{\infty}$. Thus, for $g=\partial_{w} f$ we find that

$$
\phi_{1}^{(k)}(0)=\phi_{0}^{(k)}(0)+\partial_{t}^{k-1} \partial_{w}^{2} f\left(0, x_{0}, \xi_{0}\right) w^{\prime}(0)+\cdots+\partial_{w}^{2} f\left(0, x_{0}, \xi_{0}\right) w^{(k)}(0)=\phi_{0}^{(k)}(0)=0
$$

for $k<r$, since the other terms have some factor $w^{(j)}(0)=0, j \leq k$, which implies that $\phi_{1}(t)$ vanishes of order $r$. Since $w^{\prime}=A \phi_{1}$ we find that $w^{\prime}(t)$ vanishes of order $r$, which gives the induction step and the proof.

Now, if $f\left(t, w_{0}\right)$ vanishes of order $k$ then $\partial_{w} f\left(t, w_{0}\right)$ vanishes of order $k-1$. Thus $w^{\prime}(t)$ vanishes of order $k-1$ by Lemma 5.1, and since $w(0)=w_{0}$ we find that $\Delta w(t)$ vanishes of order $k$. Thus, we find that $f(t, w(t))-f\left(t, w_{0}\right)$ vanishes of order $2 k-1$, so these terms vanish of same order if $k>1$. In the case $k=1$, we shall use an argument of Hörmander [1981] for the principal-type case. We obtain from (4-6) that $\partial_{t}(f(t, w(t)))=-\operatorname{Im} w_{0}^{\prime \prime}(t)$; thus

$$
\begin{equation*}
\operatorname{Im} w_{0}^{\prime \prime}(0)=-\partial_{t} f\left(0, w_{0}\right)-\partial_{\xi} f\left(0, w_{0}\right) \cdot \xi_{0}^{\prime}-\partial_{x} f\left(0, w_{0}\right) \cdot x_{0}^{\prime} \tag{5-5}
\end{equation*}
$$

where $\partial_{t} f\left(0, w_{0}\right)=-c<0$. We find from (4-7) that

$$
\left\{\begin{array}{l}
\xi_{0}^{\prime}(0)=\operatorname{Re} w_{2}(0) x_{0}^{\prime}(0)+\partial_{\xi} f\left(0, w_{0}\right) \operatorname{Im} w_{2}(0)  \tag{5-6}\\
x_{0}^{\prime}(0)=\left(\operatorname{Im} w_{2}(0)\right)^{-1}\left(\partial_{x} f\left(0, w_{0}\right)+\partial_{\xi} f\left(0, w_{0}\right) \operatorname{Re} w_{2}(0)\right)
\end{array}\right.
$$

If $\partial_{\xi} f\left(0, w_{0}\right)=0$ then we find that $x_{0}^{\prime}(0)=\left(\operatorname{Im} w_{2}(0)\right)^{-1} \partial_{x} f\left(0, w_{0}\right)$ and obtain

$$
\begin{equation*}
\operatorname{Im} w_{0}^{\prime \prime}(0)=c-\partial_{x} f\left(0, w_{0}\right)\left(\operatorname{Im} w_{2}(0)\right)^{-1} \partial_{x} f\left(0, w_{0}\right)>c / 2>0 \tag{5-7}
\end{equation*}
$$

by choosing $\operatorname{Im} w_{2}(0)=\kappa \operatorname{Id}$ with $\kappa \gg 1$. If $\partial_{\xi} f\left(0, w_{0}\right) \neq 0$ then we may choose $\operatorname{Re} w_{2}(0)$ so that

$$
\begin{equation*}
\partial_{x} f\left(0, w_{0}\right)+\partial_{\xi} f\left(0, w_{0}\right) \operatorname{Re} w_{2}(0)=0 \tag{5-8}
\end{equation*}
$$

Then we find $x_{0}^{\prime}(0)=0$ and we obtain

$$
\begin{equation*}
\operatorname{Im} w_{0}^{\prime \prime}(0)=c-\partial_{\xi} f\left(0, w_{0}\right) \operatorname{Im} w_{2}(0) \partial_{\xi} f\left(0, w_{0}\right)>c / 2>0 \tag{5-9}
\end{equation*}
$$

by choosing $\operatorname{Im} w_{2}(0)=\kappa \operatorname{Id}$ with $0<\kappa \ll 1$. Thus in both cases we find that $\partial_{t} f(t, w(t))=\operatorname{Im} w_{0}^{\prime}(t)<0$ at $t=0$.

We find that $t \mapsto f(t, w(t))$ changes sign from + to - of order $k$ as $t$ increases at $t=0$. We may then rewrite the equation as

$$
\begin{equation*}
\operatorname{Im} w_{0}^{\prime}(t)=t^{k} c(t) \tag{5-10}
\end{equation*}
$$

where $c(t)>0$ in a neighborhood of the origin. Since $\operatorname{Im} w_{2}(0)>0$ we find that

$$
\begin{equation*}
e^{i \lambda \omega(t, x)} \leq e^{-c_{0} \lambda\left(t^{k+1}+\left|x-x_{0}\right|^{2}\right)}, \quad\left|x-x_{0}\right| \ll 1, \quad|t| \ll 1, \quad c_{0}>0 . \tag{5-11}
\end{equation*}
$$

Thus the errors that are $\mathcal{O}\left(\left|x-x_{0}\right|^{M}\right)$ in the eikonal equation will give terms that are bounded by $C_{M} \lambda^{-M / 2}$.
We shall also consider the case when $t \mapsto f(t, x, \xi)$ changes sign from + to - of infinite order near $\Gamma$. If $\Gamma$ is not a point, then by reducing to a minimal bicharacteristic as in Section 3, we may assume that $f(t, x, \xi)$ vanishes of infinite order at $\Gamma$. We then obtain an approximate solution to the eikonal equation by solving (4-7)-(4-9) with initial data $w=(x, \xi)$ and $w_{k}(0), k \geq 2$, which gives a change of coordinates $(t, w) \mapsto(t, w(t))$. If in any neighborhood of $\Gamma=\left\{\left(t, x_{0}, \xi_{0}\right): t \in I\right\}$ there exist points in $f^{-1}(0)$ where $\partial_{t} f<0$, then as before we can construct approximate solutions in any neighborhood of $\Gamma$ satisfying (5-11) with $k=1$. If $\partial_{t} f \geq 0$ on $f^{-1}(0)$ in some neighborhood of $\Gamma$, then by the invariance of condition ( $\Psi$ ) there will still exist a change of sign of $t \mapsto f(t, w(t))$ from + to - in any neighborhood of $\Gamma$ after the change of coordinates; see [Hörmander 1985b, Lemma 26.4.11]. (Recall that conditions (2-8) and (2-11) hold in some neighborhood of $\Gamma$.) Thus if $F^{\prime}(t)=-\operatorname{Im} w_{0}^{\prime}(t)=f(t, w(t))$ then $t \mapsto F(t)$ has a local maximum at some $t=t_{0}$, and after subtraction the maximum can be assumed to be equal to 0 . By choosing suitable initial value $\left(x_{0}, \xi_{0}\right)$ for (4-7) at $t=t_{0}$, we obtain

$$
\begin{equation*}
e^{i \lambda \omega(t, x)} \leq e^{\lambda\left(F(t)-c\left|x-x_{0}\right|^{2}\right)}, \quad\left|x-x_{0}\right| \ll 1 \tag{5-12}
\end{equation*}
$$

where $F^{\prime}(t)=f(t, w(t))$ so that $\max _{I} F(t)=0$ with $F(t)<0$ for some $t \notin I$ near $\partial I$.
Proposition 5.2. Assume that $t \mapsto f\left(t, x_{0}, \xi_{0}\right)$ changes sign from + to - as $t$ increases near $I$ and that $\partial_{t}^{k} \partial_{x}^{\alpha} \partial_{\xi}^{\beta} f\left(t, x_{0}, \xi_{0}\right)=0$ for all $t \in I$ when $|I| \neq 0$. Then we may solve (4-1) in a neighborhood $\Omega$ of $\Gamma=\left\{\left(t, x_{0}, \xi_{0}\right): t \in I\right\}$ modulo $\mathcal{O}\left(\left|x-x_{0}(t)\right|^{M}\right)$ for any $M$, with $\omega(t, x)$ given by (4-2) such that the curve $t \mapsto\left(x_{0}(t), \xi_{0}(t)\right), t \in\left(t_{1}, t_{2}\right)$, is arbitrarily close to $\Gamma, w_{k}(t) \in C^{\infty}, \operatorname{Im} w_{2}(t) \geq c>0$ when $t \in\left(t_{1}, t_{2}\right)$, $\min _{\left(t_{1}, t_{2}\right)} \operatorname{Im} w_{0}(t)=0$ and $\operatorname{Im} w_{0}\left(t_{j}\right)=c>0, j=1,2$.

Observe that since $\operatorname{Im} w_{0} \geq 0$ we find that $f\left(t_{0}, x_{0}\left(t_{0}\right), \xi_{0}\left(t_{0}\right)\right)=-\operatorname{Im} w_{0}^{\prime}(t)=0$ at a minimum $t_{0} \in$ $\left(t_{1}, t_{2}\right)$. As before, the errors that are $\mathcal{O}\left(\left|x-x_{0}\right|^{M}\right)$ in the eikonal equation will give terms that are bounded by $C_{M} \lambda^{-M / 2}$ for all $M$. Observe that cutting off where $\operatorname{Im} w_{0}>0$ will give errors that are $\mathcal{O}\left(\lambda^{-M}\right)$ for all $M$.

## 6. The transport equations

Next, we shall solve the transport equations given by the following terms in (3-15):

$$
\begin{align*}
D_{t} \phi+\partial_{\eta}^{2} B\left(t, x, y, \partial_{t, x} \omega, 0\right) D_{y}^{2} \phi / 2+A(t, & \left.x, y, \partial_{t, x} \omega, 0\right) D_{y} \phi \\
& +i \partial_{\xi} f\left(t, x, \partial_{x} \omega\right) D_{x} \phi+\sum_{j \geq 0} \lambda^{-j} R_{j}\left(t, x, y, D_{t, x, y}\right) \phi \tag{6-1}
\end{align*}
$$

near $\Gamma=\left\{\left(t, x_{0}, y_{0}, 0, \xi_{0}, 0\right): t \in I\right\}$. Here $R_{0}(t, x, y)=F_{0}\left(t, x, y, \partial_{t, x} \omega, 0\right)$ and when $m>0$ we have that $R_{m}\left(t, x, y, D_{t, x, y}\right)$ are differential operators of order $j$ in $t$, order $k$ in $x$ and order $\ell$ in $y$, where $j+k+\ell \leq m+2$. Assuming the conclusions in Proposition 5.2 hold, we shall choose suitable initial values of the amplitude $\phi$ at $t=t_{0}$, which is chosen so that $\operatorname{Im} w_{0}\left(t_{0}\right)=0$. Observe that the second-order differential operator given by the first four terms in (6-1) need not be solvable in general. Instead, by Lemma 6.1 we can treat the $D_{x}$ and $D_{y}$ terms as perturbations, using condition (3-6) in the infinite vanishing case.

Since the phase function $\omega(t, x)$ is complex-valued, we will replace the values of the symbols at $(\tau, \xi)=\partial_{t, x} \omega(t, x)$ by finite Taylor expansions at $\left(\operatorname{Re} w_{0}^{\prime}(t), \xi_{0}(t)\right)$. By (4-3) and (4-4) this will give expansions in powers of $x-x_{0}(t)$ and $\operatorname{Im} w_{0}^{\prime}(t)=-f\left(t, x_{0}(t), \xi_{0}(t)\right)$. Then, we shall solve the transport equations up to arbitrarily high powers of $x-x_{0}(t)$ and $f$. Since the imaginary part of the phase function $\operatorname{Im} \omega \geq 0$ vanishes of second order at $x=x_{0}(t)$, we will obtain by Lemma 6.1 below that this will give a solution modulo any negative power of $\lambda$.

We shall use the amplitude expansion

$$
\begin{equation*}
\phi(t, x, y)=\sum_{k \geq 0} \varrho^{-k} \phi_{k}(t, x, y) \tag{6-2}
\end{equation*}
$$

and solve the transport equation recursively in $k$. Here $\phi_{k}$ depends on $\varrho$ but with uniform bounds in a suitable symbol class, and $\varrho=\lambda^{1 / N}$ with $N$ to be determined later. By doing the change of variables $(t, x, y) \mapsto\left(t-t_{0}, x-x_{0}(t), y-y_{0}\right)$, we find that $D_{t}$ changes into $D_{t}-x_{0}^{\prime}(t) D_{x}$, which does not change the order of $R_{j}$ as differential operator. Thus we may assume $t_{0}=0, x_{0}(t) \equiv 0$ and $y_{0}=0$.

Next, we apply (6-1) on $\phi$ given by (6-2). Since $\varrho=\lambda^{1 / N}$, we obtain the terms

$$
\begin{equation*}
D_{t} \phi+A_{0}(t, x) D_{x} \phi+A_{1}(t, x, y) D_{y} \phi+A_{2}(t, x, y) D_{y}^{2} \phi+\sum_{j \geq 0} \varrho^{-j N} R_{j}\left(t, x, y, D_{t, x, y}\right) \phi \tag{6-3}
\end{equation*}
$$

where

$$
\begin{align*}
A_{0}(t, x) & =i \partial_{\xi} f\left(t, x, \xi_{0}(t)+\sigma(t, x)\right)-x_{0}^{\prime}(t) \\
A_{1}(t, x, y) & =A\left(t, x, y, \partial_{t} \omega(t, x), \xi_{0}(t)+\sigma(t, x), 0\right)  \tag{6-4}\\
A_{2}(t, x, y) & =\partial_{\eta}^{2} B_{2}\left(t, x, y, \partial_{t} \omega(t, x), \xi_{0}(t)+\sigma(t, x), 0\right) / 2 \tag{6-5}
\end{align*}
$$

Here $\sigma(t, x)$ is given by (4-4) and $\partial_{t} \omega(t, x)$ by (4-3), where the expansion will be up to a sufficiently high order in $x$. Observe that after the change of variables we have $\sigma(t, 0) \equiv 0$. The values of the symbols will as before be defined by finite Taylor expansions in the $\tau$ - and $\xi$-variables, which gives expansions in powers of $x$ and $f(t, 0, \xi(t))$.

We are going to construct solutions $\phi_{k}(t, x, y)=\phi_{k}(t, x, \varrho y)$ so that $y \mapsto \phi_{k}(t, x, y) \in C_{0}^{\infty}$ uniformly in $\varrho$, which gives localization in $|y| \lesssim \varrho^{-1}$. Therefore we shall choose $\varrho y$ as new $y$-coordinates. Then (6-3) becomes
$D_{t} \phi+A_{0}(t, x) D_{x} \phi+\varrho A_{1}(t, x, y / \varrho) D_{y} \phi+\varrho^{2} A_{2}(t, x, y / \varrho) D_{y}^{2} \phi$

$$
\begin{equation*}
+\sum_{j \geq 0} \varrho^{-j N} R_{j}\left(t, y / \varrho, x, D_{t}, D_{x}, \varrho D_{y}\right) \phi \tag{6-6}
\end{equation*}
$$

By Proposition 5.2 the phase function $e^{i \lambda w(t, x)}$ gives the cut-off in $x$, and we shall expand the symbols in powers of $x$. Now the Taylor expansion of $x \mapsto \varrho^{2} A_{2}(t, x, y / \varrho)$ will give terms that are $\mathcal{O}\left(\varrho^{2} x\right)$. Therefore we take $\varrho^{2} x$ as new $x$-coordinates, which gives

$$
\begin{align*}
D_{t} \phi+\varrho^{2} A_{0}\left(t, x / \varrho^{2}\right) D_{x} \phi+\varrho A_{1}\left(t, x / \varrho^{2}, y / \varrho\right) & D_{y} \phi+\varrho^{2} A_{2}\left(t, x / \varrho^{2}, y / \varrho\right) D_{y}^{2} \phi \\
& +\sum_{j \geq 0} \varrho^{-j N} R_{j}\left(t, y / \varrho, x / \varrho^{2}, D_{t}, \varrho^{2} D_{x}, \varrho D_{y}\right) \phi \tag{6-7}
\end{align*}
$$

Now the phase function $e^{i \lambda w(t, x)}$ is $\mathcal{O}\left(e^{-c \varrho^{N-4}|x|^{2}}\right)$ in the new coordinates. So if we take $N>4$, it suffices to solve the transport equation up to a sufficiently high order of $x$; then we may cut off where $|x| \lesssim 1$, which corresponds to $|x| \lesssim \varrho^{-2}$ in the original coordinates. Thus we expand in $x$ :

$$
\begin{align*}
\phi_{k}(t, x, y) & =\sum_{k, \alpha} \phi_{k, \alpha}(t, y) x^{\alpha} \quad \phi_{k, \alpha}(t, y) \in C_{0}^{\infty},  \tag{6-8}\\
A_{0}\left(t, x / \varrho^{2}\right) D_{x} & =\sum_{\alpha, j} A_{0, \alpha, j}(t) \varrho^{-2|\alpha|} x^{\alpha} D_{x_{j}} \\
A_{j}\left(t, x / \varrho^{2}, y / \varrho\right) & =\sum_{\alpha} A_{j, \alpha}(t, y / \varrho) \varrho^{-2|\alpha|} x^{\alpha}, \quad j>0 \\
R_{k}\left(t, x / \varrho^{2}, y / \varrho, \varrho D_{y}, \varrho^{2} D_{x}\right) & =\sum_{\alpha, \ell, v, \mu} R_{k, \alpha, \ell, v, \mu}(t, y / \varrho) \varrho^{-2|\alpha|+2|v|+|\mu|} x^{\alpha} D_{t}^{\ell} D_{x}^{v} D_{y}^{\mu} \tag{6-9}
\end{align*}
$$

Here $\ell+|\nu|+|\mu| \leq k+2$ so we have at most the factor $\varrho^{2|\nu|+|\mu|} \leq \varrho^{2 k+4}$ in (6-9). When $k=0$ we have $\ell+|\nu|+|\mu|=0$ and

$$
R_{0}\left(t, x / \varrho^{2}, y / \varrho\right)=\sum_{\alpha} R_{0, \alpha}(t, y / \varrho) \varrho^{-2|\alpha|} x^{\alpha}
$$

Observe that the coefficients in the expansions are given by expansions in powers of $f(t, 0, \xi(t))$. After cut-off in $x$ we find in the original coordinates that $\phi_{k}(t, x, y)=\varphi_{k}\left(t, \varrho^{2} x, \varrho y\right)$, where $\varphi_{k}$ for any $t$ is uniformly bounded in $C_{0}^{\infty}$.

We shall first apply (6-7) on $\phi_{0}$ and expand in $x$. Then we find that the terms that are independent of $x$ are

$$
\begin{equation*}
D_{t} \phi_{0,0}-i \varrho^{2} \sum_{j} A_{0,0, j}(t) \phi_{0, e_{j}}+\varrho A_{1,0}(t, y / \varrho) D_{y} \phi_{0,0}+\varrho^{2} A_{2,0}(t, y / \varrho) D_{y}^{2} \phi_{0,0}+R_{0,0}(t, y / \varrho) \phi_{0,0} \tag{6-10}
\end{equation*}
$$

We shall need the following result, which gives estimates on $f$ and $A_{j}$ on the interval of integration. It will be proved in the next section. In the following, we shall denote $f(t)=f\left(t, 0, \xi_{0}(t)\right)$ and $F(t)=\int_{0}^{t} f(s) d s$. Observe that $f(0)=0$ since $\operatorname{Im} w_{0}^{\prime}(0)=0$.

Lemma 6.1. Assume that the conclusions in Proposition 5.2 hold and that (3-6) holds if $t \mapsto f(t)$ vanishes of infinite order at 0 . Then there exists $\varepsilon$ and $C \geq 1$ with the property that if $N \geq C, \varrho=\lambda^{1 / N} \geq C$ and

$$
\begin{equation*}
|f(t)|+\left|\int_{0}^{t}\right| A_{0}(s, 0)\left|+\left|A_{1}(s, 0, y / \varrho)\right|+\left\|A_{2}(s, 0, y / \varrho)\right\| d s\right| \geq C / \varrho^{3} \tag{6-11}
\end{equation*}
$$

holds for some $|y| \leq \varrho / C$, then $\lambda F(s) \leq-\lambda^{\varepsilon} / C$ for some $s$ in the interval connecting 0 and $t$.

Observe that if Lemma 6.1 holds for some $\varepsilon$ and $C$, then it trivially holds for smaller $\varepsilon$ and larger $C$. We shall assume that $\varepsilon<1$ and that both $N$ and $\lambda$ are large enough so that the conclusion in Lemma 6.1 holds. Since (6-11) does not hold when $t=0$, we can choose the maximal interval $I$ containing 0 such that (6-11) does not hold in $I$; thus

$$
\begin{equation*}
|f(t)|+\left|\int_{0}^{t}\right| A_{0}(s, 0)\left|+\left|A_{1}(s, 0, y / \varrho)\right|+\left\|A_{2}(s, 0, y / \varrho)\right\| d s\right|<C / \varrho^{3}, \quad t \in I \tag{6-12}
\end{equation*}
$$

when $|y| \leq \varrho / C$. By definition we obtain that (6-11) holds for some $|y| \leq \varrho / C$ when $t \in \partial I$, so Lemma 6.1 gives that $\lambda F \lesssim-\lambda^{\varepsilon}$ at $\partial I_{0}$ for some open interval $I_{0} \subseteq I$ that contains 0 . This means that $e^{i \lambda \omega(t, 0)}=$ $e^{\lambda F(t)} \leq C_{N} \lambda^{-N}$ for any $N$ at $\partial I_{0}$ when $\lambda \gg 1$. Since $F^{\prime}=f$ is uniformly bounded and the left-hand side of (6-12) is Lipschitz continuous, we may cut off near $I_{0}$ with $\chi(t) \in S\left(1, \lambda^{6 / N} d t^{2}\right) \subset S\left(1, \lambda^{2-2 \varepsilon} d t^{2}\right)$ for $N \gg 1$ so that $\chi(0) \neq 0, \lambda F(t) \lesssim-\lambda^{\varepsilon}$ in supp $\chi^{\prime}$ and (6-12) holds with some $C$ when $t \in \operatorname{supp} \chi$ and $|y| \leq \varrho / C$. Then as before, the cut-off errors can be absorbed by the exponential and the expansion in powers of $f(t, 0, \xi(0))=f(t)$ is justified. In fact, $f(t)=\mathcal{O}\left(\varrho^{-3}\right)$ in supp $\chi$, which gives errors of any negative power of $\varrho=\lambda^{1 / N}$. The bound on the integral in (6-12) means that we can ignore the $A_{j}$ terms in (6-10) in supp $\chi$ modulo lower-order terms in $\varrho$. In the following we shall change the notation and let $I=\operatorname{supp} \chi$. We need to measure the error terms in the following way.
Definition 6.2. For $a(t) \in L^{\infty}(\mathbb{R})$ and $\kappa>0$, we say that $a(t) \in I(\kappa)$ if $\int_{0}^{t} a(s) d s=\mathcal{O}(\kappa)$ for all $t \in I$.
For example, $f(t) \in I\left(\varrho^{-3}\right)$, and since the integral in (6-12) is $\mathcal{O}\left(\varrho^{-3}\right)$ in $I$, the integrand is in $I\left(\varrho^{-3}\right)$. Then according to (6-12) it suffices to solve

$$
\begin{equation*}
D_{t} \phi_{0,0}=-R_{0,0} \phi_{0,0}, \quad t \in I \tag{6-13}
\end{equation*}
$$

to obtain that the terms in (6-10) are in $I\left(\varrho^{-1}\right)$; here $R_{0,0}(t, y / \varrho) \in C^{\infty}$ uniformly since $\varrho \geq 1$. Now we can solve (6-13) with $\phi_{0,0}(0, y)=\phi(y) \in C_{0}^{\infty}$ uniformly with support where $|y| \ll 1$ such that $\phi(0)=1$. In fact, the solution is $\phi_{0,0}(t, y)=E(t, y) \phi(y)$, where

$$
E(t, y)=\exp \left(-i \int_{0}^{t} R_{0,0}(s, y / \varrho) d s\right), \quad t \in I
$$

is uniformly bounded in $C^{\infty}$. Thus $\phi_{0,0}(t, y) \in C^{\infty}$ uniformly and by choosing $\phi(y)$ with sufficiently small support, we obtain for any $t \in I$ that $\phi_{0,0}(t, \cdot)$ has support in a sufficiently small compact set in which (6-12) holds.

The coefficients of the terms in (6-7) which are homogeneous of degree $\alpha \neq 0$ in $x$ are

$$
\begin{equation*}
D_{t} \phi_{0, \alpha}+R_{0,0}(t, y / \varrho) \phi_{0, \alpha}-i \sum_{\substack{|\beta|=1 \\ j}} A_{0, \beta, j}(t)\left(\alpha_{j}+1-\beta_{j}\right) \phi_{0, \alpha+e_{j}-\beta}+\sum_{|\beta|=1} A_{2, \beta}(t, y / \varrho) D_{y}^{2} \phi_{0, \alpha-\beta} \tag{6-14}
\end{equation*}
$$

modulo $I\left(\varrho^{-1}\right)$. Letting $\Phi_{k, j}=\left\{\phi_{k, \alpha}\right\}_{|\alpha|=j}$ and $\Phi_{k}=\left\{\Phi_{k, j}\right\}_{j}$ for $k, j \geq 0$, we find that (6-14) vanishes if $\Phi_{0}$ satisfies the system

$$
\begin{equation*}
D_{t} \Phi_{0, k}=S_{0,0}^{k} \Phi_{0, k}+S_{0,1}^{k} \Phi_{0, k-1} \tag{6-15}
\end{equation*}
$$

where $S_{0,0}^{k}(t)$ is a uniformly bounded matrix depending on $t$, and $S_{0,1}^{k}\left(t, y / \varrho, D_{y}\right)$ is a system of uniformly bounded differential operators of order 2 in $y$ when $|y| \lesssim \varrho$. Let $E_{0, k}(t)$ be the fundamental solution to $D_{t} E_{0, k}=S_{0,0}^{k} E_{0, k}$ so that $E_{0, k}(0)=\mathrm{Id}$. Then letting

$$
\Phi_{0, k}(t, y)=E_{0, k}(t) \Psi_{0, k}(t, y)
$$

the system (6-15) reduces to

$$
D_{t} \Psi_{0, k}(t, y)=E_{0, k}^{-1} S_{0,1}^{k} E_{0, k} \Psi_{0, k-1}(t, y)
$$

This is a recursion equation which we can solve uniformly in $I$ with $\Psi_{0, k}(t, y)$ having initial values $\Psi_{0, k}(0, y) \equiv 0$ for $0<k \leq M$. Observe that since the initial data $\Phi_{0, k}(0, y)$ has compact support, we find that $\Phi_{0, k}(t, y) \in C^{\infty}$ uniformly. For any $t$ we find that $\Phi_{0, k}(t, y)$ has support in a sufficiently small compact set so that (6-12) holds for any $t \in I$.

We shall now apply (6-7) to $\phi$ given by the full expansion (6-8). We find that the coefficients of the terms in (6-7) which are homogeneous of degree $\alpha \neq 0$ in $x$ are equal to

$$
\begin{align*}
& \varrho^{-1}\left(D_{t} \phi_{1, \alpha}+R_{0,0}(t, y / \varrho) \phi_{1, \alpha}-i \sum_{\substack{|\beta|=1 \\
j}} A_{0, \beta, j}(t)\left(\alpha_{j}+1-\beta_{j}\right) \phi_{1, \alpha+e_{j}-\beta}+\sum_{|\beta|=1} A_{2, \beta}(t, y / \varrho) D_{y}^{2} \phi_{1, \alpha-\beta}\right. \\
& \left.\quad+\sum_{|\beta|=1} A_{1, \beta}(t, y / \varrho) D_{y}^{2} \phi_{0, \alpha-\beta}-i \varrho^{3} \sum_{j} A_{0,0, j}(t)\left(\alpha_{j}+1\right) \phi_{0, \alpha+e_{j}}+\varrho^{3} A_{2,0}(t, y / \varrho) D_{y}^{2} \phi_{0, \alpha}\right) \tag{6-16}
\end{align*}
$$

modulo $I\left(\varrho^{-2}\right)$. We find that (6-16) vanishes if $\Phi_{1}$ satisfies the system

$$
\begin{equation*}
D_{t} \Phi_{1, k}=S_{1,0}^{k} \Phi_{1, k}+S_{1,1}^{k} \Phi_{1, k-1}+\mathbb{A}_{1}^{0} \Phi_{0} \tag{6-17}
\end{equation*}
$$

where $S_{1,0}^{k}(t)$ is a uniformly bounded matrix depending on $t, S_{1,1}^{k}\left(t, y / \varrho, D_{y}\right)$ is a system of uniformly bounded differential operators of order 2 when $|y| \lesssim \varrho$ and $\mathbb{A}_{1}^{0}$ is a differential operator in $y$ of order 2 with coefficients in $I(1)$ because of (6-12). By letting $\Phi_{1, k}=E_{1, k} \Psi_{1, k}$ with the fundamental solution $E_{1, k}$ to $D_{t} E_{1, k}=S_{1,0}^{k} E_{1, k}, E_{1, k}(0)=\mathrm{Id}$, this reduces to the equation

$$
D_{t} \Psi_{1, k}=E_{1, k}^{-1} S_{1,1}^{k} E_{1, k-1} \Psi_{1, k-1}+E_{1, k}^{-1} \mathbb{A}_{1}^{0} \Phi_{0}
$$

Thus we can solve (6-17) in $I$ recursively with uniformly bounded $\Phi_{1, k}$ having initial values $\Phi_{1, k}(0, y) \equiv 0$, $k \geq 0$. But observe that $\Phi_{1}$ is not in $C^{\infty}$ uniformly; instead we have $D_{t}^{j} \Phi_{1}=\mathcal{O}\left(\varrho^{3}\right)$ if $j \geq 1$, since $\left|\partial_{t}^{j} \mathbb{A}_{1}^{0}\right| \leq C_{j} \varrho^{3}$ for all $j$ by (6-16). For that reason, we shall define $S_{\varrho}^{3} \subset C^{\infty}$ by

$$
\begin{equation*}
\left|\partial_{t}^{j} \partial_{y}^{\alpha} \phi(t, y)\right| \leq C_{j, \alpha} \varrho^{3 j} \quad \forall j, \alpha \tag{6-18}
\end{equation*}
$$

when $\phi \in S_{\varrho}^{3}$. Observe that $\phi \in S_{\varrho}^{3}$ if and only if $\phi(t, y)=\chi\left(\varrho^{3} t, y\right)$, where $\chi \in C^{\infty}$ uniformly, and that the operator $\varrho^{-3} D_{t}$ maps $S_{\varrho}^{3} \mapsto S_{\varrho}^{3}$. Note that the expansion of the symbols also contains terms with factors $\varrho^{3} f^{k}, k \geq 1$, which are uniformly bounded in $S_{\varrho}^{3}$ for $t \in I$ by (6-12). Since $\int_{0}^{t} \mathbb{A}_{1}^{0} d t \in S_{\varrho}^{3}$ in $I$, we find that $\Phi_{1} \in S_{\varrho}^{3}$ in $I$.

Recursively, the coefficients of the terms in (6-7) that are homogeneous in $x$ of degree $\alpha$ are

$$
\begin{align*}
& \varrho^{-k}\left(D_{t} \phi_{k, \alpha}-i \sum_{\beta \neq 0} A_{0, \beta, j}(t)\left(\alpha_{j}+1-\beta_{j}\right) \phi_{k+2-2|\beta|, \alpha+e_{j}-\beta}\right. \\
& \quad+\sum_{\beta \neq 0} A_{1, \beta}(t, y / \varrho) D_{y} \phi_{k+1-2|\beta|, \alpha-\beta}+\sum_{\beta \neq 0} A_{2, \beta}(t, y / \varrho) D_{y}^{2} \phi_{k+2-2|\beta|, \alpha-\beta} \\
& \quad-i \varrho^{3} \sum_{j} A_{0,0, j}(t)\left(\alpha_{j}+1\right) \phi_{k-1, \alpha+e_{j}}+\varrho^{3} A_{1,0}(t, y / \varrho) D_{y} \phi_{k-2, \alpha}+\varrho^{3} A_{2,0}(t, y / \varrho) D_{y}^{2} \phi_{k-1, \alpha} \\
& \left.\quad+\sum_{\ell+|\nu|+|\mu| \leq j+2} \varrho^{-j N} R_{j, \beta, \ell, v, \mu}(t, y / \varrho) c_{\alpha, \beta, \nu} \varrho^{-2|\beta|+2|\nu|+|\mu|+i+3 \ell}\left(\varrho^{-3} D_{t}\right)^{\ell} D_{y}^{\mu} \phi_{k-i, \alpha+\nu-\beta}\right) \tag{6-19}
\end{align*}
$$

modulo $I\left(\varrho^{-k-1}\right)$. Here the last sum has $\ell+|\nu|+|\mu|=0$ when $j=0,\left(\varrho^{-3} D_{t}\right)^{\ell} D_{y}^{\mu} \operatorname{maps} S_{\varrho}^{3} \mapsto S_{\varrho}^{3}$ and the values of the symbols are given by a finite expansion in powers of $f(t)$.

Since $\phi_{j} \in S_{\varrho}^{3}$ we obtain that the terms in (6-19) are in $I\left(\varrho^{-k-1}\right)$ if

$$
\begin{align*}
& D_{t} \phi_{k, \alpha}-i \sum_{\beta \neq 0} A_{0, \beta, j}(t)\left(\alpha_{j}+1-\beta_{j}\right) \phi_{k+2-2|\beta|, \alpha+e_{j}-\beta}+\sum_{\beta \neq 0} A_{1, \beta}(t, y / \varrho) D_{y} \phi_{k+1-2|\beta|, \alpha-\beta} \\
& \quad+\sum_{\beta \neq 0} A_{2, \beta}(t, y / \varrho) D_{y}^{2} \phi_{k+2-2|\beta|, \alpha-\beta}-i \varrho^{3} \sum_{j} A_{0,0, j}(t)\left(\alpha_{j}+1\right) \phi_{k-1, \alpha+e_{j}} \\
& \quad+\varrho^{3} A_{1,0}(t, y / \varrho) D_{y} \phi_{k-2, \alpha}+\varrho^{3} A_{2,0}(t, y / \varrho) D_{y}^{2} \phi_{k-1, \alpha} \\
& =-\sum_{\substack{i+3 \ell+2|\nu|+|\mu|=j N+2|\beta| \\
\ell+|\nu|+|\mu| \leq j+2}} R_{j, \beta, \ell, v, \mu}(t, y / \varrho) c_{\alpha, \beta, v}\left(\varrho^{-3} D_{t}\right)^{\ell} D_{y}^{\mu} \phi_{k-i, \alpha+v-\beta} . \tag{6-20}
\end{align*}
$$

When $j=0$ we find that $\ell+|\nu|+|\mu|=0, i=2|\beta|$ and we only have an expansion in $\beta$ in the last sum. Now if $j>0, \ell+|\nu|+|\mu| \leq j+2$ and $i+3 \ell+2|\nu|+|\mu|=j N+2|\beta|$ then we find that

$$
j N \leq i+3 \ell+2|v|+|\mu|<i+3(j+2)
$$

which gives $i \geq j(N-3)-6 \geq N-9 \geq 1$ if $N \geq 10$. Thus we find that (6-20) can be written as

$$
\begin{equation*}
D_{t} \Phi_{k}=\mathbb{A}_{0}^{k} \Phi_{k}+\mathbb{A}_{1}^{k} \Phi_{k-1}+\mathbb{A}_{2}^{k} \Phi_{k-2}+\cdots \tag{6-21}
\end{equation*}
$$

where $\int_{0}^{t} \mathbb{A}_{j}^{k} d t$ is a uniformly bounded differential operator on $S_{\varrho}^{3}$ for $t \in I$ and $j>0$. We have

$$
\left\{\mathbb{A}_{0}^{k} \Phi_{k}\right\}_{j}=S_{k, 0}^{j} \Phi_{k, j}+S_{k, 1}^{j} \Phi_{k, j-1}
$$

where $S_{k, 0}^{j}(t)$ is a uniformly bounded matrix depending on $t$, and $S_{k, 1}^{j}\left(t, y / \varrho, D_{y}\right)$ is a system of uniformly bounded differential operators of order 2 when $|y| \lesssim \varrho$. By letting $\Phi_{k, j}=E_{k, j} \Psi_{k, j}$ with the fundamental solution $E_{k, j}$ to $D_{t} E_{k, j}=S_{k, 0}^{j} E_{k, j}, E_{k, j}(0)=\mathrm{Id}$, (6-21) becomes a system of recursion equations in $j$ and $k$. Thus (6-21) can be solved in $I$ with $\Phi_{k} \in S_{\varrho}^{3}$ having initial values $\Phi_{k}(0) \equiv 0, k>0$. We find from (6-8) and the definition of $S_{\varrho}^{3}$ that $\phi_{k}(t, x, y)=\phi_{k}\left(\varrho^{3} t, \varrho^{2} x, \varrho y\right)$, where $\phi_{k} \in C^{\infty}$ uniformly when $t \in I$. Thus we can solve the transport equation (6-1) up to any negative power of $\lambda$. Observe that by
cutting off in $t$ and $x$, we may assume that $\phi_{k} \in C_{0}^{\infty}$ has fixed compact support in ( $x, y$ ) and support where $|t| \lesssim \varrho^{3}$. It follows that the support of $\phi_{k}$ can be chosen in an arbitrarily small neighborhood of $\Gamma$ for large enough $\lambda$. Changing to the original coordinates, we obtain the following result.

Proposition 6.3. Assume that the conclusions in Proposition 5.2 hold, and that (3-6) is satisfied near $\Gamma$ when the sign change of $t \mapsto f\left(t, x_{0}, \xi_{0}\right)$ is of infinite order. If $\varrho=\lambda^{1 / N}$ for sufficiently large $N$, then for any $K$ and $M$ we can solve the transport equations (6-20) for $k \leq K$ and $|\alpha| \leq M$ near $\left\{\left(t, x_{0}(t), y_{0}\right): t \in\left[t_{1}, t_{2}\right]\right\}$. By (6-8) this gives

$$
\phi_{k}(t, x, y)=\phi_{k}\left(\varrho^{3}\left(t-t_{0}\right), \varrho^{2}\left(x-x_{0}(t)\right), \varrho\left(y-y_{0}\right)\right), \quad k \leq K
$$

where $\phi_{k}(t, x, y) \in C^{\infty}$ uniformly, has support where $|x|+|y| \lesssim 1$ and $|t| \lesssim \varrho^{3}$, and $\phi_{0}(0,0,0)=1$ for some $t_{0} \in\left(t_{1}, t_{2}\right)$ such that $\operatorname{Im} w_{0}\left(t_{0}\right)=0$.

## 7. The rate of change of sign

We have showed that $t \mapsto f(t, x, \xi)$ changes sign from + to - on an interval $I$. Then

$$
\begin{equation*}
F(t)=\int^{t} f\left(s, x_{0}(s), \xi_{0}(s)\right) d s=\int^{t} f(s) d s \tag{7-1}
\end{equation*}
$$

has a local maximum in the interval. By choosing that maximum as the starting point, we may assume it is equal to 0 so that $F(t) \leq 0$. By changing $t$-coordinate, we may assume $F(0)=0$. We shall study how the size of the derivative $f$ affects the size of the function $F$.

Lemma 7.1. Assume that $0 \geq F(t) \in C^{\infty}$ has local maximum at $t=0$, and let $I_{t_{0}}$ be the closed interval joining 0 and $t_{0} \in \mathbb{R}$. If

$$
\max _{I_{t_{0}}}\left|F^{\prime}(t)\right|=\left|F^{\prime}\left(t_{0}\right)\right|=\kappa \leq 1
$$

with $\left|t_{0}\right| \geq \kappa^{\varrho}$ for some $\varrho>0$, then we have $\min _{I_{t_{0}}} F(t) \leq-C_{\varrho} \kappa^{1+\varrho}$. The constant $C_{\varrho}>0$ only depends on $\varrho$ and the bounds on $F$ in $C^{\infty}$.
Proof. Let $f=F^{\prime}$. Then since $F(t)=F(0)+\int_{0}^{t} f(s) d s \leq \int_{0}^{t} f(s) d s$, it is no restriction to assume the maximum $F(0)=0$. By switching $t$ to $-t$, we may assume $t_{0} \leq-\kappa^{\varrho}<0$. Let

$$
\begin{equation*}
g(t)=\kappa^{-1} f\left(t_{0}+t \kappa^{\varrho}\right) \tag{7-2}
\end{equation*}
$$

Then $|g(0)|=1,|g(t)| \leq 1$ for $0 \leq t \leq 1$ and

$$
\left|g^{(N)}(t)\right|=\kappa^{\varrho N-1}\left|f^{(N)}\left(t_{0}+t \kappa^{\varrho}\right)\right| \leq C_{N}
$$

when $N \geq 1 / \varrho$ for $0 \leq t \leq 1$. Using the Taylor expansion at $t=0$ for $N \geq 1 / \varrho$, we find

$$
\begin{equation*}
g(t)=p(t)+r(t) \tag{7-3}
\end{equation*}
$$

where $p$ is the Taylor polynomial of order $N-1$ of $g$ at 0 , and

$$
\begin{equation*}
r(t)=t^{N} \int_{0}^{1} g^{(N)}(t s)(1-s)^{N-1} d s /(N-1)! \tag{7-4}
\end{equation*}
$$

is uniformly bounded in $C^{\infty}$ for $0 \leq t \leq 1$ and $r(0)=0$. Since $g$ also is bounded on the interval, we find that $p(t)$ is uniformly bounded in $0 \leq t \leq 1$. Since all norms on the finite-dimensional space of polynomials of fixed degree are equivalent, we find that $p^{(k)}(0)=g^{(k)}(0)$ are uniformly bounded for $0 \leq k<N$ which implies that $g(t)$ is uniformly bounded in $C^{\infty}$ for $0 \leq t \leq 1$. Since $|g(0)|=1$, there exists a uniformly bounded $\delta^{-1} \geq 1$ such that $|g(t)| \geq \frac{1}{2}$ when $0 \leq t \leq \delta$; thus $g$ has the same sign in that interval. Since $g(s)=\kappa^{-1} f\left(t_{0}+s \kappa^{\varrho}\right)$, we find

$$
\begin{equation*}
\delta / 2 \leq\left|\int_{0}^{\delta} g(s) d s\right|=\left|\kappa^{-\varrho} \int_{t_{0}}^{t_{0}+\delta \kappa^{\varrho}} \kappa^{-1} f(t) d t\right| . \tag{7-5}
\end{equation*}
$$

Since $t_{0}+\delta \kappa^{\varrho} \leq 0$, we find that the variation of $F(t)$ on $\left[t_{0}, 0\right]$ is greater than $\delta \kappa^{1+\varrho} / 2$ and since $F \leq 0$, we find that the minimum of $F$ on $I_{t_{0}}$ is smaller than $-\delta \kappa^{1+\varrho} / 2$.

Proof of Lemma 6.1. As before we let $F(t)$ satisfy $F(0)=0$ and $F^{\prime}(t)=f(t)$, where $f(t)=f\left(t, 0, \xi_{0}(t)\right)$ satisfies $f(0)=0$. We have assumed that the estimate (3-6) holds near $\Gamma$ if $f(t)$ vanishes of infinite order at $t=0$. Observe that the term $x_{0}^{\prime}(t)$ in $A_{0}$ can be estimated by $\left|\partial_{w} f\left(t, 0, \xi_{0}(t)\right)\right|$ by (4-7), which gives that $\left|A_{0}(t, 0)\right| \lesssim\left|\partial_{w} f\left(t, 0, \xi_{0}(t)\right)\right|$. We find from (4-3), (4-6) and (4-7) that

$$
\left|\partial_{t} \omega(t, 0)\right| \lesssim|f(t)|+\left|\partial_{w} f\left(t, 0, \xi_{0}(t)\right)\right|
$$

Thus (6-11) follows if

$$
\begin{equation*}
|f(t)|+\left|\int_{0}^{t}\right| f(s)\left|+A_{0}(s, 0)+A_{1}(s, 0, y / \varrho)+A_{2}(s, 0, y / \varrho) d s\right| \gtrsim \varrho^{-3} \tag{7-6}
\end{equation*}
$$

where

$$
\begin{align*}
A_{0}(t) & =\left|\partial_{w} f\left(t, 0, \xi_{0}(t)\right)\right| \\
A_{1}(t, y / \varrho) & =\left|A\left(t, 0, y / \varrho, 0, \xi_{0}(t), 0\right)\right|  \tag{7-7}\\
A_{2}(t, y / \varrho) & =\left\|\partial_{\eta}^{2} B\left(t, 0, y / \varrho, 0, \xi_{0}(t), 0\right)\right\| . \tag{7-8}
\end{align*}
$$

In the following we shall suppress the $y$-variables in (7-6); the results will be uniform when $|y| \leq c \varrho$ for some $c>0$ since (3-6) holds near $\Gamma$. Observe that if $|f(s)|$ and $A_{j}(s)$ are $\ll \varrho^{-3}$ for $0 \leq j \leq 2$ when $s$ is between 0 and $t$, then (7-6) does not hold.

We shall first consider the case when $|f(t)| \cong|t|^{m}$ vanishes of finite order at $t=0$. Then the order must be odd so we find $F(t)=\int_{0}^{t} f(s) d s \leq 0$ and $c \leq|F(t)| / t^{2 k} \leq C<0$ for some $k>0$. Thus we find

$$
\begin{equation*}
\varrho^{-3} \lesssim\left|\int_{0}^{t}\right| f(s)\left|+A_{0}(s)+A_{1}(s)+A_{2}(s) d s\right| \lesssim|t| \lesssim|F(t)|^{1 / 2 k} \tag{7-9}
\end{equation*}
$$

implies that $|F(t)| \gtrsim \varrho^{-6 k}$. Since $\lambda=\varrho^{N}$, we then obtain $\lambda F(t) \lesssim-\varrho^{N-6 k} \leq-\varrho=\lambda^{1 / N}$ if $N>6 k$. The case when $|t|^{2 k-1} \cong|f(t)| \gtrsim \varrho^{-3}$ gives that $|t| \gtrsim \varrho^{-3 /(2 k-1)}$ so $\lambda F(t) \lesssim-\varrho^{N-6 k /(2 k-1)} \leq-\varrho$ if $N>6$. Now one of these cases must hold if (6-11) holds, so we get the result in the finite vanishing case.

Next, we consider the infinite vanishing case. Then we have assumed that condition (3-6) holds, which means that

$$
\sum_{j=0}^{2} A_{j}(t) \lesssim|f(t)|^{\varepsilon}
$$

which implies that $A_{j}(0)=0$ for all $j$. Now we assume that (7-6) holds at $t$; by switching $t$ and $-t$, we may assume $t>0$. Then we obtain for some $s \in[0, t]$ that $|f(s)| \geq c \varrho^{-3}$ or $A_{j}(s) \geq c \varrho^{-3}$ for some $c>0$ and $j$. Now we define $t_{0}$ as the smallest $t_{0}>0$ such that $\left|f\left(t_{0}\right)\right|=c \varrho^{-3}$ or $A_{j}\left(t_{0}\right)=c \varrho^{-3}$ for some $j$; then $t_{0} \leq t$. Then we obtain from condition (3-6) in the first case that $c \varrho^{-3}=\left|f\left(t_{0}\right)\right| \lesssim\left|f\left(t_{0}\right)\right|^{\varepsilon}$ and in the second case that

$$
\begin{equation*}
c \varrho^{-3}=A_{j}\left(t_{0}\right) \lesssim\left|f\left(t_{0}\right)\right|^{\varepsilon} \tag{7-10}
\end{equation*}
$$

Since $\varrho=\lambda^{1 / N}$, we find in both cases that

$$
\begin{equation*}
\lambda^{-3 / \varepsilon N}=\kappa \leq c\left|f\left(t_{0}\right)\right|, \quad c>0 \tag{7-11}
\end{equation*}
$$

where $\lambda \gg 1$ if and only if $\kappa \ll 1$. By taking the smallest $t_{0}>0$ such that (7-11) is satisfied, we find that $|f(t)| \leq\left|f\left(t_{0}\right)\right|$ for $0 \leq t \leq t_{0}$. Since $f(t)$ vanishes of infinite order at $t=0$, we find using Taylor's formula that $|f(t)| \leq C_{M}|t|^{M}$ for any positive integer $M$. (Actually, it suffices to take $M=1$.) Condition (7-11) then gives

$$
\begin{equation*}
\kappa^{1 / M} \lesssim\left|f\left(t_{0}\right)\right|^{1 / M} \lesssim\left|t_{0}\right| \tag{7-12}
\end{equation*}
$$

so using Lemma 7.1 with $\varrho=1 / M$, we find that

$$
\begin{equation*}
\min _{0 \leq s \leq t_{0}} F(s) \lesssim-\kappa^{1+1 / M}=-\lambda^{-3(1+1 / M) / \varepsilon N}, \quad \lambda \gg 1 \tag{7-13}
\end{equation*}
$$

Thus we find that $\min _{0 \leq s \leq t_{0}} F(s) \lesssim-\lambda^{c-1}$ for some $c>0$ if $3(1+1 / M) / \varepsilon N<1$, that is, $N>3(1+1 / M) / \varepsilon$, which gives Lemma 6.1.

## 8. The proof of Theorem 2.7

We shall use the following modification of [Hörmander 1985b, Lemma 26.4.15]. Recall that $\|u\|_{(k)}$ is the $L^{2}$ Sobolev norm of order $k$ of $u \in C_{0}^{\infty}$ and let $\mathcal{D}_{\Gamma}^{\prime}=\left\{u \in \mathcal{D}^{\prime}: \mathrm{WF}(u) \subset \Gamma\right\}$ for $\Gamma \subseteq T^{*} \mathbb{R}^{n}$.

Lemma 8.1. Let

$$
\begin{equation*}
u_{\lambda}(x)=\lambda^{(n-1) \delta / 2} \exp (i \lambda \omega(x)) \sum_{j=0}^{M} \phi_{j}\left(\lambda^{\delta} x\right) \lambda^{-j \varrho}, \quad \lambda \geq 1 \tag{8-1}
\end{equation*}
$$

with $\varrho>0,0<\delta<1, \omega \in C^{\infty}\left(\mathbb{R}^{n}\right)$ satisfying $\operatorname{Im} \omega \geq 0,|d \operatorname{Re} \omega| \geq c>0$, and $\phi_{j} \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$. Here $\omega$ and $\phi_{j}$ may depend on $\lambda$ but uniformly, and $\phi_{j}$ has fixed compact support in all but one of the variables, for which the support is bounded by $C \lambda^{\delta}$. Then for any integer $N$ we have

$$
\begin{equation*}
\left\|u_{\lambda}\right\|_{(-N)} \leq C \lambda^{-N}, \quad \lambda \geq 1 \tag{8-2}
\end{equation*}
$$

If $\phi_{0}\left(x_{0}\right) \neq 0$ and $\operatorname{Im} \omega\left(x_{0}\right)=0$ for some $x_{0}$ then there exists $c>0$ so that

$$
\begin{equation*}
\left\|u_{\lambda}\right\|_{(-N)} \geq c \lambda^{-N-n / 2+(n-1) \delta / 2}, \quad \lambda \geq 1, \quad \forall N \tag{8-3}
\end{equation*}
$$

Let $\Sigma=\bigcap_{\lambda \geq 1} \bigcup_{j} \operatorname{supp} \phi_{j}\left(\lambda^{\delta} \cdot\right)$ and let $\Gamma$ be the cone generated by

$$
\begin{equation*}
\{(x, \partial \omega(x)): x \in \Sigma, \operatorname{Im} \omega(x)=0\} \tag{8-4}
\end{equation*}
$$

Then for any $k$ we find $\lambda^{k} u_{\lambda} \rightarrow 0$ in $\mathcal{D}_{\Gamma}^{\prime}$, so $\lambda^{k} A u_{\lambda} \rightarrow 0$ in $C^{\infty}$ if $A$ is a pseudodifferential operator such that $\operatorname{WF}(A) \cap \Gamma=\varnothing$. The estimates are uniform if $\omega \in C^{\infty}$ uniformly with fixed lower bound on $|d \operatorname{Re} \omega|$, and $\phi_{j} \in C_{0}^{\infty}$ uniformly with the support condition.

In the expansion (8-1), we shall take $\varrho=1 / N$ and $\delta=3 / N$ with $N>3$, and the cone $\Gamma$ will be generated by

$$
\begin{equation*}
\left\{\left(t, x_{0}(t), y_{0}, 0, \xi_{0}(t), 0\right): t \in I\right\} \tag{8-5}
\end{equation*}
$$

where $I=\left\{t: \operatorname{Im} w_{0}(t)=0\right\}$. Observe that the phase function in (4-2) will satisfy the conditions in Lemma 8.1 near $\left\{\left(t, x_{0}(t), y_{0}\right): t \in I\right\}$ since $\xi_{0}(t) \neq 0$ and $\operatorname{Im} \omega(t, x) \geq 0$ by Proposition 5.2. Also, we find from Proposition 6.3 that the functions $\phi_{k}$ will satisfy the conditions in Lemma 8.1 with $\delta=3 / N$ after making the change of variables $(t, x, y) \mapsto\left(t-t_{0}, x-x_{0}(t), y-y_{0}\right)$ since $\phi_{0}\left(t_{0}, x_{0}\left(t_{0}\right), y_{0}\right)=1$. Observe that the conclusions of Lemma 8.1 are invariant under uniform changes of coordinates.

Proof of Lemma 8.1. We shall modify the proof of [Hörmander 1985b, Lemma 26.4.15] to this case. We have

$$
\begin{equation*}
\hat{u}_{\lambda}(\xi)=\lambda^{(n-1) \delta / 2} \sum_{j=0}^{M} \lambda^{-j \delta} \int e^{i \lambda \omega(x)-i\langle x, \xi\rangle} \phi_{j}\left(\lambda^{\delta} x\right) d x . \tag{8-6}
\end{equation*}
$$

Let $U$ be a neighborhood of the projection on the second component of the set in (8-4). When $\xi / \lambda \notin U$, for $\lambda \gg 1$ we find that

$$
\bigcup_{j} \operatorname{supp} \phi_{j}\left(\lambda^{\delta} \cdot\right) \ni x \mapsto(\lambda \omega(x)-\langle x, \xi\rangle) /(\lambda+|\xi|)
$$

is in a compact set of functions with nonnegative imaginary part with a fixed lower bound on the gradient of the real part. Thus, by integrating by parts we find for any positive integer $k$ that

$$
\begin{equation*}
\left|\hat{u}_{\lambda}(\xi)\right| \leq C_{k} \lambda{ }^{((n-1) / 2+k) \delta}(\lambda+|\xi|)^{-k}, \quad \xi / \lambda \notin U, \quad \lambda \gg 1, \tag{8-7}
\end{equation*}
$$

which gives any negative power of $\lambda$ for $k$ large enough, since $\delta<1$. If $V$ is bounded and $0 \notin \bar{V}$ then since $u_{\lambda}$ is uniformly bounded in $L^{2}$, we find

$$
\begin{equation*}
\int_{\lambda V}\left|\hat{u}_{\lambda}(\xi)\right|^{2}\left(1+|\xi|^{2}\right)^{-N} d \xi \leq C_{V} \lambda^{-2 N} \tag{8-8}
\end{equation*}
$$

which together with (8-7) gives (8-2). If $\chi \in C_{0}^{\infty}$ then we may apply (8-7) to $\chi u_{\lambda}$; thus we find for any positive integer $k$ that

$$
\begin{equation*}
\left|\widehat{\chi u}_{\lambda}(\xi)\right| \leq C \lambda^{((n-1) / 2+k) \delta}(\lambda+|\xi|)^{-k}, \quad \xi \in W, \quad \lambda \gg 1, \tag{8-9}
\end{equation*}
$$

if $W$ is any closed cone with ( $\operatorname{supp} \chi \times W) \cap \Gamma=\varnothing$. Thus we find that $\lambda^{k} u_{\lambda} \rightarrow 0$ in $\mathcal{D}_{\Gamma}^{\prime}$ for every $k$. To prove (8-3) we may assume that $x_{0}=0$ and take $\psi \in C_{0}^{\infty}$. If $\operatorname{Im} \omega(0)=0$ and $\phi_{0}(0) \neq 0$ then since $\delta<1$,
we obtain

$$
\begin{align*}
\lambda^{n-(n-1) \delta / 2} e^{-i \lambda \operatorname{Re} w(0)}\left\langle u_{\lambda}, \psi(\lambda \cdot)\right\rangle & =\int e^{i \lambda(w(x / \lambda)-\operatorname{Re} w(0))} \psi(x) \sum_{j} \phi_{j}\left(\lambda^{\delta-1} x\right) \lambda^{-j \delta} d x \\
& \rightarrow \int e^{i\left(\operatorname{Re} \partial_{x} \omega(0), x\right\rangle} \psi(x) \phi_{0}(0) d x, \quad \lambda \rightarrow \infty, \tag{8-10}
\end{align*}
$$

which is not equal to zero for some suitable $\psi \in C_{0}^{\infty}$. Since

$$
\begin{equation*}
\|\psi(\lambda \cdot)\|_{(N)} \leq C_{N} \lambda^{N-n / 2} \tag{8-11}
\end{equation*}
$$

we obtain from (8-10) that $0<c \leq \lambda^{N+n / 2-(n-1) \delta / 2}\left\|u_{\lambda}\right\|_{(-N)}$, which gives (8-3) and the lemma.
Proof of Theorem 2.7. By conjugating with elliptic Fourier integral operators and multiplying with pseudodifferential operators, we obtain that $P^{*} \in \Psi_{\mathrm{cl}}^{2}$ is of the form given by Proposition 3.2 microlocally near $\Gamma=\left\{\left(t, x_{0}, y_{0}, 0, \xi_{0}, 0\right): t \in I\right\}$. Thus we may assume

$$
\begin{equation*}
P^{*}=D_{t}+F\left(t, x, y, D_{t}, D_{x}, D_{y}\right)+R \tag{8-12}
\end{equation*}
$$

where $R \in \Psi_{\mathrm{cl}}^{2}$ satisfies $\mathrm{WF}(R) \cap \Gamma=\varnothing$.
Now we can construct approximate solutions $u_{\lambda}$ of the form (3-14) by using the expansion (3-15). By reducing to minimal bicharacteristics, we may solve first the eikonal equation by using Proposition 5.2 and then the transport equations (6-20) by using Proposition 6.3 with $\varrho=\lambda^{1 / N}$ for $N>3$. Thus after making the change of coordinates $(t, x, y) \mapsto\left(t-t_{0}, x-x_{0}(t), y-y_{0}\right)$, we obtain approximate solutions $u_{\lambda}$ of the form (8-1) in Lemma 8.1 with $\varrho=1 / N$ and $\delta=3 / N$. For $N$ large enough, we may choose $K$ and $M$ in Proposition 6.3 so that $\left|\left(D_{t}+F\right) u_{\lambda}\right| \lesssim \lambda^{-k}$ for any $k$. Now differentiation of $\left(D_{t}+F\right) u_{\lambda}$ can at most give a factor $\lambda$ since $\delta<1$, and a loss of a factor $x-x_{0}(t)$ gives at most a factor $\lambda^{1 / 2}$. Because of the bounds on the support of $u_{\lambda}$, we may obtain

$$
\begin{equation*}
\left\|\left(D_{t}+F\right) u_{\lambda}\right\|_{(\nu)}=\mathcal{O}\left(\lambda^{-N-n}\right) \tag{8-13}
\end{equation*}
$$

for any chosen $v$. Since $\phi_{0}\left(t_{0}, x_{0}\left(t_{0}\right), y_{0}\right)=1$ by Proposition 6.3 and $\operatorname{Im} w\left(t_{0}, x_{0}\left(t_{0}\right)\right)=0$ by Proposition 5.2, we find by (8-2)-(8-3) that

$$
\begin{equation*}
\lambda^{-N-n / 2} \ll \lambda^{-N-n / 2+(n-1) \delta / 2} \lesssim\|u\|_{(-N)} \lesssim \lambda^{-N} \quad \forall N, \quad \lambda \gg 1 \tag{8-14}
\end{equation*}
$$

Since $u_{\lambda}$ has support in a fixed compact set that shrinks towards $\left\{\left(t, x_{0}(t), y_{0}\right): t \in I\right\}$ as $\lambda \rightarrow \infty$, we find from Lemma 8.1 that $\|R u\|_{(\nu)}$ and $\|A u\|_{(0)}$ are $\mathcal{O}\left(\lambda^{-N-n}\right)$ if $\operatorname{WF}(A)$ does not intersect $\Gamma$. Thus we find from (8-13) and (8-14) that (2-17) does not hold when $\lambda \rightarrow \infty$, so $P$ is not solvable at $\Gamma$ by Remark 2.10.

## References

[Cardoso and Trèves 1974] F. Cardoso and F. Trèves, "A necessary condition of local solvability for pseudo-differential equations with double characteristics", Ann. Inst. Fourier (Grenoble) 24:1 (1974), 225-292. MR Zbl
[Colombini et al. 2003] F. Colombini, L. Pernazza, and F. Trèves, "Solvability and nonsolvability of second-order evolution equations", pp. 111-120 in Hyperbolic problems and related topics, edited by F. Colombini and T. Nishitani, International Press, Somerville, MA, 2003. MR
[Colombini et al. 2010] F. Colombini, P. D. Cordaro, and L. Pernazza, "Local solvability for a class of evolution equations", $J$. Funct. Anal. 258:10 (2010), 3469-3491. MR Zbl
[Dencker 2006] N. Dencker, "The resolution of the Nirenberg-Treves conjecture", Ann. of Math. (2) 163:2 (2006), 405-444. MR Zbl
[Dencker 2016] N. Dencker, "Solvability and limit bicharacteristics", J. Pseudo-Differ. Oper. Appl. 7:3 (2016), 295-320. MR Zbl
[Egorov 1977] J. V. Egorov, "Solvability conditions for equations with double characteristics", Dokl. Akad. Nauk SSSR 234:2 (1977), 280-282. In Russian; translated in Sov. Math. Dokl. 18 (1977), 632-635. MR Zbl
[Gilioli and Trèves 1974] A. Gilioli and F. Trèves, "An example in the solvability theory of linear PDE's", Amer. J. Math. 96 (1974), 367-385. MR Zbl
[Goldman 1975] R. Goldman, "A necessary condition for the local solvability of a pseudodifferential equation having multiple characteristics", J. Differential Equations 19:1 (1975), 176-200. MR Zbl
[Hörmander 1977] L. Hörmander, "The Cauchy problem for differential equations with double characteristics", J. Analyse Math. 32 (1977), 118-196. MR Zbl
[Hörmander 1981] L. Hörmander, "Pseudodifferential operators of principal type", pp. 69-96 in Singularities in boundary value problems (Maratea, Italy, 1980), edited by H. G. Garnir, NATO Adv. Study Inst. Ser. C: Math. Phys. Sci. 65, Reidel, Dordrecht, 1981. MR
[Hörmander 1983] L. Hörmander, The analysis of linear partial differential operators, I: Distribution theory and Fourier analysis, Grundlehren der Mathematischen Wissenschaften 256, Springer, Berlin, 1983. MR Zbl
[Hörmander 1985a] L. Hörmander, The analysis of linear partial differential operators, III: Pseudodifferential operators, Grundlehren der Mathematischen Wissenschaften 274, Springer, Berlin, 1985. MR Zbl
[Hörmander 1985b] L. Hörmander, The analysis of linear partial differential operators, IV: Fourier integral operators, Grundlehren der Mathematischen Wissenschaften 275, Springer, Berlin, 1985. MR Zbl
[Mendoza 1984] G. A. Mendoza, "A necessary condition for solvability for a class of operators with involutive double characteristics", pp. 193-197 in Microlocal analysis (Boulder, CO, 1983), Contemp. Math. 27, Amer. Math. Soc., Providence, RI, 1984. MR Zbl
[Mendoza and Uhlmann 1983] G. A. Mendoza and G. A. Uhlmann, "A necessary condition for local solvability for a class of operators with double characteristics", J. Funct. Anal. 52:2 (1983), 252-256. MR Zbl
[Mendoza and Uhlmann 1984] G. A. Mendoza and G. A. Uhlmann, "A sufficient condition for local solvability for a class of operators with double characteristics", Amer. J. Math. 106:1 (1984), 187-217. MR Zbl
[Nishitani 2004] T. Nishitani, "Effectively hyperbolic Cauchy problem", pp. 363-449 in Phase space analysis of partial differential equations, vol. 2, Scuola Norm. Sup., Pisa, 2004. MR Zbl
[Popivanov 1974] P. R. Popivanov, "The local solvability of a certain class of pseudodifferential equations with double characteristics", C. R. Acad. Bulgare Sci. 27 (1974), 607-609. In Russian. MR
[Trèves 1973] F. Trèves, "Concatenations of second-order evolution equations applied to local solvability and hypoellipticity", Comm. Pure Appl. Math. 26 (1973), 201-250. MR Zbl
[Trèves 1980] F. Trèves, Introduction to pseudodifferential and Fourier integral operators, vol. 2: Fourier integral operators, Plenum Press, New York, 1980. MR Zbl
[Wenston 1977] P. R. Wenston, "A necessary condition for the local solvability of the operator $P_{m}^{2}(x, D)+P_{2 m-1}(x, D)$ ", $J$. Differential Equations 25:1 (1977), 90-95. MR Zbl
[Wenston 1978] P. R. Wenston, "A local solvability result for operators with characteristics having odd order multiplicity", $J$. Differential Equations 28:3 (1978), 369-380. MR Zbl
[Wittsten 2012] J. Wittsten, "On some microlocal properties of the range of a pseudodifferential operator of principal type", Anal. PDE 5:3 (2012), 423-474. MR Zbl
[Yamasaki 1980] A. Yamasaki, "On a necessary condition for the local solvability of pseudodifferential operators with double characteristics", Comm. Partial Differential Equations 5:3 (1980), 209-224. MR Zbl
[Yamasaki 1983] A. Yamasaki, "On the local solvability of $D_{1}^{2}+A\left(x_{2}, D_{2}\right)$ ", Math. Japon. 28:4 (1983), 479-485. MR Zbl Received 23 Nov 2015. Revised 18 Sep 2016. Accepted 1 Nov 2016.

NILS DENCKER: dencker@maths.lth.se
Centre for Mathematical Sciences, University of Lund, Box 118, SE-221 00 Lund, Sweden

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