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 Journal of MathematicsTHE UNITARY DUAL OF $p$-ADIC $\widetilde{\operatorname{Sp(2)}}$<br>Marcela Hanzer and Ivan Matić

# THE UNITARY DUAL OF $p$-ADIC $\widetilde{\mathbf{S p ( 2 )}}$ 

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#### Abstract

We investigate the composition series of the induced admissible representations of the metaplectic group $\widetilde{\mathbf{S p ( 2 )}}$ over a $p$-adic field $F$. In this way, we determine the nonunitary and unitary duals of $\widetilde{\mathbf{S p ( 2 )}}$ modulo cuspidal representations.


## 1. Introduction

The admissible representations of reductive groups over $p$-adic fields have been studied intensively by many authors, but knowledge about the unitary dual of such groups is still incomplete. Besides some results concerning specific parts of the unitary dual of some classical and exceptional groups (that is, spherical, generic [Lapid et al. 2004] and so on), there are some situations where, for some low rank groups, the complete unitary dual is described [Sally and Tadić 1993; Muić 1997; Hanzer 2006; Matić 2010].

In this paper, we completely describe the noncuspidal unitary dual of the double cover of the symplectic group of split rank two. Although this is not an algebraic group, some recent results enabled us to study this group in the same spirit as the classical split groups. More concretely, Hanzer and Muić [2009] related reducibilities of the induced representations of metaplectic groups with those of the odd orthogonal groups (using theta correspondence), while their paper [2010] describes the extension of the Jacquet module techniques of Tadić for classical groups to metaplectic groups. More specifically, Tadić's structure formula for symplectic and odd-orthogonal groups [1995] (which is a version of a geometric lemma of [Bernstein and Zelevinsky 1977]) is extended to metaplectic groups. These ingredients made the determination of the irreducible subquotients of the principal series for $\widetilde{\mathrm{Sp}(2)}$ very similar to the one obtained in [Matić $\geq 2010$ ] for $\mathrm{SO}(5)$, but this happens to be insufficient tool in some cases. In these cases, we will use the theta correspondence to again obtain the formal similarity to the $\mathrm{SO}(5)$ case. This similarity was expected; see for example [Zorn 2010]. After determining complete nonunitary dual, modulo cuspidal representations, the unitary dual follows in the almost the same way as in [Matić 2010], but after discussion of some exceptional

[^0]cases (for example, the discussion of the unitary principal series): In the case of the odd orthogonal group $\mathrm{SO}(5)$, the irreducibility of the unitary principal series follows from the considerations about $R$-groups, and in the case of $\overparen{\operatorname{Sp}(2)}$, since the $R$-group theory for metaplectic groups is not available in its full generality, irreducibility is obtained using theta correspondence. In the forthcoming paper [Hanzer and Matić 2010], we extend the methods used here to prove for general $n$ the irreducibility of unitary principal series for $\widetilde{\mathrm{Sp}(n)}$. We hope these results will have applications in the theory of automorphic forms.

We now describe the content of the paper. In Section 2, we recall the definition of the metaplectic double cover $\widetilde{\mathrm{Sp}(n)}$. We also recall the notions of parabolic subgroups, Jacquet functor, and parabolic induction in the context of metaplectic groups. We then recall the notion of the dual pair, and the lifts of an irreducible representations of one member of the pair dual to the Weil representation of the ambient metaplectic group. We recall the criteria for the square integrability and temperedness of the irreducible representations of metaplectic groups, due to Ban and Jantzen [2009] and recall the classification of the irreducible genuine representations of $\widehat{\mathrm{Sp}(n)}$ obtained in [Hanzer and Muić 2010]. In Section 3, we analyze the principal series for $\widetilde{\mathrm{Sp}(2)}$, using both theta correspondence and Tadić's methods applied to metaplectic groups. In Section 4, we determine the unitary dual of $\widetilde{\mathrm{Sp}(2)}$ supported in the minimal parabolic subgroup. In Section 5, we describe irreducible representations of $\widehat{\mathrm{Sp}(2)}$ supported on maximal parabolic subgroups, and the unitary dual of $\widetilde{\mathrm{Sp}(2)}$ supported on maximal parabolic subgroups.

## 2. Preliminaries

Let $\widetilde{\mathrm{Sp}(2)}$ be the unique nontrivial two-fold central extension of symplectic group $\operatorname{Sp}(2, F)$, where $F$ is a non-Archimedean local field of characteristic different from two. In other words, we have maps

$$
1 \rightarrow \mu_{2} \rightarrow \widetilde{\mathrm{Sp}(2)} \rightarrow \mathrm{Sp}(2, F) \rightarrow 1
$$

The multiplication in $\widetilde{\mathrm{Sp}(2)}$, which is as a set given by $\operatorname{Sp}(2, F) \times \mu_{2}$, is given the cocycle of [Ranga Rao 1993]. The topology of $\widehat{\operatorname{Sp}(n)}$ is explained in detail in [Hanzer and Muić 2010, Section 3.3]. There exist compact open subgroups of $\mathrm{Sp}(n)$ that split in $\widetilde{\mathrm{Sp}(n)}$. Recall that a maximal good compact subgroup $\mathrm{Sp}\left(O_{F}\right)$ splits if the residual characteristic of $F$ is odd (here $O_{F}$ denotes the ring of integers on $F$ ). In [Hanzer and Muić 2010], the metaplectic group $\widetilde{\mathrm{Sp}(2)}$ was denoted by $\widetilde{\operatorname{Sp}\left(W_{2}\right)}$. We say that the representation of $\widetilde{\mathrm{Sp}(2)}$ (or, more generally, $\widetilde{\mathrm{Sp}(n))}$ is smooth if, for every vector $v$ in the representation space $V$, there exists a compact open subgroup $K_{1}$ of $\mathrm{Sp}(2)$ that splits in $\widetilde{\mathrm{Sp}(2)}$ and fixes $v$. The representation is admissible if for every $K_{1}$ as above, the space $V^{K_{1}}$ is finite-dimensional.

Lemma 2.1. Irreducible smooth representations of $\widetilde{\mathrm{Sp}(n)}$ are admissible.
Proof. First, we prove that an irreducible, smooth, cuspidal representation of $\widetilde{\operatorname{Sp}(n)}$ is admissible. We can proceed as in the corollary on [Bernstein 1992, page 36], so we have to prove that an irreducible cuspidal representation of $\widehat{\operatorname{Sp}(n)}$ is compact, and that was proved in proving [Hanzer and Muić 2009, Lemma 3.1]. Then, on this compact irreducible representation we can apply [Bernstein 1992, Proposition 11], which is formulated for a general totally disconnected group (so the metaplectic groups satisfy the conditions), and says that finitely generated compact representations are admissible. The claim follows since every irreducible smooth representation can be embedded in the representations parabolically induced from the cuspidal representations of Levi subgroups (and for the representations of Levi subgroups, the same reasoning as above shows that these representations are also admissible), which was proved in [Hanzer and Muić 2008, Proposition 4.4].

In this paper we are interested only in genuine representations of $\widetilde{\operatorname{Sp}(n)}$ (that is, those that do not factor through $\left.\mu_{2}\right)$. So, let $R(n)$ be the Grothendieck group of the category of all admissible genuine representations of finite length of $\overline{\mathrm{Sp}(n)}$ (that is, a free abelian group over the set of all irreducible genuine representations of $\widetilde{\mathrm{Sp}(n)})$, and define $R=\bigoplus_{n \geq 0} R(n)$. By $v$ we denote a character of $\operatorname{GL}(k, F)$ defined by $\mid \operatorname{det}^{F}$. Further, for an ordered partition $s=\left(n_{1}, n_{2}, \ldots, n_{j}\right)$ of some $m \leq n$, we denote by $P_{s}$ a standard parabolic subgroup of $\operatorname{Sp}(n, F)$ (consisting of block upper-triangular matrices), whose Levi factor equals $\mathrm{GL}\left(n_{1}\right) \times \mathrm{GL}\left(n_{2}\right) \times \cdots \times$ $\operatorname{GL}\left(n_{j}\right) \times \operatorname{Sp}(n-|s|, F)$, where $|s|=\sum_{i=1}^{j} n_{i}$. By a standard parabolic subgroup $\widetilde{P}_{s}$ of $\overline{\mathrm{Sp}(n)}$ we mean the preimage of $P_{s}$ in $\widetilde{\mathrm{Sp}(n)}$. We have the analogous notation for the Levi subgroups of the metaplectic groups, and, for the completeness, we explicitly describe the structure of the parabolic and Levi subgroups, as explained in [Hanzer and Muić 2010, Section 2.2]. There is a natural splitting from the unipotent radical of $N_{s}$ of the corresponding standard parabolic subgroup $P_{s}$ of $\operatorname{Sp}(n, F)$ to its cover [Mœglin et al. 1987, Lemma 2.9 on page 43]; let $N_{s}^{\prime}$ be the image of that homomorphism. We then have $\widetilde{P}_{s} \cong \widetilde{M}_{s} \ltimes N_{s}^{\prime}$.

We can explicitly describe $\widetilde{M}_{s}$ as follows. There is a natural epimorphism

$$
\begin{equation*}
\phi: \widetilde{\mathrm{GL}\left(n_{1}, F\right)} \times \cdots \times \widetilde{\mathrm{GL}\left(n_{k}, F\right)} \times \mathrm{Sp} \widetilde{\left(W_{n-|s|}\right)} \rightarrow \widetilde{M}_{s} \tag{1}
\end{equation*}
$$

given by

$$
\begin{equation*}
\left(\left[g_{1}, \epsilon_{1}\right], \ldots,\left[g_{k}, \epsilon_{k}\right],[h, \epsilon]\right) \mapsto\left[\left(g_{1}, g_{2}, \ldots, g_{k}, h\right), \epsilon_{1} \cdots \epsilon_{k} \epsilon \beta\right] \tag{2}
\end{equation*}
$$

with $\beta=\prod_{i<j}\left(\operatorname{det} g_{i}, \operatorname{det} g_{j}\right)_{F}\left(\prod_{i=1}^{k}\left(\operatorname{det} g_{i}, x(h)\right)_{F}\right)$, where $x(h)$ is defined in [Ranga Rao 1993, Lemma 5.1] and $(\cdot, \cdot)_{F}$ denotes the Hilbert symbol of the field $F$. Although $\widetilde{M}$ is not exactly the product at left in (1), it differs from it by a finite subgroup that enables us to write every irreducible representation $\pi$
of $\tilde{M}$ in the form $\pi_{1} \otimes \cdots \otimes \pi_{k} \otimes \pi^{\prime}$, where the representations $\pi_{1}, \ldots, \pi_{k}, \pi^{\prime}$ are either all genuine or none genuine. This simple property enables us to set up Tadić's machinery [Tadić 1995; Hanzer and Muić 2008] of parabolic induction and Jacquet functors. Recall that the irreducible representations in this paper, unless mentioned otherwise, are assumed to be genuine (that is, nontrivial on $\mu_{2}$ ). Also, the cuspidality of representations is defined in the same way as for the reductive groups (because of the splitting of the unipotent radical) and characterized in terms if the support of the matrix coefficients also as for the reductive groups.

Let $\sigma$ be a representation of $\widetilde{\operatorname{Sp}(2)}$. Following the notation introduced in [Hanzer and Muić 2010], we denote by $R_{\widetilde{P_{(1,1)}}}(\sigma)$ the normalized Jacquet module with respect to $\widetilde{M}_{(1,1)}$; by $R_{\widetilde{P}_{1}}(\sigma)$ the normalized Jacquet module with respect to $\widetilde{M}_{(1)}$; and by $R_{\widetilde{P}_{2}}(\sigma)$ the normalized Jacquet module with respect to $\widetilde{M}_{(2)}$.

We fix a nontrivial additive character $\psi$ of $F$ and let $\omega_{n, r}$ be the pullback of the Weil representation $\omega_{n(2 r+1), \psi}$ of the group $\operatorname{Sp}(\widetilde{n(2 r+1))}$, restricted to the dual pair $\widehat{\mathrm{Sp}(n)} \times O(2 r+1)$ [Kudla 1996, Chapter II]. Here $O(2 r+1)$ denotes the split odd-orthogonal group of the split rank $r$, with the one-dimensional anisotropic space sitting at the bottom of the orthogonal tower [Kudla 1996, Chapter III.1]. The standard parabolic subgroups (containing the upper triangular Borel subgroup) of $O(2 r+1)$ have the analogous description as the standard parabolic subgroups of $\operatorname{Sp}(n, F)$; we use the analogous notation for the normalized Jacquet functors.

Let $\sigma$ be an irreducible smooth genuine representation of $\overline{\mathrm{Sp}(n)}$. We write $\Theta(\sigma, r)$ for the smooth isotypic component of $\sigma$ in $\omega_{n, r}$ (we view it as a representation of $O(2 r+1)$ ). Denote by $r_{0}$ the smallest $r$ such that $\Theta(\sigma, r) \neq 0$. When $\sigma$ is cuspidal, we know that $\Theta\left(\sigma, r_{0}\right)$ is an irreducible cuspidal representation of $O(2 r+1)$.

Let $\mathrm{GL}(n, F)$ be a double cover of $\mathrm{GL}(n, F)$, where the multiplication is given by

$$
\left(g_{1}, \epsilon_{1}\right)\left(g_{2}, \epsilon_{2}\right)=\left(g_{1} g_{2}, \epsilon_{1} \epsilon_{2}\left(\operatorname{det} g_{1}, \operatorname{det} g_{2}\right)_{F}\right)
$$

Here $\epsilon_{i} \in \mu_{2}$ for $i=1,2$ and $(\cdot, \cdot)_{F}$ denotes the Hilbert symbol of the field $F$, and this cocycle on $\mathrm{GL}(n, F)$ is actually a restriction of Ranga Rao's cocyle on $\operatorname{Sp}(n, F)$ to $\operatorname{GL}(n, F)$, if we view this group as the Siegel Levi subgroup of $\operatorname{Sp}(n, F)$ [Kudla 1986, page 235]. Now we fix a character $\chi_{V, \psi}(g, \epsilon)=$ $\chi_{V}(\operatorname{det} g) \epsilon \gamma\left(\operatorname{det} g, \psi_{1 / 2}\right)^{-1}$ of G $\widetilde{\operatorname{LL}(n, F)}$. Here $\gamma$ denotes the Weil invariant, while $\chi_{V}$ is a character related to the quadratic form on $O(2 r+1)$ [Kudla 1996, pages 17 and 37], and $\psi_{a}(x)=\psi(a x)$ for $a \in F^{*}$. We may suppose $\chi_{V} \equiv 1$ (but the arguments that follow are valid without this assumption). We write $\alpha=\chi_{V, \psi}^{2}$ and observe that $\alpha$ is a quadratic character on $\operatorname{GL}(n, F)$.

The following fact, follows directly from [Hanzer and Muić 2010], and we use it frequently while determining composition series of induced representations: For
an irreducible genuine representation $\pi$ of $\widetilde{\mathrm{GL}(k, F)}$ and an irreducible genuine representation $\sigma$ of $\widetilde{\operatorname{Sp}(n)}$ we have $\pi \rtimes \sigma=\tilde{\pi} \alpha \rtimes \sigma$ (in $R$ ), where $\pi \rtimes \sigma$ denotes the representation of the group $\mathrm{Sp}(n+k)$ parabolically induced from the representation $\pi \otimes \sigma$ of the maximal Levi subgroup $\widetilde{M}_{(k)}$. We follow here the usual notation for parabolic induction for classical groups, adapted to the metaplectic case [Tadić 1994; Hanzer and Muić 2010]. We also freely use Zelevinsky's notation [1980] for the parabolic induction for general linear groups. We denote the Steinberg representation of the reductive algebraic group $G$ by $\mathrm{St}_{G}$ and the trivial representation of that group by $1_{G}$. Following [Kudla 1996], we let $\omega_{\psi_{a}, n}^{+}$denote the even part of the Weil representation of $\widetilde{\operatorname{Sp}(n)}$ determined by the additive character $\psi_{a}$. The nontrivial character of $\mu_{2}$, when we view it as a representation of $\widetilde{\mathrm{Sp}(0)}$, is denoted by $\omega_{0}$.

If $\zeta$ is a quadratic character of $F^{\times}$, we can write $\zeta(x)=(x a)_{F}$ for some $a \in F^{\times}$. Let $s p_{\zeta, 1}$ be an irreducible (square-integrable, according to the criterion for the square-integrability which we recall below) subrepresentation of $\chi_{V, \psi} \zeta \nu^{1 / 2} \rtimes \omega_{0}$. Then, as in [Kudla 1996, page 89], we have the exact sequence

$$
0 \rightarrow s p_{\zeta, 1} \longrightarrow \chi_{V, \psi} \zeta \nu^{1 / 2} \rtimes \omega_{0} \longrightarrow \omega_{\psi_{a}, 1}^{+} \rightarrow 0
$$

The results of [Ban and Jantzen 2009] imply that Casselman's criteria for squareintegrability and temperedness hold for metaplectic groups in a similar form as for the classical groups (for example symplectic). We now recall these criteria.

Let $\pi$ be an admissible irreducible (genuine) representation of $\widetilde{\mathrm{Sp}(n)}$ and let $\widetilde{P}_{s}$ be any standard parabolic subgroup minimal with respect to the property that $R_{\widetilde{P}_{s}}(\pi) \neq 0$. Write $s=\left(n_{1}, \ldots, n_{k}\right)$ and let $\sigma$ be any irreducible subquotient of $R_{\widetilde{P}_{s}}(\pi)$. As we saw above, we can write $\sigma=\rho_{1} \otimes \rho_{2} \otimes \cdots \otimes \rho_{k} \otimes \rho$, where $\rho_{i}$ is an irreducible genuine cuspidal representation of some $\operatorname{GL(n_{i},F)}$ for $i=1, \ldots, k$ and $\rho$ is an irreducible genuine cuspidal representation of some $\widetilde{\mathrm{Sp}(n-l)}$. Define $e\left(\rho_{i}\right)$ by $\rho_{i}=v^{e\left(\rho_{i}\right)} \rho_{i}^{u}$, where $\rho_{i}^{u}$ is unitary for $1 \leq i \leq n$.

Assume that the inequalities

$$
\begin{aligned}
n_{1} e\left(\rho_{1}\right) & >0, \\
n_{1} e\left(\rho_{1}\right)+n_{2} e\left(\rho_{2}\right) & >0, \\
& \vdots \\
n_{1} e\left(\rho_{1}\right)+n_{2} e\left(\rho_{2}\right)+\cdots+n_{k} e\left(\rho_{k}\right) & >0 .
\end{aligned}
$$

hold for every $s$ and $\sigma$ as above. Then $\pi$ is a square integrable representation. For such $s$ and $\sigma$, these inequalities also hold if $\pi$ is a square integrable representation.

The criterion for tempered representations is given by replacing every $>$ with $\geq$.
We recall the definition of a negative representation [Hanzer and Muić 2010, Definition 4.1].

Let $\sigma$ be an admissible irreducible genuine representation of $\widetilde{\mathrm{Sp}(n)}$. Then $\sigma$ is a strongly negative (respectively, negative) representation if and only if for every embedding $\sigma \hookrightarrow \rho_{1} \times \rho_{2} \times \cdots \times \rho_{k} \rtimes \rho$, where $\rho_{i}$ for $1 \leq i \leq k$ and $\rho$ are irreducible genuine supercuspidal representations of some of the $\widetilde{\mathrm{GL}}$ and of some $\widetilde{\mathrm{Sp}(n-l)}$, we have

$$
\begin{array}{rlrl}
n_{1} e\left(\rho_{1}\right) & <0 & & (\text { respectively, } \leq 0), \\
n_{1} e\left(\rho_{1}\right)+n_{2} e\left(\rho_{2}\right) & <0 & \text { (respectively, } \leq 0), \\
& \vdots & \\
n_{1} e\left(\rho_{1}\right)+n_{2} e\left(\rho_{2}\right)+\cdots+n_{k} e\left(\rho_{k}\right)<0 & \text { (respectively, } \leq 0) .
\end{array}
$$

As soon as $\sigma$ as above is genuine, the $\rho_{i}$ and $\rho$ are also necessarily genuine. For notation, we recall [Hanzer and Muić 2010, Theorems 4.5 and 4.6]. Recall that, for a cuspidal representation $\rho$ of some $\operatorname{GL}\left(m_{\rho}, F\right)$, a segment $\Delta$ is a set of cuspidal representations $\Delta=\left\{\rho, v \rho, \ldots, v^{k-1} \rho\right\}$ and $\langle\Delta\rangle$ is a unique irreducible subrepresentation of $\rho \times v \rho \times \cdots \times v^{k-1} \rho$. We use the same notation for genuine cuspidal representations of $\operatorname{GL}\left(m_{\rho}, F\right)$ since the transfer from nongenuine to genuine representations in the case of $\mathrm{GL}\left(m_{\rho}, F\right)$ is particularly simple (obtained by multiplication with the character $\chi_{V, \psi}(g, \epsilon)$ defined above). Now the two theorems above follow from the analogous results in the case of classical reductive groups of [Hanzer and Muić 2008], since the analogous calculations with Jacquet modules are possible, due to results in [Hanzer and Muić 2010].

- Suppose that $\Delta_{1}, \ldots, \Delta_{k}$ is a sequence of segments (of genuine representations) such that $e\left(\Delta_{1}\right) \geq \cdots \geq e\left(\Delta_{k}\right)>0$ (we also allow $k=0$ ). Let $\sigma_{\text {neg }}$ be a negative (genuine) representation. Then the induced representation $\left\langle\Delta_{1}\right\rangle \times\left\langle\Delta_{2}\right\rangle \times \cdots \times\left\langle\Delta_{k}\right\rangle \rtimes \sigma_{\text {neg }}$ has a unique irreducible subrepresentation; we denote it by $\left\langle\Delta_{1}, \ldots, \Delta_{k} ; \sigma_{\text {neg }}\right\rangle$.
- If $\sigma$ is an irreducible admissible genuine representation of $\widetilde{\operatorname{Sp}(n)}$, then there exist a sequence of segments (of genuine representations) $\Delta_{1}, \ldots, \Delta_{k}$ such that $e\left(\Delta_{1}\right) \geq \cdots \geq e\left(\Delta_{k}\right)>0$ and a negative (genuine) representation $\sigma_{\text {neg }}$ such that $\sigma \simeq\left\langle\Delta_{1}, \ldots, \Delta_{k} ; \sigma_{\text {neg }}\right\rangle$.

We can carry over Tadić's structure formula for classical groups to the metaplectic case [Hanzer and Muić 2010, Proposition 4.5], which enables us to calculate Jacquet modules of the induced representations. In more detail, let

$$
\left.R^{\mathrm{gen}}=\bigoplus_{n} R(\widetilde{\mathrm{GL}(n, F)})\right)_{\mathrm{gen}},
$$

where $R(\widetilde{\mathrm{GL}(n, F)})$ gen denotes the Grothendieck group of finite length, smooth, genuine representations of $\widehat{\operatorname{GL(n,F)}}$. We denote by $\times$ the linear extension to
$R^{\text {gen }} \otimes R^{\text {gen }}$ of the parabolic induction (from a maximal parabolic subgroup). We can easily check that if $\sigma$ is an irreducible genuine representation of $\overline{\operatorname{Sp}\left(W_{n}\right)}$, then $r_{k}(\sigma)$, the normalized Jacquet module of $\sigma$ with respect to the standard maximal parabolic $\widetilde{P}_{k}$, is a genuine representation of $\widetilde{M_{(k)}}$ and as such can be interpreted as a (genuine) representation of $\mathrm{GL(k,F)} \times \widehat{\mathrm{Sp}\left(W_{n-k}\right)}$, that is, as an element of $R^{\text {gen }} \otimes R$, with $R$ defined as above. So for irreducible genuine $\sigma$, we can introduce $\mu^{*}(\sigma) \in R^{\text {gen }} \otimes R$ by

$$
\mu^{*}(\sigma)=\sum_{k=0}^{n} \text { s.s. }\left(r_{k}(\sigma)\right)
$$

where s.s. stands for semisimplification. We can extend $\mu^{*}$ linearly to the whole $R$. Using Jacquet modules for the maximal parabolic subgroups of G $\widetilde{\mathrm{L}(n, F)}$ we can analogously define

$$
m^{*}(\pi)=\sum_{k=0}^{n} \text { s.s. }\left(r_{k}(\pi)\right) \in R^{\mathrm{gen}} \otimes R^{\mathrm{gen}}
$$

for a genuine, irreducible representation $\pi$ of $\widehat{\mathrm{GL}(n, F)}$ and then extend $m^{*}$ linearly to the whole $R^{\text {gen }}$. Let $\kappa: R^{\text {gen }} \otimes R^{\text {gen }} \rightarrow R^{\text {gen }} \otimes R^{\text {gen }}$ be defined by $\kappa(x \otimes y)=y \otimes x$. We extend the contragredient ${ }^{\sim}$ to an automorphism of $R^{\text {gen }}$ naturally. Finally, we define

$$
M^{*}=(m \otimes \mathrm{id}) \circ\left(\sim \alpha \otimes m^{*}\right) \circ \kappa \circ m^{*} .
$$

Here ${ }^{\sim} \alpha$ means taking contragredient of a representation, and then multiplying by the character $\alpha$, acting on the general linear group as $\alpha(g)=(\operatorname{det} g,-1)_{F}$.

For $\pi$ in $R^{\text {gen }}$ and $\sigma$ from $R$, we have

$$
\mu^{*}(\pi \rtimes \sigma)=M^{*}(\pi) \rtimes \mu^{*}(\sigma)
$$

Using this formula for the induced representations of $\widetilde{\mathrm{Sp}(2)}$, we get the following:

- Fix an admissible representation $\pi$ of $\widetilde{\mathrm{GL}(2)}$, and suppose that $\pi$ is of finite length. Let $m^{*}(\pi)=1 \otimes \pi+\sum_{i} \pi_{i}^{1} \otimes \pi_{i}^{2}+\pi \otimes 1$, where $\sum_{i} \pi_{i}^{1} \otimes \pi_{i}^{2}$ is a decomposition into a sum of irreducible representations. Now we have

$$
\begin{aligned}
\mu^{*}\left(\pi \rtimes \omega_{0}\right)=1 \otimes \pi \rtimes \omega_{0}+\sum_{i} \pi_{i}^{1} \otimes & \pi_{i}^{2} \rtimes \omega_{0}+\sum_{i} \alpha \widetilde{\pi}_{i}^{2} \otimes \pi_{i}^{1} \rtimes \omega_{0} \\
& +\pi \otimes \omega_{0}+\alpha \widetilde{\pi} \otimes \omega_{0}+\sum_{i} \pi_{i}^{1} \times \alpha \widetilde{\pi}_{i}^{2} \otimes \omega_{0}
\end{aligned}
$$

- Fix an admissible representation $\pi$ of $\widetilde{\text { GL(1) }}$ and an admissible representation $\sigma$ of $\widetilde{\operatorname{Sp}(1)}$. If we have

$$
\mu^{*}(\sigma)=1 \otimes \sigma+\sum_{i} \sigma_{i}^{1} \otimes \sigma_{i}^{2}
$$

where $\sigma_{i}^{1}$ and $\sigma_{i}^{2}$ are irreducible representations, then

$$
\begin{aligned}
\mu^{*}(\pi \rtimes \sigma)=1 \otimes \pi \rtimes & \sigma+\pi \otimes \sigma+\alpha \tilde{\pi} \otimes \sigma \\
& +\sum_{i} \sigma_{i}^{1} \otimes \pi \rtimes \sigma_{i}^{2}+\sum_{i} \pi \times \sigma_{i}^{1} \otimes \sigma_{i}^{2}+\sum_{i} \sigma_{i}^{1} \times \alpha \tilde{\pi} \otimes \sigma_{i}^{2}
\end{aligned}
$$

From now on, $\widehat{F^{\times}}$denotes the set of the unitary characters of $F^{\times}$, while $\widetilde{F^{\times}}$ denotes those that are not necessarily unitary.

## 3. Principal series

We first state an important reducibility result that follows directly from [Hanzer and Muić 2009, Theorems 3.5. and 4.2].

Proposition 3.1. Let $\chi \in \widehat{F^{\varnothing}}$ and let $s \in \mathbb{R}$ be nonnegative. The representation $\chi_{V, \psi} \nu^{s} \chi \rtimes \omega_{0}$ of $\widetilde{\mathrm{Sp}(1)}$ reduces if and only if $\chi^{2}=1_{F^{\times}}$and $s=1 / 2$.

Let $\zeta \in \widehat{F^{×}}$ such that $\zeta^{2}=1_{F^{\times}}$. In $R$ we have (see [Kudla 1996, page 89]

$$
\chi_{V, \psi} v^{1 / 2} \zeta \rtimes \omega_{0}=s p_{\zeta, 1}+\omega_{\psi_{a}, 1}^{+}
$$

The following proposition is well known and follows easily from the analogous results for the split $\mathrm{SO}(3)$ and $\mathrm{SO}(5)$.

Proposition 3.2. (1) Let $\chi \in \widehat{F^{\varnothing}}$ and suppose $s \in \mathbb{R}$ is nonnegative. The representation $\nu^{s} \chi \rtimes 1$ of $O(3)$ reduces if and only if $\chi^{2}=1_{F^{\times}}$and $s=1 / 2$. In that situation, the length of $\nu^{1 / 2} \chi \rtimes 1$ is two, and this representation has the unique subrepresentation that is square integrable.
(2) Let $\zeta_{1}, \zeta_{2} \in \widehat{F^{×}}$. Then, the unitary principal series $\zeta_{1} \times \zeta_{2} \rtimes 1$ of $O(5)$ is irreducible.

We use these two propositions in the sequel without explicitly mentioning them.
3.1. Unitary principal series. In this subsection we prove irreducibility of the unitary principal series $\chi_{V, \psi} \chi_{1} \times \chi_{V, \psi} \chi_{2} \rtimes \omega_{0}$, where $\chi_{i} \in \widehat{F^{\times}}$for $i=1,2$.

Let $\Pi$ denote the representation $\chi_{V, \psi} \chi_{1} \times \chi_{V, \psi} \chi_{2} \rtimes \omega_{0}$. Using the structure formula for $\mu^{*}(\Pi)$ from the end of the previous section, we get

$$
\begin{aligned}
R_{\widetilde{P}_{1}}(\Pi)=\chi_{V, \psi} \chi_{1}^{-1} \otimes \chi_{V, \psi} & \chi_{2} \rtimes \omega_{0}+\chi_{V, \psi} \chi_{1} \otimes \chi_{V, \psi} \chi_{2} \rtimes \omega_{0} \\
& +\chi_{V, \psi} \chi_{2}^{-1} \otimes \chi_{V, \psi} \chi_{1} \rtimes \omega_{0}+\chi_{V, \psi} \chi_{2} \otimes \chi_{V, \psi} \chi_{1} \rtimes \omega_{0}
\end{aligned}
$$

Remark. Let $\pi$ be an irreducible subrepresentation of $\Pi$. Because of irreducibility of the representations $\chi_{V, \psi} \chi_{1} \times \chi_{V, \psi} \chi_{2}$ and $\chi_{V, \psi} \chi_{i} \rtimes \omega_{0}$ for $i=1$, 2 , we get

$$
\pi \hookrightarrow \Pi \simeq \chi_{V, \psi} \chi_{1}^{-1} \times \chi_{V, \psi} \chi_{2} \rtimes \omega_{0} \simeq \chi_{V, \psi} \chi_{2}^{-1} \times \chi_{V, \psi} \chi_{1} \rtimes \omega_{0}
$$

If $\chi_{i} \neq \chi_{i}^{-1}$ holds for both $i=1,2$ and $\chi_{1} \neq \chi_{2}^{ \pm 1}$, then Frobenius reciprocity implies that $R_{\widetilde{P}_{1}}(\pi)=R_{\widetilde{P}_{1}}(\Pi)$, so $\pi=\Pi$ and the representation $\Pi$ is irreducible.

Now we prove the irreducibility of the unitary principal series for general unitary characters. Let $\zeta_{1}, \zeta_{2}$ be the unitary characters of $F^{\times}$. We prove irreducibility of the representation $\chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}$ using theta correspondence, beginning with this lemma:

Lemma 3.3. Let $\pi_{1}$ be an irreducible subrepresentation of $\chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}$. Then $\Theta\left(\pi_{1}, 2\right)=\zeta_{1} \times \zeta_{2} \rtimes 1$.

Proof. According to the stable range condition [Kudla 1996, page 48], $\Theta\left(\pi_{1}, 4\right) \neq 0$ (observe that $\Theta\left(\pi_{1}, 4\right)$ is a smooth representation of $O(9)$ ). We have epimorphisms $\omega_{2,4} \rightarrow \pi_{1} \otimes \Theta\left(\pi_{1}, 4\right)$ and $R_{P_{1}}\left(\omega_{2,4}\right) \rightarrow \pi_{1} \otimes R_{P_{1}}\left(\Theta\left(\pi_{1}, 4\right)\right)$. If $\tau$ is an irreducible quotient of $\Theta\left(\pi_{1}, 4\right)$, then [Kudla 1986, Corollary 2.6] implies [ $\tau$ ] = [ $\left.v^{-3 / 2}, v^{-1 / 2}, \zeta_{1}, \zeta_{2} ; 1\right]$, where [ $\tau$ ] denotes the cuspidal support of $\tau$. Clearly, $R_{P_{(1,1,1,1)}}(\tau) \geq v^{l_{1} / 2} \otimes v^{l_{2} / 2} \otimes \zeta_{1}^{ \pm 1} \otimes \zeta_{2}^{ \pm 1}$ or $R_{P_{(1,1,1,1)}}(\tau) \geq \zeta_{1}^{ \pm 1} \otimes v^{l_{1} / 2} \otimes \zeta_{2}^{ \pm 1} \otimes v^{l_{2} / 2}$ (or we have some order of factors) for some $l_{1}, l_{2} \in\{ \pm 1, \pm 3\}$. If we assume that in the Jacquet module $R_{P_{(1,1,1,1)}}(\tau)$ there is an irreducible subquotient as above whose first factor consists of a unitary character, then, using [Bernstein 1992, Lemma 26] together with Frobenius reciprocity, easily follows that

$$
\operatorname{Hom}\left(\tau, \zeta_{1}^{ \pm 1} \times v^{l_{1} / 2} \times \zeta_{2}^{ \pm 1} \times v^{l_{2} / 2} \rtimes 1\right) \neq 0
$$

But since $\zeta_{i}^{ \pm 1} \times \nu^{l_{i} / 2} \cong \nu^{l_{i} / 2} \times \zeta_{i}^{ \pm 1}$, we have $\operatorname{Hom}\left(\tau, \nu^{l_{1} / 2} \times \zeta_{1}^{ \pm 1} \times \zeta_{2}^{ \pm 1} \times v^{l_{2} / 2} \rtimes 1\right) \neq 0$. So, there is an irreducible subquotient $\tau^{\prime}$ of $\zeta_{1} \times \zeta_{2} \times v^{l_{2} / 2} \rtimes 1$ such that $\tau$ is a subrepresentation of $\nu^{l_{1} / 2} \rtimes \tau^{\prime}$. This implies that $R_{P_{1}}(\tau)\left(\nu^{l_{1} / 2}\right)$, the isotypic component of $R_{P_{1}}(\tau)$ along the generalized character $\nu^{l_{1} / 2}$, is nonzero, as is $R_{P_{1}}\left(\Theta\left(\pi_{1}, 4\right)\right)\left(\nu^{l_{1} / 2}\right)$.

Observations above imply that there is an irreducible representation $\tau_{1}$ of $O$ (3) such that the mappings $R_{P_{1}}\left(\omega_{2,4}\right) \rightarrow \pi_{1} \otimes R_{P_{1}}\left(\Theta\left(\pi_{1}, 4\right)\right) \rightarrow \pi_{1} \otimes \nu^{l_{1} / 2} \otimes \tau_{1}$ are epimorphisms. We denote the epimorphism $\left.R_{P_{1}}\left(\omega_{2,4}\right)\right) \rightarrow \pi_{1} \otimes \nu^{l_{1} / 2} \otimes \tau_{1}$ by $T$. Now $R_{P_{1}}\left(\omega_{2,4}\right)$ has the filtration in which

- $I_{10}=v^{-3 / 2} \otimes \omega_{2,3}$ is the quotient and
- $I_{11}=\operatorname{Ind}_{\mathrm{GL}(1) \times \widetilde{P}_{1} \times O(3)}^{M_{1} \times \widetilde{\operatorname{Sp}(2)}}\left(\chi_{V, \psi} \Sigma_{1}^{\prime} \otimes \omega_{1,3}\right)$ is the subrepresentation.

See [Kudla 1996, page 57] and [Hanzer and Muić 2009, Proposition 3.3], where the notation is explained in detail.

Suppose $\left.T\right|_{I_{11}} \neq 0$. Because $\chi_{V, \psi} \nu^{-l_{1} / 2}$ is the isotypic component of $\nu^{l_{1} / 2}$ in the $\widehat{\mathrm{GL}(1, F)} \times \mathrm{GL}(1, F)$-module $\chi_{V, \psi} \Sigma_{1}^{\prime}$, by applying the second Frobenius we get a nonzero $\widehat{\mathrm{GL}(1, F)} \times \mathrm{GL}(1, F) \times \overparen{\mathrm{Sp}(1)} \times O(3)$-homomorphism

$$
v^{l_{1} / 2} \otimes \chi_{V, \psi} v^{-l_{1} / 2} \otimes \omega_{1,3} \rightarrow v^{l_{1} / 2} \otimes \tau_{1} \otimes \widetilde{R_{\widetilde{P}_{1}}\left(\widetilde{\pi}_{1}\right)}
$$

which implies that $\widetilde{R_{\widetilde{P}_{1}}\left(\widetilde{\pi}_{1}\right)}\left(\chi_{V, \psi} v^{-l_{1} / 2}\right) \neq 0$. Because $l_{1} \neq 0$, this contradicts our assumption $\pi_{1} \hookrightarrow \chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}$; hence $\left.T\right|_{I_{11}}=0$. Therefore, we can consider $T$ as an epimorphism $I_{10} \rightarrow \pi_{1} \otimes v^{l_{1} / 2} \otimes \tau_{1}$. Consequently, $l_{1}=-3$ and there is an epimorphism $\omega_{2,3} \rightarrow \pi \otimes \tau_{1}$. Obviously, $\Theta\left(\pi_{1}, 3\right) \neq 0$.

Repeating the same procedure once again, we obtain $\Theta\left(\pi_{1}, 2\right) \neq 0$. Since the cuspidal support of each irreducible quotient of $\Theta\left(\pi_{1}, 2\right)$ equals [ $\left.\zeta_{1}, \zeta_{2} ; 1\right]$, all of the irreducible quotients of $\Theta\left(\pi_{1}, 2\right)$ are equal to $\zeta_{1} \times \zeta_{2} \rtimes 1$.
Proposition 3.4. Let $\zeta_{1}, \zeta_{2} \in \widehat{F^{\times}}$. Then the unitary principal series representation $\chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}$ is irreducible.

We present two proofs of this proposition, both based on the previous lemma. The first proof is much simpler than the second - it also uses some known results about Whittaker models for the principal series for metaplectic groups, but we have to assume that the residue characteristic of $F$ is odd. The second proof is more technical, but it doesn't depend on the residue characteristic of $F$. We feel that presenting both proofs may be useful.

First proof of Proposition 3.4. We denote the representation $\chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}$ by $\Pi$. Suppose that the residue characteristic of $F$ is not 2 . Howe's duality conjecture and lemma then implies that the representation $\Theta\left(\zeta_{1} \times \zeta_{2} \rtimes 1,2\right)$ has a unique irreducible quotient, so, by Lemma 3.3, all the irreducible subrepresentations of $\Pi$ are isomorphic, that is,

$$
\begin{equation*}
\Pi=\pi \oplus \cdots \oplus \pi \tag{3}
\end{equation*}
$$

Now, observe that the representation $\Pi$ has a unique Whittaker model [Banks 1998; Szpruch 2007]. In more words, for a nondegenerate character $\theta$ of the unipotent radical $U$ of Borel subgroup of $\mathrm{Sp}(n)$ (observe that $\widetilde{\mathrm{Sp}(n)}$ splits over $U$, and the mapping $n \mapsto(n, 1)$ is the splitting) and a genuine character $\chi_{V, \psi} \zeta_{1} \otimes \cdots \otimes \chi_{V, \psi} \zeta_{n}$ of $\widetilde{T}$ (where $\widetilde{T}$ denotes the preimage of maximal diagonal torus in $\operatorname{Sp}(n)$ ), we have

$$
\operatorname{dim}_{\mathbb{C}} \operatorname{Hom}_{\widetilde{\mathrm{Pp}_{n}}}\left(\chi_{V, \psi} \zeta_{1} \times \cdots \times \chi_{V, \psi} \zeta_{n} \rtimes \omega_{0}, \operatorname{Ind}_{U}^{\widetilde{\mathrm{Sp}(n)}}(\theta)\right)=1
$$

This forces that the number of copies of $\pi$ in (3) to be one, and this finishes the first proof.

Second proof of Proposition 3.4. We have already seen that there is an epimorphism

$$
\begin{aligned}
R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right) \rightarrow & \chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \rtimes \omega_{0} \otimes \zeta_{1} \times \zeta_{2} \rtimes 1, \text { so } \\
& \Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \zeta_{1} \times \zeta_{2} \rtimes 1, R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right)\right) \neq 0 .
\end{aligned}
$$

$R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right)$ has the filtration in which

- $J_{10}=\chi_{V, \psi} \nu^{1 / 2} \otimes \omega_{1,2}$ is the quotient and
- $J_{11}=\operatorname{Ind} \frac{\widetilde{M_{1}} \times O(2)}{\operatorname{GL}(1) \times P_{1} \times \widetilde{\operatorname{Sp}(1)}}\left(\chi_{V, \psi} \Sigma_{1}^{\prime} \otimes \omega_{1,1}\right)$ is the subrepresentation.

Lemma 3.5. There is an isomorphism
$\operatorname{Hom}_{\widetilde{M}_{1}}\left(R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right), \chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}\right) \cong \operatorname{Hom}_{\widetilde{M}_{1}}\left(J_{11}, \chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}\right)$ of vector spaces that is given by restriction (that is, $\left.T \mapsto T\right|_{J_{11}}$ ).
Proof of Lemma 3.5. The map obtained by the restriction is obviously a homomorphism, while the injectivity follows directly. Surjectivity is proved as follows:

We consider the filtration $0 \subseteq W_{2} \subseteq W_{1} \subseteq R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right)$, where $W_{1}$ is the representation $J_{11}$, and $W_{1} / W_{2} \cong \chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, J_{11}\right)$. Observe that

$$
\left(R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right) / W_{2}\right) /\left(W_{1} / W_{2}\right) \cong R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right) / W_{1} \cong J_{10}
$$

Using standard argument, it can be proved that the representation $R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right) / W_{2}$ is $\widetilde{\mathrm{GL}(1)}$-finite. Then, using the decomposition along the generalized central characters, which in this case coincide with the central characters because $W_{1} / W_{2}$ and $J_{10}$ have different central characters, we obtain

$$
R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right) / W_{2} \cong W_{1} / W_{2} \oplus J_{10}
$$

Now an element of $\operatorname{Hom}_{\widetilde{M}_{1}}\left(J_{11}, \chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}\right)$ is trivial on $W_{2}$, so it can be extended to $R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right)$ in an obvious way and surjectivity is proved.

Using a standard relation between taking a smooth part of the isotypic component of a representation and the homomorphism functor [Hanzer and Muić 2009, page 10], it follows from Lemma 3.5 that

$$
\Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right)\right) \cong \Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, J_{11}\right)
$$

if we can prove that $\Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, J_{11}\right)$ is admissible.
Lemma 3.6. We have $\Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, J_{11}\right)=\zeta_{1} \times \zeta_{2} \rtimes 1$.
Proof of Lemma 3.6. Since $\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \zeta_{1} \times \zeta_{2} \rtimes 1$ is a quotient of $J_{11}$, there is an epimorphism $\Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, J_{11}\right) \rightarrow \zeta_{1} \times \zeta_{2} \rtimes 1$.

Applying [Hanzer and Muić 2009, Lemma 3.2], we have

$$
\begin{aligned}
& \operatorname{Hom}_{\widetilde{M}_{1} \times O(2)}\left(J_{11}, \chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, J_{11}\right)\right) \\
& \cong \operatorname{Hom}_{\widetilde{M}_{1} \times M(1)}\left(\chi_{V, \psi} \Sigma_{1}^{\prime} \otimes \omega_{1,1},\right. \\
& \left.\quad \chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes R_{\bar{P}_{1}}\left(\Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, J_{11}\right)\right)\right) .
\end{aligned}
$$

For every intertwining map $T$ from the first space, let $T_{0}$ be the corresponding intertwining map from the second space. Let $\varphi$ be a natural homomorphism belonging to the first space.

Since $\chi_{V, \psi} \zeta_{1} \otimes \zeta_{1}^{-1}$ (respectively, $\chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \zeta_{2} \rtimes 1$ ) are the corresponding isotypic components in the $\widehat{\mathrm{GL}(1, F)} \times \mathrm{GL}(1, F)$-module $\chi_{V, \psi} \Sigma_{1}^{\prime}$ (respectively, in
the $\widetilde{\operatorname{Sp}(1)} \times O(3)$-module $\left.\omega_{1,1}\right)$, irreducibility of these isotypic components implies that the image of $\varphi_{0}$ is isomorphic to $\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \zeta_{1}^{-1} \otimes \zeta_{2} \rtimes 1$. Now, we write $\varphi_{0}=\varphi^{\prime \prime} \circ \varphi^{\prime}$, where $\varphi^{\prime}$ is a canonical epimorphism

$$
\chi_{V, \psi} \Sigma_{1}^{\prime} \otimes \omega_{1,1} \rightarrow \chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \zeta_{1}^{-1} \otimes \zeta_{2} \rtimes 1
$$

and $\varphi^{\prime \prime}$ is an inclusion of the representation

$$
\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \zeta_{1}^{-1} \otimes \zeta_{2} \rtimes 1
$$

in $\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes R_{\bar{P}_{1}}\left(\Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, J_{11}\right)\right)$. Observe that $\operatorname{Ind}\left(\varphi^{\prime}\right)$ is a homomorphism

$$
\operatorname{Ind} \underset{\widetilde{\mathcal{M L}_{1} \times O}(1) \times P_{1} \times \widetilde{\operatorname{Sp}(1)}}{ }\left(\chi_{V, \psi} \Sigma_{1}^{\prime} \otimes \omega_{1,1}\right) \rightarrow \chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \zeta_{1}^{-1} \times \zeta_{2} \rtimes 1
$$

Let $\varphi_{1}$ be an operator belonging to
$\operatorname{Hom}\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \zeta_{1} \times \zeta_{2} \rtimes 1\right.$,

$$
\left.\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, J_{11}\right)\right)
$$

such that $\left(\varphi_{1}\right)_{0}=\varphi^{\prime \prime}$.
Lemma 3.7. Under the assumptions above, $\left(\varphi_{1} \circ \operatorname{Ind}\left(\varphi^{\prime}\right)\right)_{0}=\varphi_{0}$.
Proof of Lemma 3.7. We prove it much more generally. Let $(\pi, V)$ be a smooth representation of some Levi subgroup $M^{\prime}$ in the parabolic $P^{\prime}$ and the opposite parabolic $\overline{P^{\prime}}$ of the group $G^{\prime}$ (which is one of the groups we are considering, that is, metaplectic or odd orthogonal) and let $(\Pi, W)$ be a smooth representation of $G^{\prime}$. Then the second Frobenius isomorphism asserts

$$
\operatorname{Hom}_{G^{\prime}}\left(\operatorname{Ind}_{M^{\prime}}^{G^{\prime}}(\pi), \Pi\right) \cong \operatorname{Hom}_{M^{\prime}}\left(\pi, R_{\overline{P^{\prime}}}(\Pi)\right)
$$

Let $\psi \hookrightarrow R_{\overline{P^{\prime}}}\left(\operatorname{Ind}_{M^{\prime}}^{G^{\prime}}(\pi)\right)$ be an embedding corresponding to the open cell $P^{\prime} \overline{P^{\prime}}$ in $G^{\prime}$ given in the following way:

For an open compact subgroup $K$ of $G^{\prime}$ that has Iwahori decomposition with respect to both $P^{\prime}$ and $\overline{P^{\prime}}$, and for $v \in V^{K \cap M^{\prime}}$, we define

$$
f_{v, K}(g)=\frac{1}{\operatorname{meas}_{\overline{N^{\prime}}}\left(K \cap \overline{N^{\prime}}\right)} \begin{cases}0 & \text { if } g \notin P^{\prime} K \\ \delta_{P}^{1 / 2}(m) \pi(m) v & \text { if } g=m n k \\ & \text { for } m \in M^{\prime}, n \in N^{\prime}, k \in K\end{cases}
$$

Then $\psi: v \mapsto f_{v, K}+\operatorname{Ind}_{M^{\prime}}^{G^{\prime}}(\pi)\left(\overline{N^{\prime}}\right)$ is independent on the choice of $K$.
For $\varphi \in \operatorname{Hom}_{G^{\prime}}\left(\operatorname{Ind}_{M^{\prime}}^{G^{\prime}}(\pi), \Pi\right)$, we take $\varphi_{0}$ to be the corresponding element of $\operatorname{Hom}_{M^{\prime}}\left(\pi, R_{\bar{P}}(\Pi)\right)$. It follows that $\varphi_{0}(v)=\varphi\left(f_{v, K}\right)+\Pi\left(\overline{N^{\prime}}\right)$. Write $\varphi_{0}=\varphi^{\prime \prime} \circ \varphi^{\prime}$, where $\varphi^{\prime}$ denotes the canonical epimorphism $\pi \rightarrow \pi / \operatorname{Ker} \varphi_{0}$ and $\varphi^{\prime \prime}$ denotes the
embedding $\pi / \operatorname{Ker} \varphi_{0} \hookrightarrow R_{\overline{P^{\prime}}}$. So, we are able to construct the mapping $\operatorname{Ind}\left(\varphi^{\prime}\right)$ : $\operatorname{Ind}_{M^{\prime}}^{G^{\prime}}(\pi) \rightarrow \operatorname{Ind}_{M^{\prime}}^{G^{\prime}}\left(\pi / \operatorname{Ker} \varphi_{0}\right)$. Since

$$
\operatorname{Hom}_{M^{\prime}}\left(\pi / \operatorname{Ker} \varphi_{0}, R_{\bar{P}^{\prime}}(\Pi)\right) \cong \operatorname{Hom}_{G^{\prime}}\left(\operatorname{Ind}_{M^{\prime}}^{G^{\prime}}\left(\pi / \operatorname{Ker} \varphi_{0}\right), \Pi\right)
$$

analogously as above, we see there is an element $\varphi_{1} \in \operatorname{Hom}_{G^{\prime}}\left(\operatorname{Ind}_{M^{\prime}}^{G^{\prime}}\left(\pi / \operatorname{Ker} \varphi_{0}\right), \Pi\right)$ such that $\left(\varphi_{1}\right)_{0}=\varphi^{\prime \prime}$.

To prove $\left(\varphi_{1} \circ \operatorname{Ind}\left(\varphi^{\prime}\right)\right)_{0}=\varphi_{0}$, it is enough to prove $\left(\varphi_{1} \circ \operatorname{Ind}\left(\varphi^{\prime}\right)\right)_{0}=\left(\varphi_{1}\right)_{0} \circ \varphi^{\prime}$.
Let $v \in V$. Clearly, $\varphi^{\prime}(v)=v+\operatorname{Ker} \varphi_{0}$. Further,

$$
\begin{aligned}
\left(\varphi_{1}\right)_{0}\left(\varphi^{\prime}(v)\right) & =\varphi_{1}\left(f_{v+\operatorname{Ker} \varphi_{0}, K}\right)+\Pi\left(\overline{N^{\prime}}\right) \\
\left(\varphi_{1} \circ \operatorname{Ind}\left(\varphi^{\prime}\right)\right)_{0}(v) & =\varphi_{1}\left(\operatorname{Ind}\left(\varphi^{\prime}\right) f_{v, K}\right)+\Pi\left(\overline{N^{\prime}}\right)
\end{aligned}
$$

It follows easily that $f_{v+\operatorname{Ker} \varphi_{0}, K}=f_{v, K}+\operatorname{Ker} \varphi_{0} \operatorname{and} \operatorname{Ind}\left(\varphi^{\prime}\right) f_{v, K}=f_{v, K}+\operatorname{Ker} \varphi_{0}$, and the lemma follows.

We can complete the proof of Lemma 3.6. Lemma 3.7 gives $\varphi_{1} \circ \operatorname{Ind}\left(\varphi^{\prime}\right)=\varphi$, so the image of $\varphi$ is a quotient of $\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes \zeta_{1}^{-1} \times \zeta_{2} \rtimes 1$. This implies that $\Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, J_{11}\right)$ is a quotient of $\zeta_{1}^{-1} \times \zeta_{2} \rtimes 1$. Since $\zeta_{1}^{-1} \times \zeta_{2} \rtimes 1 \simeq \zeta_{1} \times \zeta_{2} \rtimes 1$ is an irreducible representation,

$$
\Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, J_{11}\right)=\zeta_{1} \times \zeta_{2} \rtimes 1
$$

Lemma 3.8. There is an epimorphism $\Theta\left(\zeta_{1} \times \zeta_{2} \rtimes 1,2\right) \rightarrow \chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}$.
Proof of Lemma 3.8. We have an isomorphism

$$
\operatorname{Hom}_{O(2)}\left(\omega_{2,2}, \zeta_{1}^{-1} \times \zeta_{2} \rtimes 1\right) \cong \operatorname{Hom}\left(R_{P_{1}}\left(\omega_{2,2}\right), \zeta_{1}^{-1} \otimes \zeta_{2} \rtimes 1\right)
$$

of vector spaces, which also an isomorphism of $\widetilde{\operatorname{Sp}(2)}$ modules. By taking the smooth parts, we obtain

$$
\operatorname{Hom}_{\widetilde{\operatorname{Spp}(2)} \times O(2)}\left(\omega_{2,2}, \zeta_{1}^{-1} \times \zeta_{2} \rtimes 1\right)_{\infty} \cong \operatorname{Hom}\left(R_{P_{1}}\left(\omega_{2,2}\right), \zeta_{1}^{-1} \otimes \zeta_{2} \rtimes 1\right)_{\infty}
$$

so that $\Theta\left(\zeta_{1}^{-1} \times \zeta_{2} \rtimes 1,2\right)^{\sim} \cong \Theta\left(\zeta_{1}^{-1} \otimes \zeta_{2} \rtimes 1, R_{P_{1}}\left(\omega_{2,2}\right)\right)^{\sim}$.
In the same way as before, we get

$$
\Theta\left(\zeta_{1}^{-1} \otimes \zeta_{2} \rtimes 1, R_{P_{1}}\left(\omega_{2,2}\right)\right)^{\sim} \cong \Theta\left(\zeta_{1}^{-1} \otimes \zeta_{2} \rtimes 1, I_{11}\right)^{\sim}
$$

Now, the epimorphism $I_{11} \rightarrow \zeta_{1}^{-1} \otimes \zeta_{2} \rtimes 1 \otimes \chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}$ gives an epimorphism $\left.\Theta\left(\zeta_{1}^{-1} \otimes \zeta_{2} \rtimes 1, I_{11}\right)\right) \rightarrow \chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}$. Since the representations $\zeta_{1}^{-1} \times \zeta_{2} \rtimes 1$ and $\zeta_{1} \times \zeta_{2} \rtimes 1$ are isomorphic, we obtain the epimorphism $\Theta\left(\zeta_{1} \times \zeta_{2} \rtimes 1,2\right) \rightarrow \chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}$, which proves the lemma.

Now we finish the second proof of Proposition 3.4. Suppose that the representation $\chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}$ reduces. Suppose also that it is the representation of length 2 and write $\chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}=\pi_{1} \oplus \pi_{2}$. Obviously, $R_{\widetilde{P}_{1}}\left(\chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}\right)=R_{\widetilde{P}_{1}}\left(\pi_{1}\right) \oplus R_{\widetilde{P}_{1}}\left(\pi_{2}\right)$.

We have, by Lemma 3.8, an epimorphism

$$
\omega_{2,2} \rightarrow \zeta_{1} \times \zeta_{2} \rtimes 1 \otimes \chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}
$$

which leads to the epimorphisms $R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right) \rightarrow \zeta_{1} \times \zeta_{2} \rtimes 1 \otimes\left(R_{\widetilde{P}_{1}}\left(\pi_{1}\right) \oplus R_{\widetilde{P}_{1}}\left(\pi_{2}\right)\right)$ and $R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right) \rightarrow \chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0} \otimes\left(\zeta_{1} \times \zeta_{2} \rtimes 1 \oplus \zeta_{1} \times \zeta_{2} \rtimes 1\right)$.

Finally, we obtain an epimorphism

$$
\Theta\left(\chi_{V, \psi} \zeta_{1} \otimes \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}, R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right)\right) \rightarrow \zeta_{1} \times \zeta_{2} \rtimes 1 \oplus \zeta_{1} \times \zeta_{2} \rtimes 1
$$

which contradicts Lemmas 3.5 and 3.6.
The same proof remains valid if we suppose that $\chi_{V, \psi} \zeta_{1} \times \chi_{V, \psi} \zeta_{2} \rtimes \omega_{0}$ is the representation of the length 4.
3.2. Nonunitary principal series. First we determine the reducibility points of the representations with cuspidal support in the minimal parabolic subgroup $\widetilde{P_{(1,1)}}$.

Let $\chi_{1}, \chi_{2} \in \widehat{F^{\times}}$and $s_{i} \geq 0$ for $i=1,2$, such that $s_{i}>0$ for at least one $i$. Define $\Pi=\chi_{V, \psi} \nu^{s_{1}} \chi_{1} \times \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \rtimes \omega_{0}$. We have

$$
\begin{aligned}
\mu^{*}(\Pi)= & \chi_{V, \psi} \nu^{s_{1}} \chi_{1} \otimes \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \rtimes \omega_{0}+\chi_{V, \psi} v^{-s_{1}} \chi_{1}^{-1} \otimes \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \rtimes \omega_{0} \\
& +\chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \chi_{V, \psi} \nu^{s_{1}} \chi_{1} \rtimes \omega_{0}+\chi_{V, \psi} \nu^{-s_{2}} \chi_{2}^{-1} \otimes \chi_{V, \psi} \nu^{s_{1}} \chi_{1} \rtimes \omega_{0} \\
& +\chi_{V, \psi} \nu^{-s_{1}} \chi_{1}^{-1} \times \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \omega_{0}+\chi_{V, \psi} \nu^{s_{1}} \chi_{1} \times \chi_{V, \psi} \nu^{-s_{2}} \chi_{2}^{-1} \otimes \omega_{0} \\
& +\chi_{V, \psi} \nu^{s_{1}} \chi_{1} \times \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \omega_{0}+\chi_{V, \psi} \nu^{-s_{1}} \chi_{1}^{-1} \times \chi_{V, \psi} \nu^{-s_{2}} \chi_{2}^{-1} \otimes \omega_{0} \\
& +1 \otimes \chi_{V, \psi} \nu^{s_{1}} \chi_{1} \times \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \rtimes \omega_{0} .
\end{aligned}
$$

We prove that irreducibility of all the representations above implies irreducibility of the representation $\Pi$. We keep this assumption throughout this subsection.

First, suppose that $v^{s_{1}} \chi_{1} \neq v^{-s_{1}} \chi_{1}^{-1}, v^{s_{2}} \chi_{2} \neq v^{-s_{2}} \chi_{2}^{-1}$ and $v^{s_{1}} \chi_{1} \neq v^{ \pm s_{2}} \chi_{2}^{ \pm 1}$ (that is, Jacquet modules of $\Pi$ are multiplicity one).

Let $\tau$ be an irreducible subquotient of $\Pi$ such that

$$
\chi_{V, \psi} \nu^{-s_{1}} \chi_{1}^{-1} \times \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \omega_{0} \leq R_{\widetilde{P}_{2}}(\tau)
$$

From transitivity of Jacquet modules we get

$$
\chi_{V, \psi} \nu^{-s_{1}} \chi_{1}^{-1} \otimes \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \omega_{0}+\chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \chi_{V, \psi} \nu^{-s_{1}} \chi_{1}^{-1} \otimes \omega_{0} \leq R_{\widetilde{P_{(1,1)}}}(\tau)
$$

This implies

$$
\chi_{V, \psi} v^{-s_{1}} \chi_{1}^{-1} \otimes \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \rtimes \omega_{0}+\chi_{V, \psi} v^{s_{2}} \chi_{2} \otimes \chi_{V, \psi} v^{s_{1}} \chi_{1} \rtimes \omega_{0} \leq R_{\widetilde{P}_{1}}(\tau) .
$$

We get directly that

$$
\begin{aligned}
R_{\widetilde{P}_{2}}(\tau)= & \chi_{V, \psi} \nu^{-s_{1}} \chi_{1}^{-1} \times \chi_{V, \psi} v^{s_{2}} \chi_{2} \otimes \omega_{0}+\chi_{V, \psi} \nu^{s_{1}} \chi_{1} \times \chi_{V, \psi} \nu^{-s_{2}} \chi_{2}^{-1} \otimes \omega_{0} \\
& +\chi_{V, \psi} \nu^{s_{1}} \chi_{1} \times \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \omega_{0}+\chi_{V, \psi} \nu^{-s_{1}} \chi_{1}^{-1} \times \chi_{V, \psi} v^{-s_{2}} \chi_{2}^{-1} \otimes \omega_{0}
\end{aligned}
$$

so $\tau=\Pi$ and $\Pi$ is irreducible.
Now we assume that there is some $i$ such that $\nu^{s_{i}} \chi_{i} \neq v^{-s_{i}} \chi_{i}^{-1}$. Without loss of generality, let $i=1$. So, $s_{1}=0$ and $\chi_{1}=\chi_{1}^{-1}$, that is, $\chi_{1}^{2}=1_{F^{\times}}$. We prove that in this case $\Pi$ is also irreducible. Again, we start by writing corresponding Jacquet modules:

$$
\begin{aligned}
R_{\widetilde{P}_{1}}(\Pi)=2 \chi_{V, \psi} \chi_{1} \otimes \chi_{V, \psi} v^{s_{2}} \chi_{2} \rtimes \omega_{0} & +\chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \chi_{V, \psi} \chi_{1} \rtimes \omega_{0} \\
& +\chi_{V, \psi} v^{-s_{2}} \chi_{2}^{-1} \otimes \chi_{V, \psi} \chi_{1} \rtimes \omega_{0} \\
R_{\widetilde{P}_{2}}(\Pi)=2 \chi_{V, \psi} \chi_{1} \times \chi_{V, \psi} v^{s_{2}} \chi_{2} \otimes \omega_{0} & +\chi_{V, \psi} \chi_{1} \times \chi_{V, \psi} v^{-s_{2}} \chi_{2}^{-1} \otimes \omega_{0}
\end{aligned}
$$

Let $\tau$ be an irreducible subquotient of $\Pi$ such that

$$
R_{\widetilde{P}_{1}}(\tau) \geq \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \chi_{V, \psi} \chi_{1} \rtimes \omega_{0} .
$$

Of course, $R_{\widetilde{P_{(1,1)}}}(\tau) \geq 2 \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \chi_{V, \psi} \chi_{1} \otimes \omega_{0}$, so

$$
R_{\widetilde{P}_{2}}(\tau) \geq 2 \chi_{V, \psi} \chi_{1} \times \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \omega_{0}
$$

Continuing in the same way, we get

$$
\begin{aligned}
R_{\widetilde{P_{(1,1)}}}(\tau) & \geq 2 \chi_{V, \psi} \chi_{1} \otimes \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \omega_{0}+2 \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \chi_{V, \psi} \chi_{1} \otimes \omega_{0}, \\
R_{\widetilde{P}_{1}}(\tau) & \geq 2 \chi_{V, \psi} \chi_{1} \otimes \chi_{V, \psi} v^{s_{2}} \chi_{2} \rtimes \omega_{0}+\chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \chi_{V, \psi} \chi_{1} \rtimes \omega_{0} .
\end{aligned}
$$

Finally,

$$
\begin{aligned}
& R_{\widetilde{P_{(1,1)}}}(\tau) \geq 2 \chi_{V, \psi} \chi_{1} \otimes \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \omega_{0}+2 \chi_{V, \psi} \chi_{1} \otimes \chi_{V, \psi} v^{-s_{2}} \chi_{2}^{-1} \otimes \omega_{0} \\
&+2 \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \chi_{V, \psi} \chi_{1} \otimes \omega_{0} \\
& R_{\widetilde{P}_{2}}(\tau) \geq 2 \chi_{V, \psi} \chi_{1} \times \chi_{V, \psi} \nu^{s_{2}} \chi_{2} \otimes \omega_{0}+\chi_{V, \psi} \chi_{1} \times \chi_{V, \psi} v^{-s_{2}} \chi_{2}^{-1} \otimes \omega_{0}=R_{\widetilde{P}_{2}}(\Pi) .
\end{aligned}
$$

So, $\Pi=\tau$ and $\Pi$ is irreducible.
If $v^{s_{1}} \chi_{1}=v^{s_{2}} \chi_{2}$ or $v^{s_{1}} \chi_{1}=v^{-s_{2}} \chi_{2}^{-1}$, then the irreducibility of $\Pi$ follows in the same way as above. Observe that equalities $v^{s_{1}} \chi_{1}=v^{-s_{1}} \chi_{1}^{-1}$ and $v^{s_{2}} \chi_{2}=v^{-s_{2}} \chi_{2}^{-1}$ lead to unitary principal series.

In this way we have proved irreducibility of the principal series, with these exceptions:

- Some of the representations $\chi_{V, \psi} \nu^{s_{1}} \chi_{1} \rtimes \omega_{0}$ or $\chi_{V, \psi} \nu^{s_{2}} \chi_{2} \rtimes \omega_{0}$ reduce (the so-called $\widetilde{\mathrm{Sp}(1)}$ reducibility).
- Some of the representations

$$
\begin{array}{ll}
\chi_{V, \psi} \nu^{-s_{1}} \chi_{1}^{-1} \times \chi_{V, \psi} \nu^{s_{2}} \chi_{2}, & \chi_{V, \psi} \nu^{s_{1}} \chi_{1} \times \chi_{V, \psi} \nu^{-s_{2}} \chi_{2}^{-1}, \\
\chi_{V, \psi} \nu^{s_{1}} \chi_{1} \times \chi_{V, \psi} \nu^{s_{2}} \chi_{2}, & \chi_{V, \psi} \nu^{-s_{1}} \chi_{1}^{-1} \times \chi_{V, \psi} \nu^{-s_{2}} \chi_{2}^{-1}
\end{array}
$$

reduce (the so-called $\widetilde{\mathrm{GL}(2)}$ reducibility).
3.2.1. $\widetilde{\mathrm{Sp}(1)}$ reducibility. Let $\chi, \zeta \in \widehat{F^{\varnothing}}, \zeta^{2}=1_{F^{\times}}$, and $s \geq 0$. It is well known that, in $R$,

$$
\chi_{V, \psi} \nu^{s} \chi \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}=\chi_{V, \psi} \nu^{s} \rtimes s p_{\zeta, 1}+\chi_{V, \psi} \nu^{s} \rtimes \omega_{\psi_{a}, 1}^{+}
$$

Let $\Pi$ denote $\chi_{V, \psi} \nu^{s} \rtimes s p_{\zeta, 1}$.
Calculating Jacquet modules, we find
$R_{\widetilde{P}_{1}}(\Pi)=\chi_{V, \psi} \nu^{-s} \chi^{-1} \otimes s p_{\zeta, 1}+\chi_{V, \psi} \nu^{s} \chi \otimes s p_{\zeta, 1}+\chi_{V, \psi} \nu^{1 / 2} \zeta \otimes \chi_{V, \psi} \nu^{s} \chi \rtimes \omega_{0}$, $R_{\widetilde{P}_{2}}(\Pi)=\chi_{V, \psi} \nu^{-s} \chi^{-1} \times \chi_{V, \psi} \nu^{1 / 2} \zeta \otimes \omega_{0}+\chi_{V, \psi} \nu^{s} \chi \times \chi_{V, \psi} \nu^{1 / 2} \zeta \otimes \omega_{0}$.

If the representation $\chi_{V, \psi} \nu^{s} \chi \rtimes \omega_{0}$ is irreducible (that is, when $\nu^{s} \chi \neq \nu^{ \pm 1 / 2} \zeta_{2}$, where $\zeta_{2}^{2}=1_{F^{\times}}$), we proceed in the following way:

Let $\rho$ be an irreducible subquotient of $\Pi$ such that

$$
\chi_{V, \psi} \nu^{1 / 2} \zeta \otimes \chi_{V, \psi} \nu^{s} \chi \rtimes \omega_{0} \leq s_{1}(\rho)
$$

We directly get that

$$
\chi_{V, \psi} \nu^{1 / 2} \zeta \otimes \chi_{V, \psi} \nu^{s} \chi \otimes \omega_{0}+\chi_{V, \psi} \nu^{1 / 2} \zeta \otimes \chi_{V, \psi} \nu^{-s} \chi^{-1} \otimes \omega_{0} \leq R_{P_{(1,1)}}(\rho)
$$

If both $\chi_{V, \psi} \nu^{-s} \chi^{-1} \times \chi_{V, \psi} \nu^{1 / 2} \zeta$ and $\chi_{V, \psi} \nu^{s} \chi \times \chi_{V, \psi} \nu^{1 / 2} \zeta$ are irreducible, $\Pi$ is also irreducible.

For the reducibility of the $\widetilde{\mathrm{Sp}(1)}$ part we still have to determine the composition factors of the representations
(i) $\chi_{V, \psi} \nu^{1 / 2} \zeta_{1} \times \chi_{V, \psi} \nu^{1 / 2} \zeta_{2} \rtimes \omega_{0}$,
(ii) $\chi_{V, \psi} \nu^{1 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}$, and
(iii) $\chi_{V, \psi} \nu^{3 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}$, where $\zeta^{2}=\zeta_{1}^{2}=\zeta_{2}^{2}=1_{F^{*}}$.

Thus, we have proved the following result:
Proposition 3.9. Let $\chi \in \widehat{F^{\times}}$, a nonnegative $s \in \mathbb{R}$, and $\zeta \in \widehat{F^{\times}}$with $\zeta^{2}=1_{F^{\times}}$. The representations $\chi_{V, \psi} \nu^{s} \chi \rtimes s p_{\zeta, 1}$ and $\chi_{V, \psi} \nu^{s} \chi \rtimes \omega_{\psi_{a}, 1}^{+}$are irreducible unless $(s, \chi)=(3 / 2, \zeta)$ or $\left(1 / 2, \zeta_{1}\right)$, where $\zeta_{1}^{2}=1_{F^{\times}}$. In $R$, we have

$$
\chi_{V, \psi} \nu^{s} \chi \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}=\chi_{V, \psi} \nu^{s} \chi \rtimes s p_{\zeta, 1}+\chi_{V, \psi} \nu^{s} \chi \rtimes \omega_{\psi_{a}, 1}^{+}
$$

Also, if $(s, \chi) \neq(3 / 2, \zeta)$ and $(s, \chi) \neq\left(1 / 2, \zeta_{1}\right)$, then

$$
\begin{aligned}
\chi_{V, \psi} \nu^{s} \chi \rtimes s p_{\zeta, 1} & = \begin{cases}\left\langle\chi_{V, \psi} \nu^{1 / 2} \zeta ; \chi_{V, \psi} \chi \rtimes \omega_{0}\right\rangle & \text { if } s=0, \\
\left\langle\chi_{V, \psi} \nu^{1 / 2} \zeta, \chi_{V, \psi} \nu^{s} \chi ; \omega_{0}\right\rangle & \text { if } 0<s \leq 1 / 2, \\
\left\langle\chi_{V, \psi} \nu^{s} \chi, \chi_{V, \psi} \nu^{1 / 2} \zeta ; \omega_{0}\right\rangle & \text { if } s>1 / 2,\end{cases} \\
\chi_{V, \psi} \nu^{s} \chi \rtimes \omega_{\psi_{a}, 1}^{+} & = \begin{cases}\chi_{V, \psi} \chi \rtimes \omega_{\psi_{a}, 1}^{+} & \text {if } s=0, \\
\left\langle\chi_{V, \psi} \nu^{s} \chi ; \omega_{\psi_{a}, 1}^{+}\right\rangle & \text {if } s>0 .\end{cases}
\end{aligned}
$$

3.2.2. $\widehat{\mathrm{GL}(2)}$ reducibility. Let $\chi \in \widehat{F^{\times}}$and $s \in \mathbb{R}$ be nonnegative. In $R$, we have $\chi_{V, \psi} \nu^{s+1 / 2} \chi \times \chi_{V, \psi} \nu^{s-1 / 2} \chi \rtimes \omega_{0}=\chi_{V, \psi} \nu^{s} \chi \mathrm{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}+\chi_{V, \psi} \nu^{s} \chi 1_{\mathrm{GL}(2)} \rtimes \omega_{0}$.

Let $\Pi$ denote the representation $\chi_{V, \psi} \nu^{s} \chi \operatorname{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}$. Calculation of $\mu^{*}(\Pi)$ gives

$$
\begin{aligned}
R_{\widetilde{P}_{1}}(\Pi)=\chi_{V, \psi} \nu^{s+1 / 2} \chi \otimes \chi_{V, \psi} \nu^{s-1 / 2} \chi & \rtimes \omega_{0} \\
& +\chi_{V, \psi} \nu^{1 / 2-s} \chi^{-1} \otimes \chi_{V, \psi} \nu^{s+1 / 2} \chi \rtimes \omega_{0},
\end{aligned}
$$

$$
\begin{aligned}
& R_{\widetilde{P}_{2}}(\Pi)=\chi_{V, \psi} \nu^{s} \chi \operatorname{St}_{\mathrm{GL}(2)} \otimes \omega_{0}+\chi_{V, \psi} v^{-s} \chi^{-1} \mathrm{St}_{\mathrm{GL}(2)} \otimes \omega_{0} \\
&+\chi_{V, \psi} \nu^{s+1 / 2} \chi \times \chi_{V, \psi} v^{1 / 2-s} \chi^{-1} \otimes \omega_{0}
\end{aligned}
$$

Looking at Jacquet modules with respect to different parabolic subgroups we can conclude, in the same way as in the $\widetilde{\mathrm{Sp}(1)}$ reducibility case, that if

$$
\chi_{V, \psi} \nu^{s-1 / 2} \chi \rtimes \omega_{0}, \quad \chi_{V, \psi} \nu^{s+1 / 2} \chi \rtimes \omega_{0}, \quad \chi_{V, \psi} \nu^{s+1 / 2} \chi \times \chi_{V, \psi} \nu^{1 / 2-s} \chi^{-1}
$$

are irreducible representations, then the representation $\Pi$ is also irreducible.
Observe that the representation $\chi_{V, \psi} \nu^{s+1 / 2} \chi \rtimes \omega_{0}$ reduces for $(\chi, s)=(\zeta, 0)$, while $\chi_{V, \psi} \nu^{s-1 / 2} \chi \rtimes \omega_{0}$ reduces for $(\chi, s)=(\zeta, 1)$, where $\zeta^{2}=1_{F^{\times}}$.

The representation $\chi_{V, \psi} \nu^{s+1 / 2} \chi \times \chi_{V, \psi} \nu^{1 / 2-s} \chi^{-1}$ reduces for $(\chi, s)=(\zeta, 1 / 2)$, where $\zeta^{2}=1_{F^{\times}}$. These observations imply this:
Proposition 3.10. Let $\chi \in \widehat{F^{\varnothing}}$ and $s \in \mathbb{R}$ be nonnegative. The representations $\chi_{V, \psi} \nu^{s} \chi \operatorname{St}_{\mathrm{GL}(2)} \rtimes 1$ and $\chi_{V, \psi} \nu^{s} \chi 1_{\mathrm{GL}(2)} \rtimes 1$ are irreducible except in the cases that $(s, \chi)=(1 / 2, \zeta),(s, \chi)=(1, \zeta)$ or $(s, \chi)=(0, \zeta)$, where $\zeta^{2}=1_{F^{\times}}$. In $R$, we have
$\chi_{V, \psi} \nu^{s+1 / 2} \chi \times \chi_{V, \psi} \nu^{s-1 / 2} \chi \rtimes \omega_{0}=\chi_{V, \psi} \nu^{s} \chi \operatorname{St}_{G L(2)} \rtimes \omega_{0}+\chi_{V, \psi} \nu^{s} \chi 1_{\mathrm{GL}(2)} \rtimes \omega_{0}$. Also, if $\chi_{V, \psi} \nu^{s+1 / 2} \chi \times \chi_{V, \psi} \nu^{s-1 / 2} \chi \rtimes \omega_{0}$ is a representation of length 2 , then $\chi_{V, \psi} \nu^{s} \chi 1_{\mathrm{GL}(2)} \rtimes \omega_{0}=\left\langle\chi_{V, \psi} \nu^{s} \chi 1_{\mathrm{GL}(2)} ; \omega_{0}\right\rangle$ and

$$
\chi_{V, \psi} v^{s} \chi \operatorname{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}= \begin{cases}\left\langle\chi_{V, \psi} \nu^{s+1 / 2} \chi, \chi_{V, \psi} v^{1 / 2-s} \chi ; \omega_{0}\right\rangle & \text { if } s<1 / 2 \\ \left\langle\chi_{V, \psi} v^{s+1 / 2} \chi ; \chi_{V, \psi} v^{s-1 / 2} \chi \rtimes \omega_{0}\right\rangle & \text { if } s=1 / 2 \\ \left\langle\chi_{V, \psi} v^{s+1 / 2} \chi, \chi_{V, \psi} v^{s-1 / 2} \chi ; \omega_{0}\right\rangle & \text { if } s>1 / 2 .\end{cases}
$$

For the reducibility of the $\widetilde{\text { GL(2) }}$ part, we still have to determine the composition factors of the representations
(i) $\chi_{V, \psi} \nu^{1 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}$,
(ii) $\chi_{V, \psi} \nu^{3 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}$, and
(iii) $\chi_{V, \psi} \nu \zeta \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$, where $\zeta^{2}=1_{F^{\times}}$.

Altogether, this leaves us four exceptional cases of the representations whose composition series we have to determine:
(a) $\chi_{V, \psi} \nu^{1 / 2} \zeta_{1} \times \chi_{V, \psi} \nu^{1 / 2} \zeta_{2} \rtimes \omega_{0}$,
(b) $\chi_{V, \psi} \nu^{1 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}$,
(c) $\chi_{V, \psi} \nu^{3 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}$,
(d) $\chi_{V, \psi} \nu \zeta \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$, where $\zeta^{2}=\zeta_{1}^{2}=\zeta_{2}^{2}=1_{F^{\times}}$and $\zeta_{1} \neq \zeta_{2}$.

These cases are treated in the following subsection.
3.2.3. Exceptional cases. All the equalities that follow are given in semisimplifications. We obtain desired composition series using case-by-case examination:
(a) Write $\chi_{V, \psi} \nu^{1 / 2} \zeta_{i} \rtimes \omega_{0}=\chi_{V, \psi} s p_{\zeta_{i}, 1}+\omega_{\psi_{a_{i}}, 1}^{+}$for $i=1$, 2. In $R$, we have

$$
\begin{aligned}
\chi_{V, \psi} \nu^{1 / 2} \zeta_{1} \times \chi_{V, \psi} \nu^{1 / 2} \zeta_{2} \rtimes \omega_{0} & =\chi_{V, \psi} \nu^{1 / 2} \zeta_{2} \times \chi_{V, \psi} \nu^{1 / 2} \zeta_{1} \rtimes \omega_{0} \\
& =\chi_{V, \psi} \nu^{1 / 2} \zeta_{1} \rtimes s p_{\zeta_{2}, 1}+\chi_{V, \psi} v^{1 / 2} \zeta_{1} \rtimes \omega_{\psi_{a_{2}, 1}}^{+} \\
& =\chi_{V, \psi} \nu^{1 / 2} \zeta_{2} \rtimes s p_{\zeta_{1}, 1}+\chi_{V, \psi} \nu^{1 / 2} \zeta_{2} \rtimes \omega_{\psi_{a_{1}, 1}}^{+}
\end{aligned}
$$

Using standard calculations, we obtain

$$
\begin{aligned}
& R_{\widetilde{P}_{1}}\left(\chi_{V, \psi} v^{1 / 2} \zeta_{1} \rtimes s p_{\zeta_{2}, 1}\right)=\chi_{V, \psi} v^{-1 / 2} \zeta_{1} \otimes s p_{\zeta_{1}, 1}+\chi_{V, \psi} v^{1 / 2} \zeta_{1} \otimes s p_{\zeta_{2}, 1} \\
&+\chi_{V, \psi} v^{1 / 2} \zeta_{2} \otimes s p_{\zeta_{1}, 1}+\chi_{V, \psi} v^{1 / 2} \zeta_{2} \otimes \omega_{\psi_{a_{1}, 1}^{+}}^{+}
\end{aligned}
$$

and

$$
\begin{aligned}
& R_{\widetilde{P}_{2}}\left(\chi_{V, \psi} v^{1 / 2} \zeta_{1} \rtimes s p_{\zeta_{2}, 1}\right)=\chi_{V, \psi} v^{-1 / 2} \zeta_{1} \times \chi_{V, \psi} v^{1 / 2} \zeta_{2} \otimes \omega_{0} \\
&+\chi_{V, \psi} v^{1 / 2} \zeta_{1} \times \chi_{V, \psi} v^{1 / 2} \zeta_{2} \otimes \omega_{0}
\end{aligned}
$$

The last equality implies that the length of $\chi_{V, \psi} \nu^{1 / 2} \zeta_{1} \rtimes s p_{\zeta_{2}, 1}$ is no more than 2. If $\chi_{V, \psi} \nu^{1 / 2} \zeta_{1} \rtimes s p_{\zeta_{2}, 1}$ were an irreducible representation, then it would have to be equal either to $\chi_{V, \psi} \nu^{1 / 2} \zeta_{2} \rtimes s p_{\zeta_{1}, 1}$ or to $\chi_{V, \psi} \nu^{1 / 2} \zeta_{2} \rtimes \omega_{\psi_{a_{1}, 1}}^{+}$, but Jacquet modules of those two representations show that this is not the case. So, we write $\chi_{V, \psi} \nu^{1 / 2} \zeta_{1} \rtimes s p_{\zeta_{2}, 1}=\rho_{1}+\rho_{2}$, where $\rho_{1}$ and $\rho_{2}$ are irreducible representations
such that

$$
\begin{aligned}
& R_{\widetilde{P}_{2}}\left(\rho_{1}\right)=\chi_{V, \psi} v^{-1 / 2} \zeta_{1} \times \chi_{V, \psi} v^{1 / 2} \zeta_{2} \otimes \omega_{0}, \\
& R_{\widetilde{P}_{2}}\left(\rho_{2}\right)=\chi_{V, \psi} v^{1 / 2} \zeta_{1} \times \chi_{V, \psi} v^{1 / 2} \zeta_{2} \otimes \omega_{0} .
\end{aligned}
$$

Clearly, $\rho_{2}$ is square-integrable (since $\rho_{2}=\left\langle\chi_{V, \psi} \nu^{1 / 2} \zeta_{1}, \chi_{V, \psi} \nu^{1 / 2} \zeta_{2} ; \omega_{0}\right\rangle$ ) and $\rho_{1}=\left\langle\chi_{V, \psi} \nu^{1 / 2} \zeta_{2} ; \omega_{\psi_{a_{1}, 1}}^{+}\right\rangle$.

Reasoning in the same way, we obtain that $\chi_{V, \psi} \nu^{1 / 2} \zeta_{1} \rtimes \omega_{\psi_{a_{2}, 1}}^{+}=\rho_{3}+\rho_{4}$, where $\rho_{3}$ and $\rho_{4}$ are irreducible representations such that

$$
\begin{aligned}
& R_{\widetilde{P}_{2}}\left(\rho_{3}\right)=\chi_{V, \psi} v^{-1 / 2} \zeta_{1} \times \chi_{V, \psi} v^{-1 / 2} \zeta_{2} \otimes \omega_{0} \\
& R_{\widetilde{P}_{2}}\left(\rho_{4}\right)=\chi_{V, \psi} v^{1 / 2} \zeta_{1} \times \chi_{V, \psi} v^{-1 / 2} \zeta_{2} \otimes \omega_{0}
\end{aligned}
$$

So, $\rho_{3}$ is a strongly negative representation, while $\rho_{4}=\left\langle\chi_{V, \psi} \nu^{1 / 2} \zeta_{1} ; \omega_{\psi_{a_{2}, 1}}^{+}\right\rangle$. Using Jacquet modules again, we easily obtain the composition factors of the representations above. Thus, we conclude:

Proposition 3.11. Let $\zeta_{1}, \zeta_{2} \in \widehat{F^{\times}}$such that $\zeta_{i}^{2}=1_{F^{\times}}$for $i=1,2\left(\right.$ with $\left.\zeta_{1} \neq \zeta_{2}\right)$. Than the representations

$$
\begin{array}{ll}
\chi_{V, \psi} v^{1 / 2} \zeta_{2} \rtimes \omega_{\psi_{a_{1,1}}}^{+}, & \chi_{V, \psi} v^{1 / 2} \zeta_{2} \rtimes s p_{\zeta_{1}, 1}, \\
\chi_{V, \psi} v^{1 / 2} \zeta_{1} \rtimes \omega_{\psi_{a_{2}, 1}}^{+}, & \chi_{V, \psi} v^{1 / 2} \zeta_{1} \rtimes s p_{\zeta_{2}, 1}
\end{array}
$$

are reducible and $\chi_{V, \psi} \nu^{1 / 2} \zeta_{1} \times \chi_{V, \psi} \nu^{1 / 2} \zeta_{2} \rtimes \omega_{0}$ is a representation of length 4 . The representations $\chi_{V, \psi} \nu^{1 / 2} \zeta_{1} \rtimes s p_{\zeta_{2}, 1}$ and $\chi_{V, \psi} \nu^{1 / 2} \zeta_{2} \rtimes s p_{\zeta_{1}, 1}$ have exactly one irreducible subquotient in common; that subquotient is square-integrable, and we denote it with $\sigma$ (that is, $\sigma=\left\langle\chi_{V, \psi} \nu^{1 / 2} \zeta_{1}, \chi_{V, \psi} \nu^{1 / 2} \zeta_{2} ; \omega_{0}\right\rangle$ ). Also, the unique irreducible common subquotient of

$$
\chi_{V, \psi} v^{1 / 2} \zeta_{1} \rtimes \omega_{\psi_{a_{2}, 1}}^{+} \quad \text { and } \quad \chi_{V, \psi} v^{1 / 2} \zeta_{2} \rtimes \omega_{\psi_{a_{1}, 1}}^{+}
$$

is a strongly negative representation; we denote it by $\rho_{\text {sneg. }}$. In $R$, we have

$$
\begin{aligned}
\chi_{V, \psi} v^{1 / 2} \zeta_{1} \rtimes s p_{\zeta_{2}, 1} & =\sigma+\left\langle\chi_{V, \psi} v^{1 / 2} \zeta_{2} ; \omega_{\psi_{a_{1}, 1}}^{+}\right\rangle \\
\chi_{V, \psi} v^{1 / 2} \zeta_{1} \rtimes \omega_{\psi_{a_{2}, 1}}^{+} & =\left\langle\chi_{V, \psi} v^{1 / 2} \zeta_{1} ; \omega_{\psi_{a_{2}, 1}}^{+}\right\rangle+\rho_{\text {sneg }}, \\
\chi_{V, \psi} v^{1 / 2} \zeta_{2} \rtimes s p_{\zeta_{1}, 1} & =\sigma+\left\langle\chi_{V, \psi} v^{1 / 2} \zeta_{1} ; \omega_{\psi_{a_{2}, 1}}^{+}\right\rangle \\
\chi_{V, \psi} v^{1 / 2} \zeta_{2} \rtimes \omega_{\psi_{a_{1}, 1}}^{+} & =\left\langle\chi_{V, \psi} v^{1 / 2} \zeta_{2} ; \omega_{\psi_{a_{1}, 1}}^{+}\right\rangle+\rho_{\text {sneg }} .
\end{aligned}
$$

(b) In this case, we have

$$
\begin{aligned}
\chi_{V, \psi} v^{1 / 2} \zeta \times \chi_{V, \psi} \nu^{-1 / 2} \zeta \rtimes \omega_{0} & =\chi_{V, \psi} v^{1 / 2} \zeta \rtimes s p_{\zeta, 1}+\chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{\psi_{a}, 1}^{+} \\
& =\chi_{V, \psi} \zeta \operatorname{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}+\chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}
\end{aligned}
$$

From Jacquet modules, we get

$$
\begin{array}{r}
R_{\widetilde{P}_{1}}\left(\chi_{V, \psi} v^{1 / 2} \zeta \rtimes s p_{\zeta, 1}\right)=2 \chi_{V, \psi} v^{1 / 2} \zeta \otimes s p_{\zeta, 1}+\chi_{V, \psi} v^{-1 / 2} \zeta \otimes s p_{\zeta, 1} \\
\quad+\chi_{V, \psi} v^{1 / 2} \zeta \otimes \omega_{\psi_{a}, 1}^{+} \\
R_{\widetilde{P}_{1}}\left(\chi_{V, \psi} v^{1 / 2} \zeta \rtimes \omega_{\psi_{1}}^{+}\right)=2 \chi_{V, \psi} v^{-1 / 2} \zeta \otimes \omega_{\psi_{a}, 1}^{+} \\
+\chi_{V, \psi} v^{1 / 2} \zeta \otimes \omega_{\psi_{a}, 1}^{+} \\
\\
+\chi_{V, \psi} v^{-1 / 2} \zeta \otimes s p_{\zeta, 1} \\
R_{\widetilde{P}_{1}}\left(\chi_{V, \psi} \zeta \mathrm{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}\right)=2 \chi_{V, \psi} v^{1 / 2} \zeta \otimes s p_{\zeta, 1}+2 \chi_{V, \psi} v^{1 / 2} \zeta \otimes \omega_{\psi_{a}, 1}^{+} \\
R_{\widetilde{P}_{1}}\left(\chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}\right)=2 \chi_{V, \psi} v^{-1 / 2} \zeta \otimes s p_{\zeta, 1}+2 \chi_{V, \psi} v^{-1 / 2} \zeta \otimes \omega_{\psi_{a}, 1}^{+}
\end{array}
$$

From preceding Jacquet modules we conclude, as in [Tadić 1994, Chapter 3], that $\chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{\psi_{a}, 1}^{+}$and $\chi_{V, \psi} \zeta \mathrm{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}$ have an irreducible subquotient in common, which is different from both $\chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{\psi_{a}, 1}^{+}$and $\chi_{V, \psi} \zeta \operatorname{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}$. For simplicity of the notation, we let $\rho_{1}$ stand for this subquotient. Thus $R_{\widetilde{P}_{1}}\left(\rho_{1}\right)=$ $\chi_{V, \psi} \nu^{1 / 2} \zeta \otimes \omega_{\psi_{a}, 1}^{+}$.

In the same way, let $\rho_{2}$ be an irreducible common subquotient that

$$
\chi_{V, \psi} v^{1 / 2} \zeta \rtimes s p_{\zeta, 1} \quad \text { and } \quad \chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}
$$

have in common. Then $R_{\widetilde{P}_{1}}\left(\rho_{2}\right)=\chi_{V, \psi} v^{-1 / 2} \zeta \otimes s p_{\zeta, 1}$.
The representations $\chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} \otimes \omega_{0}$ and $\chi_{V, \psi} \zeta \mathrm{St}_{\mathrm{GL}(2)} \otimes \omega_{0}$ are irreducible and unitary. The multiplicity of $\chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} \otimes \omega_{0}$ in $R_{\widetilde{P}_{2}}\left(\chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}\right)$ is equal to 2 , which implies that length of $\chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}$ is 2 . Analogously, the length of $\chi_{V, \psi} \zeta \mathrm{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}$ also equals 2 .

Now we write $\chi_{V, \psi} \zeta \operatorname{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}=\rho_{1}+\rho_{3}$ and $\chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}=\rho_{2}+\rho_{4}$. Observe that

$$
\begin{aligned}
& R_{\widetilde{P}_{1}}\left(\rho_{3}\right)=2 \chi_{V, \psi} v^{1 / 2} \zeta \otimes s p_{\zeta, 1}+\chi_{V, \psi} v^{1 / 2} \zeta \otimes \omega_{\psi_{a}, 1}^{+} \\
& R_{\widetilde{P}_{1}}\left(\rho_{4}\right)=\chi_{V, \psi} v^{-1 / 2} \zeta \otimes s p_{\zeta, 1}+2 \chi_{V, \psi} v^{-1 / 2} \zeta \otimes \omega_{\psi_{a}, 1}^{+}
\end{aligned}
$$

We immediately get this:
Proposition 3.12. Let $\zeta \in \widehat{F}^{\rtimes}$ such that $\zeta^{2}=1_{F^{\times}}$. Then the representations
$\zeta 1_{\widetilde{\mathrm{GL}(2)}} \rtimes \omega_{0}, \quad \chi_{V, \psi} \zeta \mathrm{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}, \quad \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{\psi_{a}, 1}^{+} \quad \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes s p_{\zeta, 1}$
are reducible and $\chi_{V, \psi} \nu^{1 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}$ is a representation of length 4 . The representations

$$
\chi_{V, \psi} \zeta \operatorname{St}_{\mathrm{GL}(2)} \rtimes \omega_{0} \quad \text { and } \quad v^{1 / 2} \chi_{V, \psi} \zeta \rtimes \omega_{\psi_{a}, 1}^{+}
$$

(respectively $\chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes s p_{\zeta, 1}$ ) have exactly one irreducible subquotient in common, which is tempered and denoted by $\tau_{1}$ (respectively $\tau_{2}$ ). Observe that

$$
\tau_{1}=\left\langle\chi_{V, \psi} v^{1 / 2} \zeta ; \omega_{\psi_{a}, 1}^{+}\right\rangle \quad \text { and } \quad \tau_{2}=\left\langle\chi_{V, \psi} v^{1 / 2} \zeta, \chi_{V, \psi} v^{1 / 2} \zeta ; \omega_{0}\right\rangle
$$

## Also, the unique irreducible common subquotient of

$$
\chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0} \quad \text { and } \quad \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{\psi_{a}, 1}^{+}
$$

is a negative representation, which we denote by $\rho_{\text {neg. }}$. In $R$, we have

$$
\begin{array}{ll}
\chi_{V, \psi} \zeta \mathrm{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}=\tau_{1}+\tau_{2}, & \chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}=\rho_{\mathrm{neg}}+\left\langle\chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} ; \omega_{0}\right\rangle \\
\chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{\psi_{a}, 1}^{+}=\tau_{1}+\rho_{\mathrm{neg}}, & \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes s p_{\zeta, 1}=\tau_{2}+\left\langle\chi_{V, \psi} \zeta 1_{\mathrm{GL}(2)} ; \omega_{0}\right\rangle .
\end{array}
$$

(c) In this case we have

$$
\begin{aligned}
\chi_{V, \psi} \nu^{3 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0} & =\chi_{V, \psi} \nu^{3 / 2} \zeta \rtimes s p_{\zeta, 1}+\chi_{V, \psi} \nu^{3 / 2} \zeta \rtimes \omega_{\psi_{a}, 1}^{+} \\
& =\chi_{V, \psi} \nu \zeta \operatorname{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}+\chi_{V, \psi} \nu \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}
\end{aligned}
$$

Observe that $\chi_{V, \psi} \nu^{3 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}$ is a regular representation. So, it is a representation of the length $2^{2}=4$ by [Tadić 1998b] (there only the techniques of Jacquet modules were used, so they can be applied in our case). Since the irreducible subquotients of $\chi_{V, \psi} \nu^{3 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}$ are $\left\langle\chi_{V, \psi} \nu^{3 / 2} \zeta, \chi_{V, \psi} \nu^{1 / 2} \zeta ; \omega_{0}\right\rangle, \quad\left\langle\chi_{V, \psi} \nu \zeta 1_{\mathrm{GL}(2)} ; \omega_{0}\right\rangle, \quad \omega_{\psi_{a}, 2}^{+}, \quad\left\langle\chi_{V, \psi} \nu^{3 / 2} \zeta ; \omega_{\psi_{a}, 1}^{+}\right\rangle$, using Jacquet modules we easily obtain the following proposition:
Proposition 3.13. Let $\zeta \in \widehat{F^{×}}$ such that $\zeta^{2}=1_{F^{\star}}$. Then the representations

$$
\begin{array}{ll}
\chi_{V, \psi} \nu^{3 / 2} \zeta \rtimes s p_{\zeta, 1}, & \chi_{V, \psi} \nu^{3 / 2} \zeta \rtimes \omega_{\psi_{a}, 1}^{+} \\
\chi_{V, \psi} \nu \zeta \mathrm{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}, & \chi_{V, \psi} \nu \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}
\end{array}
$$

are reducible and $\chi_{V, \psi} \nu^{3 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}$ is a representation of length 4. The unique irreducible common subquotient of the representations $\chi_{V, \psi} \nu^{3 / 2} \zeta \rtimes s p_{\zeta, 1}$ and $\chi_{V, \psi} \nu \zeta \mathrm{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}$ is square-integrable. In $R$, we have

$$
\begin{aligned}
\chi_{V, \psi} \nu^{3 / 2} \zeta \rtimes s p_{\zeta, 1} & =\left\langle\chi_{V, \psi} \nu^{3 / 2} \zeta, \chi_{V, \psi} \nu^{1 / 2} \zeta ; \omega_{0}\right\rangle+\left\langle\chi_{V, \psi} \nu \zeta 1_{\mathrm{GL}(2)} ; \omega_{0}\right\rangle \\
\chi_{V, \psi} \nu^{3 / 2} \zeta \rtimes \omega_{\psi_{a}, 1}^{+} & =\omega_{\psi_{a}, 2}^{+}+\left\langle\chi_{V, \psi} \nu^{3 / 2} \zeta ; \omega_{\psi_{a}, 1}^{+}\right\rangle \\
\chi_{V, \psi} \nu \zeta \operatorname{St}_{\mathrm{GL}(2)} \rtimes \omega_{0} & =\left\langle\chi_{V, \psi} \nu^{3 / 2} \zeta, \chi_{V, \psi} \nu^{1 / 2} \zeta ; \omega_{0}\right\rangle+\left\langle\chi_{V, \psi} \nu^{3 / 2} \zeta ; \omega_{\psi_{a}, 1}^{+}\right\rangle \\
\chi_{V, \psi} \nu \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0} & =\left\langle\chi_{V, \psi} \nu \zeta 1_{\mathrm{GL}(2)} ; \omega_{0}\right\rangle+\omega_{\psi_{a}, 2}^{+}
\end{aligned}
$$

(d) In this case,

$$
\chi_{V, \psi} \nu \zeta \times \chi_{V, \psi} \zeta \rtimes \omega_{0}=\chi_{V, \psi} \nu^{1 / 2} \zeta \operatorname{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}+\chi_{V, \psi} \nu^{1 / 2} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}
$$

Since it isn't known yet if the results related to the $R$-groups [Goldberg 1994] also hold for metaplectic groups, this case will not be solved using only the method Jacquet modules. Tadić [1998a] used a combination of Jacquet modules techniques and knowledge about $R$-groups for symplectic groups to determine the composition
series of the representations similar to this one (for symplectic groups). We resolve this case using theta correspondence.
Lemma 3.14. The following equalities hold:
(1) $\Theta\left(\zeta \nu \otimes \zeta \rtimes 1, R_{P_{1}}\left(\omega_{2,2}\right)\right)=\chi_{V, \psi} \nu^{-1} \zeta \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$,
(2) $\Theta\left(\zeta \nu^{-1} \otimes \zeta \rtimes 1, R_{P_{1}}\left(\omega_{2,2}\right)\right)=\chi_{V, \psi} \nu \zeta \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$,
(3) $\Theta\left(\zeta \otimes \zeta \nu \rtimes 1, R_{P_{1}}\left(\omega_{2,2}\right)\right)=\chi_{V, \psi} \zeta \times \chi_{V, \psi} \nu \zeta \rtimes \omega_{0}$,
(4) $\Theta\left(\chi_{V, \psi} \nu \zeta \otimes \chi_{V, \psi} \zeta \rtimes \omega_{0}, R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right)\right)=\zeta \nu^{-1} \times \zeta \rtimes 1$,
(5) $\Theta\left(\chi_{V, \psi} v^{-1} \zeta \otimes \chi_{V, \psi} \zeta \rtimes \omega_{0}, R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right)\right)=\zeta v \times \zeta \rtimes 1$,
(6) $\Theta\left(\chi_{V, \psi} \zeta \otimes \chi_{V, \psi} \nu \zeta \rtimes \omega_{0}, R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right)\right)=\zeta \times \zeta \nu \rtimes 1$.

Proof. Recall that $R_{P_{1}}\left(\omega_{2,2}\right)$ has the filtration in which

- $I_{10}=v^{1 / 2} \otimes \omega_{2,1}$ is the quotient, and
- $I_{11}=\operatorname{Ind}_{G L(1) \times \widetilde{P_{1}} \times O(3)}^{M_{1} \times \widetilde{\operatorname{Sp}(2)}}\left(\chi_{V, \psi} \Sigma_{1}^{\prime} \otimes \omega_{1,1}\right)$ is the subrepresentation.

We will prove (1); the proofs of (2)-(6) are analogous. In the same way as in the second proof of Proposition 3.4, we get

$$
\Theta\left(\zeta v \otimes \zeta \rtimes 1, R_{P_{1}}\left(\omega_{2,2}\right)\right)=\Theta\left(\zeta v \otimes \zeta \rtimes 1, I_{11}\right)
$$

so it is sufficient to show $\Theta\left(\zeta \nu \otimes \zeta \rtimes 1, I_{11}\right)=\chi_{V, \psi} \nu^{-1} \zeta \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$. It can be seen easily that there is an GL(1) $\times \widetilde{M}_{1} \times O(3)$-invariant epimorphism

$$
\chi_{V, \psi} \Sigma_{1}^{\prime} \otimes \omega_{1,1} \rightarrow \chi_{V, \psi} \zeta \nu^{-1} \otimes \chi_{V, \psi} \zeta \rtimes \omega_{0} \otimes \zeta \nu \otimes \zeta \rtimes 1 .
$$

Consequently, we get an $M_{1} \times \widetilde{\mathrm{Sp(2)}}$-invariant epimorphism
$I_{11}=\operatorname{Ind}_{\mathrm{GL}(1) \times \widetilde{P}_{1} \times O(3)}^{M_{1} \widetilde{\operatorname{Sp}^{2}(2)}}\left(\chi_{V, \psi} \Sigma_{1}^{\prime} \otimes \omega_{1,1}\right) \rightarrow \zeta \nu \otimes \zeta \rtimes 1 \otimes \chi_{V, \psi} \zeta v^{-1} \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$, so we conclude that $\chi_{V, \psi} \zeta \nu^{-1} \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$ is a quotient of $\Theta\left(\zeta \nu \otimes \zeta \rtimes 1, I_{11}\right)$.

We prove that $\Theta\left(\zeta \nu \otimes \zeta \rtimes 1, I_{11}\right)$ is also a quotient of $\chi_{V, \psi} \zeta \nu^{-1} \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$. Let $\varphi \in \operatorname{Hom}\left(I_{11}, \zeta \nu \otimes \zeta \rtimes 1 \otimes \Theta\left(\zeta v \otimes \zeta \rtimes 1, I_{11}\right)\right)$. Using the second Frobenius reciprocity, as before, we get

$$
\begin{aligned}
& \operatorname{Hom}\left(I_{11}, \zeta \nu \otimes \zeta \rtimes 1 \otimes \Theta\left(\zeta \nu \otimes \zeta \rtimes 1, I_{11}\right)\right) \\
& \cong \operatorname{Hom}\left(\chi_{V, \psi} \Sigma_{1}^{\prime} \otimes \omega_{1,1}, \zeta \nu \otimes \zeta \rtimes 1 \otimes R_{\widetilde{P}_{1}}\left(\Theta\left(\zeta \nu \otimes \zeta \rtimes 1, I_{11}\right)\right)\right)
\end{aligned}
$$

let $\varphi_{0}$ be an element corresponding to $\varphi$. Since the representations $\zeta \rtimes 1$ and $\chi_{V, \psi} \zeta \rtimes \omega_{0}$ are irreducible, the image of $\varphi_{0}$ equals $\zeta \nu \otimes \zeta \rtimes 1 \otimes \chi_{V, \psi} \zeta \nu^{-1} \otimes$ $\chi_{V, \psi} \zeta \rtimes \omega_{0}$. Reasoning as before, we get that the image of $\varphi$ is a quotient of $\zeta \nu \otimes \zeta \rtimes 1 \otimes \chi_{V, \psi} \zeta \nu^{-1} \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$. Finally, $\Theta\left(\zeta \nu \otimes \zeta \rtimes 1, I_{11}\right)$ is a quotient of $\chi_{V, \psi} \zeta \nu^{-1} \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$. Hence $\Theta\left(\zeta \nu \otimes \zeta \rtimes 1, I_{11}\right)=\chi_{V, \psi} \zeta \nu^{-1} \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$.

Proposition 3.15. Let $\zeta \in \widehat{F^{\rtimes}}$ such that $\zeta^{2}=1_{F^{\times}}$. Then the representations

$$
\chi_{V, \psi} \nu^{1 / 2} \zeta \mathrm{St}_{\mathrm{GL}(2)} \rtimes \omega_{0} \quad \text { and } \quad \chi_{V, \psi} \nu^{1 / 2} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}
$$

are irreducible and $\chi_{V, \psi} \nu \zeta \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$ is a representation of length 2 . Also

$$
\begin{aligned}
\chi_{V, \psi} v^{1 / 2} \zeta \mathrm{St}_{\mathrm{GL}(2)} \rtimes \omega_{0} & =\left\langle\chi_{V, \psi} \nu \zeta ; \chi_{V, \psi} \zeta \rtimes \omega_{0}\right\rangle, \\
\chi_{V, \psi} v^{1 / 2} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0} & =\left\langle\chi_{V, \psi} v^{1 / 2} \zeta 1_{\mathrm{GL}(2)} ; \omega_{0}\right\rangle .
\end{aligned}
$$

Proof. Suppose on the contrary that the representation $\chi_{V, \psi} \nu^{1 / 2} \zeta \operatorname{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}$ reduces. Jacquet modules imply that length of this representation is at most 2. Choose $\pi_{1}$ and $\pi_{2}$ so the equality $\chi_{V, \psi} \nu^{1 / 2} \zeta \operatorname{St} \mathrm{GL}_{(2)} \rtimes \omega_{0}=\pi_{1}+\pi_{2}$ holds in $R$. Also suppose $R_{\widetilde{P}_{1}}\left(\pi_{1}\right)=\chi_{V, \psi} \zeta \otimes \chi_{V, \psi} \zeta \nu \rtimes \omega_{0}$ and $R_{\widetilde{P}_{1}}\left(\pi_{2}\right)=\chi_{V, \psi} \zeta \nu \otimes \chi_{V, \psi} \zeta \rtimes \omega_{0}$. Frobenius reciprocity implies

$$
\operatorname{Hom}\left(\omega_{2,2}, \pi_{1} \otimes \zeta \times \zeta v^{-1} \rtimes 1\right) \cong \operatorname{Hom}\left(R_{P_{1}}\left(\omega_{2,2}\right), \pi_{1} \otimes \zeta \otimes \zeta v^{-1} \rtimes 1\right)
$$

Using Lemma 3.14 we obtain
$\operatorname{Hom}\left(R_{P_{1}}\left(\omega_{2,2}\right), \pi_{1} \otimes \zeta \otimes \zeta \nu^{-1} \rtimes 1\right) \cong \operatorname{Hom}\left(\chi_{V, \psi} \zeta \times \chi_{V, \psi} \nu^{-1} \zeta \rtimes \omega_{0}, \pi_{1}\right) \neq 0$, because $\pi_{1}$ is a quotient of $\chi_{V, \psi} \zeta \times \chi_{V, \psi} \nu \zeta \rtimes \omega_{0}$. So, $\Theta\left(\pi_{1}, 2\right) \neq 0$.

The representation $\chi_{V, \psi} \zeta \otimes \chi_{V, \psi} \zeta \nu \rtimes \omega_{0} \otimes \Theta\left(\pi_{1}, 2\right)$ is a quotient of $R_{\widetilde{P}_{1}}\left(\omega_{2,2}\right)$. Lemma 3.14 implies that $\Theta\left(\pi_{1}, 2\right)$ is a quotient of $\zeta \times \zeta \nu \rtimes 1$. Listing quotients of $\zeta \times \zeta \nu \rtimes 1$ we get the possibilities
(a) $\Theta\left(\pi_{1}, 2\right)=\zeta \times \zeta \nu \rtimes 1$,
(b) $\Theta\left(\pi_{1}, 2\right)=v^{1 / 2} \zeta \mathrm{St}_{\mathrm{GL}(2)} \rtimes 1$,
(c) $\Theta\left(\pi_{1}, 2\right)=v^{-1 / 2} \zeta 1_{\mathrm{GL}(2)} \rtimes 1$.

Suppose that (a) holds. Obviously, $\pi_{1} \otimes \zeta v^{-1} \otimes \zeta \rtimes 1$ is then a quotient of $R_{P_{1}}\left(\omega_{2,2}\right)$, since it is a quotient of $\pi_{1} \otimes R_{P_{1}}(\zeta \times \zeta \nu \rtimes 1)$. This implies that $\pi_{1}$ is a quotient of $\chi_{V, \psi} \zeta \nu \times \chi_{V, \psi} \zeta \rtimes \omega_{0}$ and $R_{\widetilde{P}_{1}}\left(\pi_{1}\right)$ contains $\chi_{V, \psi} \zeta \nu^{-1} \otimes \chi_{V, \psi} \zeta \rtimes \omega_{0}$. This contradicts our assumption on $\pi_{1}$.

Similarly, using Jacquet modules, we obtain contradiction with (b) and (c). So, $\chi_{V, \psi} \nu^{1 / 2} \zeta \mathrm{St}_{\mathrm{GL}(2)} \rtimes \omega_{0}$ is irreducible.

Irreducibility of $\chi_{V, \psi} \nu^{1 / 2} \zeta 1_{\mathrm{GL}(2)} \rtimes \omega_{0}$ can be proved in the same way.

## 4. Unitary dual supported in minimal parabolic subgroup

Let $\pi$ be an irreducible genuine admissible representation of $\widetilde{\operatorname{Sp}(n)}$. We recall that the contragredient representation is denoted by $\tilde{\pi}$. We write $\bar{\pi}$ for the complex conjugate representation of the representation $\pi$. The representation $\pi$ is called Hermitian if $\pi \simeq \widetilde{\bar{\pi}}$. It is well known that every unitary representation is Hermitian. For a deeper discussion, we refer the reader to [Muić and Tadić 2007].

Suppose that $\Delta_{1}, \ldots, \Delta_{k}$ is a sequence of segments such that $e\left(\Delta_{1}\right) \geq \cdots \geq$ $e\left(\Delta_{k}\right)>0$, let $\sigma_{\text {neg }}$ be a negative representation of some $\widehat{\operatorname{Sp}\left(n^{\prime}\right)}$. From [Hanzer and Muić 2010, Theorem 4.5(v)], we directly get

$$
\overline{\left\langle\Delta_{1}, \ldots, \Delta_{k} ; \sigma_{\mathrm{neg}}\right\rangle}=\left\langle\bar{\Delta}_{1}, \ldots, \bar{\Delta}_{k} ; \bar{\sigma}_{\mathrm{neg}}\right\rangle .
$$

Also, we have an epimorphism $\left\langle\widetilde{\Delta_{1}}\right\rangle \times \cdots \times\left\langle\widetilde{\Delta_{k}}\right\rangle \rtimes \widetilde{\sigma}_{\text {neg }} \rightarrow\left\langle\Delta_{1}, \ldots, \Delta_{k} ; \sigma_{\text {neg }}\right\rangle^{\sim}$. We know that the group $\operatorname{GSp}(n)$ acts on $\widetilde{\operatorname{Sp}(n)}$, by [Mœglin et al. 1987, II.1(3)]. Moreover, by [ibid., page 92], this action extends to the action on irreducible representations, which is equivalent to taking contragredients. We choose an element $\eta^{\prime}=(1, \eta) \in \operatorname{GSp}(n)$, where $\eta \in \operatorname{GSp}\left(n^{\prime}\right)$ is an element with similitude equal to -1 , and 1 denotes the identity acting on the GL part. Thus, we obtain an epimorphism

$$
\alpha\left\langle\widetilde{\Delta_{1}}\right\rangle \times \cdots \times \alpha\left\langle\widetilde{\Delta_{k}}\right\rangle \rtimes \widetilde{\sigma}_{\text {neg }}^{\eta} \rightarrow\left\langle\Delta_{1}, \ldots, \Delta_{k} ; \sigma_{\text {neg }}\right\rangle^{\sim \eta^{\prime}}
$$

Since $\widetilde{\sigma}_{\text {neg }}^{\eta} \simeq \sigma_{\text {neg }}$, we have $\left\langle\Delta_{1}, \ldots, \Delta_{k} ; \sigma_{\text {neg }}\right\rangle^{\sim}=\left\langle\alpha \Delta_{1}, \ldots, \alpha \Delta_{k} ; \widetilde{\sigma}_{\text {neg }}\right\rangle$.
Remark. When we are dealing with the action of the group of similitudes on the symplectic groups, $\alpha$ does not appear in the situation similar to the one above. However, since the action of $\operatorname{GSp}(n)$ on the metaplectic group is not trivial on its center $\mu_{2}$, one has to compare the action $\eta$ on the metaplectic part of the Levi subgroup with the action of $\eta^{\prime}$ on the whole Levi subgroup. The calculation is not very complicated and resembles the calculations in [Hanzer and Muić 2010, Lemma 3.2].

First, we classify Hermitian irreducible genuine representations.
Proposition 4.1. Let $\chi, \zeta, \zeta_{1}, \zeta_{2} \in \widehat{F}^{\times}$such that $\zeta^{2}=\zeta_{i}^{2}=1_{F^{\times}}$for $i=1,2$, with $\zeta_{1}$ and $\zeta_{2}$ not necessarily different. Let $s, s_{1}, s_{2}>0$. The following families of representations are Hermitian and exhaust all irreducible Hermitian genuine representations of $\widetilde{\mathrm{Sp}(2)}$ supported in the minimal parabolic subgroup $\widetilde{P_{(1,1)}}$ :
(1) irreducible tempered representations supported in $\widetilde{P_{(1,1)}}$,
(2) $\left\langle\chi_{V, \psi} \nu^{s} \chi, \chi_{V, \psi} \nu^{s} \chi^{-1} ; \omega_{0}\right\rangle$,
(3) $\left\langle\chi_{V, \psi} \nu^{s_{1}} \zeta_{1}, \chi_{V, \psi} \nu^{s_{2}} \zeta_{2} ; \omega_{0}\right\rangle$,
(4) $\left\langle\chi_{V, \psi} \nu^{s} \zeta 1_{\mathrm{GL}(2)} ; \omega_{0}\right\rangle$,
(5) $\left\langle\chi_{V, \psi} \nu^{s} \zeta ; \chi_{V, \psi} \chi \rtimes \omega_{0}\right\rangle$,
(6) $\left\langle\chi_{V, \psi} \nu^{s} \zeta ; \omega_{\psi_{a}, 1}^{+}\right\rangle$,
(7) $\omega_{\psi_{a}, 2}^{+}$.

Proof. Using the reasoning before this proposition, we see that a representation $\left\langle\Delta_{1}, \ldots, \Delta_{k} ; \sigma_{\text {neg }}\right\rangle$ is Hermitian, if and only if

$$
\left\langle\Delta_{1}, \ldots, \Delta_{k} ; \sigma_{\text {neg }}\right\rangle=\left\langle\alpha \bar{\Delta}_{1}, \ldots, \alpha \bar{\Delta}_{k} ; \overline{\widetilde{\sigma}}_{\text {neg }}\right\rangle
$$

The representation $\sigma_{\text {neg }}$ also has to be Hermitian. Now we just check this requirement on the set of all irreducible representations of $\widetilde{\operatorname{Sp(} 2)}$ with the support in the minimal parabolic subgroup; we have classified them in the previous section. For example, if we analyze the representation $\Pi=\left\langle\chi_{V, \psi} \nu^{s_{1}} \chi_{1}, \chi_{V, \psi} \nu^{s_{2}} \chi_{2} ; \omega_{0}\right\rangle$ with $s_{2} \geq s_{1}>0$, we have

$$
\widetilde{\Pi}=\left\langle\alpha \overline{\chi_{V, \psi} \nu^{s_{1}} \chi_{1}}, \alpha \overline{\chi_{V, \psi} \nu^{s_{2}} \chi_{2}} ; \omega_{0}\right\rangle .
$$

Now we see that this representation is isomorphic to $\Pi$ if and only if $\chi_{1}^{2}=1=\chi_{2}^{2}$ or $s_{1}=s_{2}$ and $\chi_{1}^{-1}=\chi_{2}$. This gives us the second and the third case from the proposition. All other cases are dealt with analogously.
Theorem 4.2. Let $\chi, \zeta, \zeta_{1}, \zeta_{2} \in\left(\widehat{F^{\varnothing}}\right)$ such that $\zeta^{2}=\zeta_{i}^{2}=1_{F^{\times}}$for $i=1$, 2 , with $\zeta_{1}$ and $\zeta_{2}$ not necessarily different. The following families of representations are unitary and exhaust all irreducible unitary genuine representations of $\widetilde{\operatorname{Sp}(2)}$ that are supported in the minimal parabolic subgroup $\widetilde{P_{(1,1)}}$ :
(1) irreducible tempered representations supported in $\widetilde{P_{(1,1)}}$.
(2) $\left\langle\chi_{V, \psi} \nu^{s} \chi, \chi_{V, \psi} \nu^{s} \chi^{-1} ; \omega_{0}\right\rangle$ for $0<s \leq 1 / 2$,
(3) $\left\langle\chi_{V, \psi} \nu^{s_{1}} \zeta_{1}, \chi_{V, \psi} \nu^{s_{2}} \zeta_{2} ; \omega_{0}\right\rangle$ for $s_{2} \leq s_{1}$ and $0<s_{1} \leq 1 / 2$,
(4) $\left\langle\chi_{V, \psi} v^{s} \zeta ; \chi_{V, \psi} \chi \rtimes \omega_{0}\right\rangle$ for $0<s \leq 1 / 2$,
(5) $\left\langle\chi_{V, \psi} \nu^{s} \zeta ; \omega_{\psi_{a}, 1}^{+}\right\rangle$for $s \leq 1 / 2$,
(6) $\omega_{\psi_{a}, 2}^{+}$.

Proof. We first review some basic facts of representation theory of reductive groups, which directly carry over to the case of metaplectic groups.

Unitarizability of the complementary series. As explained in detail in [Tadić 1993, Section 3], it is enough to have a continuous family of $\widetilde{\operatorname{Sp}(2)}$-invariant hermitian forms (and the representations should be realized on one space - a compact picture). Then, a linear algebra argument (involving finite-dimensional representations of a compact subgroup) ensures that if this family of hermitian forms is positive definite at one point, it has to be positive definite everywhere, and this finishes the argument. So, we first have to show that, the restriction of an irreducible admissible (hermitian) representation $\pi$ of $\widetilde{\mathrm{Sp}(2)}$ to the inverse image $\widetilde{K}$ in $\widetilde{\mathrm{Sp}(2)}$ of a maximal, good compact subgroup of $\mathrm{Sp}(2)$ (for example, $K=\mathrm{Sp}\left(2, O_{F}\right)$, where $O_{F}$ is the ring of integers of $F$ ) decomposes into a direct sum of irreducible representations of $\widetilde{K}$ with finite multiplicities. But this follows directly from the admissibility of the representation $\pi$. Second, we have to have a way to form a continuous families of hermitian forms. This is obtained using intertwining operators, in the same way as for the algebraic groups. To define them, we note that the unipotent radicals of the standard parabolic subgroups of $\mathrm{Sp}(2)$ are split in $\widetilde{\mathrm{Sp}(2)}$.

Then we can define standard intertwining operators for the complex argument deep enough in the Weyl chamber in the same way as in [Shahidi 1981]. These operators can be meromorphically continued, using the results on filtration via Bruhat cells; see [Casselman 1995, Section 6 and 7] or [Muić 2008]. The arguments in the last reference carry over to the metaplectic groups without change, through the splitting of the unipotent radicals and Frobenius reciprocity, also valid for the metaplectic group. This passage in the construction of the intertwining operators from the linear case to the case of metaplectic groups is explained in detail also in [Zorn 2007; 2010]. We now illustrate how the hermitian form is defined. For example, suppose that $\chi_{1}, \chi_{2} \in \hat{F}^{\times}$such that $\chi_{1}^{2}=\chi_{2}^{2}=1$, so that, for the longest element $w_{0}$ of the Weyl group, we have a map $A\left(s_{1}, s_{2}, \chi_{1}, \chi_{2}, w_{0}\right)$ from

$$
\chi_{V, \psi} \chi_{1} v^{s_{1}} \times \chi_{V, \psi} \chi_{2} v^{s_{2}} \rtimes \omega_{0} \quad \text { to } \quad \chi_{V, \psi} v^{-s_{1}} \chi_{1} \times \chi_{V, \psi} v^{-s_{2}} \chi_{2} \rtimes \omega_{0}
$$

Let $f_{s_{1}, s_{2}}, g_{s_{1}, s_{2}}$ be sections from the compact picture of the induced representation $\chi_{V, \psi} \chi_{1} \nu^{s_{1}} \times \chi_{V, \psi} \chi_{2} \nu^{s_{2}} \rtimes \omega_{0}$. Then, a hermitian form indexed by $\left(s_{1}, s_{2}\right)$ is defined by

$$
(f, g)_{\left(s_{1}, s_{2}\right)}=\int_{\widetilde{K}} A\left(s_{1}, s_{2}, \chi_{1}, \chi_{2}, w_{0}\right) f_{s_{1}, s_{2}}(\widetilde{k}) \overline{g_{s_{1}, s_{2}}(\widetilde{k})} \mathrm{d} \widetilde{k}
$$

The $\widetilde{\mathrm{Sp}(2)}$-invariance of this form follows from [Casselman 1981, Theorem 2.4.2] (in the context of totally disconnected groups) and then from Proposition 3.1.3 therein, after normalizing the measure on $\widetilde{P}$ so that $\widetilde{P} \cap \widetilde{K}$ is of volume one (since $\widetilde{P} \widetilde{K}=\widetilde{G})$.

Unitarizability of the ends of the complementary series. For the reductive algebraic groups, the unitarizability of the ends of the complementary series is proved by Miličić [1973] using $C^{*}$-algebra arguments. To avoid that (although this argument may also apply in the case of metaplectic groups), we use a similar result, that is, [Tadić 1986, Theorem 2.5]. The proof of this result relies on calculations of the limits of the operators acting on the finite-dimensional complex vector spaces, and the only requirements are admissibility of the irreducible smooth representations in question (our Lemma 2.1) and a result of Bernstein about uniform admissibility. But, since we do not require the generality in which that theorem is posed, we actually do not need Bernstein's argument, since we are dealing with the family of representations in the complementary series - all of them have the same restriction to the compact open subgroup $K_{1}$ (which splits in $\widetilde{\mathrm{Sp}}$ ), and the requirement labeled $(*)$ there is automatically fulfilled. Hecke algebra $H\left(\widetilde{\mathrm{Sp}(2)}, K_{1}\right)$ is defined in the same way as in the case of reductive groups.

The asymptotics of the matrix coefficients of the representations of the metaplectic group. These can also be estimated in terms of Jacquet modules of the representations [Casselman 1981, Section 4]. Indeed, the arguments there rely on the calculation of the spaces of the coinvariants for the unipotent subgroups (which split in $\widetilde{\mathrm{Sp}(2)})$ and spaces of vectors fixed by some small compact subgroups of $\widetilde{\mathrm{Sp}(2)}$. These subgroups can always be taken to belong to the maximal compact subgroup of $\operatorname{Sp}(2)$ that splits in $\widetilde{\mathrm{Sp}(2)}$ (if the residual characteristic is odd) or to some smaller open compact subgroup that splits, so we actually take the fixed vectors by these splittings of compact subgroups.

On the other hand, the reducibility points of the principal series for $\mathrm{SO}(5)$ are analogous to those for $\widetilde{\mathrm{Sp}(2)}$, so the unbounded areas of [Matić 2010, Figure 1 of Theorem 3.5], through the asymptotics explained above, give rise to the representations with unbounded matrix coefficients. Thus none of these representations are unitarizable (because of the continuity of the hermitian forms on these unbounded parts).

The arguments above (plus the irreducibility of the unitary principal series) were the main tools in the proof of [Matic 2010, Theorem 3.5]; there was only the problem of how to deal with certain isolated representations.

Recall that in $R$ we have $\chi_{V, \psi} \nu^{3 / 2} \zeta \times \chi_{V, \psi} \nu^{1 / 2} \zeta \rtimes \omega_{0}$ is equal to $\left\langle\chi_{V, \psi} \nu^{3 / 2} \zeta, \chi_{V, \psi} \nu^{1 / 2} \zeta ; \omega_{0}\right\rangle+\omega_{\psi_{a}, 2}^{+}+\left\langle\chi_{V, \psi} \nu^{3 / 2} \zeta ; \omega_{\psi_{a}, 1}^{+}\right\rangle+\left\langle\chi_{V, \psi} \nu \zeta 1_{\mathrm{GL}(2)} ; \omega_{0}\right\rangle$, where $\left\langle\chi_{V, \psi} \nu^{3 / 2} \zeta, \chi_{V, \psi} \nu^{1 / 2} \zeta ; \omega_{0}\right\rangle$ and $\omega_{\psi_{a}, 2}^{+}$are unitarizable. Observe that the representation $\left\langle\chi_{V, \psi} \nu^{3 / 2} \zeta ; \omega_{\psi_{a}, 1}^{+}\right\rangle$(respectively, $\left.\left\langle\chi_{V, \psi} \nu \zeta 1_{\mathrm{GL}(2)} ; \omega_{0}\right\rangle\right)$ has Jacquet modules analogous to those of the representation $L\left(\delta\left(\left[\nu^{1 / 2}, v^{3 / 2}\right]\right), 1\right)$ (respectively, $\left.L\left(v^{3 / 2}, \mathrm{St}_{\mathrm{SO}(3)}\right)\right)$ of the group $\mathrm{SO}(5)$. Hence, nonunitarizability of these two representations can be proved analagously to the nonunitarizability of the representations $L\left(\delta\left(\left[v^{1 / 2}, v^{3 / 2}\right]\right), 1\right)$ and $L\left(v^{3 / 2}, \mathrm{St}_{\mathrm{SO}(3)}\right)$, which is a special case of [Hanzer and Tadić 2010, Propositions 4.1 and 4.6]. The arguments used there rely on the Jacquet modules method, which also applies to group $\widetilde{\mathrm{Sp}(2)}$, and the simple fact that every unitary representation is also semisimple.

## 5. Unitary dual supported in maximal parabolic subgroups

5.1. The Siegel case. Using [Hanzer and Muić 2009], [Matić 2010, Proposition 4.1] and previously discussed issues of complementary series and nonunitarizability of the representations indexed by the (geometrically) unbounded pieces of the plane, we directly get the following:

Proposition 5.1. Let $\rho$ be an irreducible cuspidal representation of $\operatorname{GL}(2, F)$. There is at most one $s \geq 0$ such that $\chi_{V, \psi} \nu^{s} \rho \rtimes \omega_{0}$ reduces. One of the following holds:
(1) If $\rho$ is not self-dual, then $\chi_{V, \psi} \rho \rtimes \omega_{0}$ is irreducible and unitarizable. Also, the representations $\chi_{V, \psi} \nu^{s} \rho \rtimes \omega_{0}$ are irreducible and nonunitarizable for $s>0$.
(2) If $\rho$ is self-dual and $\omega_{\rho}=1$, where $\omega_{\rho}$ denotes the central character of $\rho$, then the representation $\chi_{V, \psi} \rho \rtimes \omega_{0}$ reduces, while all of the representations $\chi_{V, \psi} \nu^{s} \rho \rtimes \omega_{0}$ are nonunitarizable for $s>0$.
(3) If $\rho$ is self-dual and $\omega_{\rho} \neq 1$, then the unique $s \geq 0$ such that $\chi_{V, \psi} \nu^{s} \rho \rtimes \omega_{0}$ reduces is equal to $1 / 2$. For $0 \leq s \leq 1 / 2$, the representations $\chi_{V, \psi} \nu^{s} \rho \rtimes \omega_{0}$ are all unitarizable; for $s>1 / 2$, the representations $\chi_{V, \psi} \nu^{s} \rho \rtimes \omega_{0}$ are all nonunitarizable. All irreducible subquotients of $\chi_{V, \psi} \nu^{1 / 2} \rho \rtimes \omega_{0}$ are unitarizable.
5.2. The non-Siegel case. Hanzer and Muić [2009, Section 5.2] determine the reducibility points of the representations $\chi_{V, \psi} \nu^{s} \zeta \rtimes \pi$, where $s \in \mathbb{R}, \zeta \in \widehat{F^{x}}$ and $\pi$ is an irreducible cuspidal representation of $\widetilde{\operatorname{Sp}(1)}$. After determining the reducibility points, the unitarizability of the induced representations and irreducible subquotients follow in the same way as in Proposition 5.1. For the convenience of the reader, we write down all the results.

To the fixed quadratic character $\chi_{V}$ we attach, as in [Kudla 1996, Chapter V], two odd-orthogonal towers, the + -tower and the --tower. We denote by $\Theta^{ \pm}(\pi)$ the first appearance of the representation $\Theta(\pi)$ in the respective $\pm$-tower. Analogously, for $r \geq 0$, we denote by $\Theta^{ \pm}(\pi, r)$ the lift of the representation $\pi$ to the $r$-th level of the respective $\pm$-tower.

Since the representation $\chi_{V, \psi} \nu^{s} \zeta \rtimes \pi$ is irreducible for $\zeta^{2} \neq 1$, we suppose $\zeta^{2}=1$ and consider two cases:
(a) $\zeta \neq 1$. Applying [Hanzer and Muić 2009, Theorem 3.5] we see that $\chi_{V, \psi} \nu^{s} \zeta \rtimes \pi$ reduces if and only if $\zeta \nu^{s} \rtimes \Theta^{+}(\pi)$ reduces (in the + -tower) if and only if $\zeta \nu^{s} \rtimes$ $\Theta^{-}(\pi)$ reduces (in the --tower).

Now $\Theta^{+}(\pi)$ is an irreducible cuspidal representation of some of the groups $O(1), O(3)$ or $O(5)$. Let $r$ denote the first occurrence of a nonzero lift of $\pi$ in the odd orthogonal + -tower. We have several cases depending on $r$ :

- If $r=0$, that is, if $\pi$ equals $\omega_{\psi_{a}^{-}, 1}$, which is an odd part of the Weil representation attached to additive character $\psi$, then $\Theta^{+}(\pi, 0)=\operatorname{sgn}_{O(1)}$, so the representation $\zeta \nu^{s} \rtimes \operatorname{sgn}_{O(1)}$ reduces if and only if $\zeta \nu^{s} \rtimes 1$ reduces in SO (3). It is well known that this representation reduces when $s=1 / 2$.
- If $r=1$, the representation $\zeta \nu^{s} \rtimes \pi$ reduces if and only if the representation $\left.\zeta \nu^{s} \rtimes \Theta^{+}(\pi, 1)\right|_{\mathrm{SO}(3)}$ reduces. As in [Matić 2010], we obtain that the unique $s$ such that $\zeta \nu^{s} \rtimes \pi$ reduces is equal to $1 / 2$.
- If $r=2$, the representation $\zeta \nu^{s} \rtimes \pi$ reduces if and only if $\zeta \nu^{s} \rtimes \Theta^{+}(\pi, 2)$ reduces, and that is if and only if $\left.\zeta \nu^{s} \rtimes \Theta^{+}(\pi, 2)\right|_{\operatorname{SO}(5)}$ reduces. We do
not know if the representation $\Theta^{+}(\pi, 2)$ is generic, so we turn our attention to the representation $\zeta \nu^{s} \rtimes \Theta^{-}(\pi, 0)$, because we know that $\Theta^{-}(\pi, 0)$ is a nonzero representation of $O(1)$ (since $\pi$ is cuspidal, the dichotomy conjecture holds). Recall that $\zeta \nu^{s} \rtimes \Theta^{-}(\pi, 0)$ reduces for $s=0$ if and only if $\mu\left(s, \zeta \otimes \Theta^{-}(\pi, 0)\right) \neq 0$ for $s=0$ and that $\zeta \nu^{s} \rtimes \Theta^{-}(\pi, 0)$ reduces for $s_{0}>0$ if and only if $\mu\left(s, \zeta \otimes \Theta^{-}(\pi, 0)\right)$ has a pole for $s=s_{0}$. In the same way as in [Hanzer and Muić 2009, Section 5.2, case 3], we obtain $\mu\left(s, \zeta \otimes \Theta^{-}(\pi, 0)\right) \cong \mu\left(s, \zeta \otimes J L\left(\Theta^{-}(\pi, 0)\right)\right)$, where $\left.J L\left(\Theta^{-}(\pi, 0)\right)\right)$ denotes the Jacquet-Langlands lift of $\Theta^{-}(\pi, 0)$ ). Now we consider two possibilities:
(I) $\Theta^{-}(\pi, 0)$ is not one-dimensional. In this case, $J L\left(\Theta^{-}(\pi, 0)\right)$ is a cuspidal generic representation of $\mathrm{SO}(3, F)$ and the reducibility point is $s=1 / 2$.
(II) $\Theta^{-}(\pi, 0)=\zeta_{1} \circ v_{D}$, where $\zeta_{1}$ is a quadratic character of $F^{\times}$, while $v_{D}$ is a reduced norm on $D^{\times}$(here $D$ is a nonsplit quaternion algebra over $F$ ). We have $J L\left(\Theta^{-}(\pi, 0)\right)=\zeta_{1} \mathrm{St}_{\mathrm{GL}(2, F)}$. If $\zeta_{1}=\zeta$, then the reducibility point is $s=1 / 2$, otherwise the reducibility point is $s=3 / 2$.
(b) $\zeta=1$. This case can be completely solved using [Hanzer and Muić 2009, Theorem 4.2]. We again denote by $r$ the first occurrence of nonzero lift of representation $\pi$ in the odd orthogonal + -tower and consider all the possible cases:
- If $r=0$, then $\pi$ equals $\omega_{\psi_{a}^{-}, 1}$ and the representation $\chi_{V, \psi} \nu^{s} \rtimes \omega_{\psi_{a}^{-}, 1}$ reduces for $s= \pm 3 / 2$.
- If $r=1$, the representation $\chi_{V, \psi} \nu^{s} \rtimes \pi$ reduces for $s=1 / 2$.
- If $r=2$, the representation $\chi_{V, \psi} \nu^{s} \rtimes \pi$ reduces for $s=1 / 2$.


## Acknowledgments

We thank the Erwin Schrödinger Institute in Vienna and J. Schwermer for their hospitality during the workshop "Representation theory of reductive groups local and global aspects", where work on this paper began. We also thank M. Tadić for showing us his proof of the unitarizability of the ends of complementary series in [Tadić 1986].

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Received July 16, 2009. Revised November 10, 2009.
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The subscription price for 2010 is US $\$ 420 /$ year for the electronic version, and $\$ 485 /$ year for print and electronic.
Subscriptions, requests for back issues from the last three years and changes of subscribers address should be sent to Pacific Journal of Mathematics, P.O. Box 4163, Berkeley, CA 94704-0163, U.S.A. Prior back issues are obtainable from Periodicals Service Company, 11 Main Street, Germantown, NY 12526-5635. The Pacific Journal of Mathematics is indexed by Mathematical Reviews, Zentralblatt MATH, PASCAL CNRS Index, Referativnyi Zhurnal, Current Mathematical Publications and the Science Citation Index.
The Pacific Journal of Mathematics (ISSN 0030-8730) at the University of California, c/o Department of Mathematics, 969 Evans Hall, Berkeley, CA 94720-3840, is published monthly except July and August. Periodical rate postage paid at Berkeley, CA 94704, and additional mailing offices. POSTMASTER: send address changes to Pacific Journal of Mathematics, P.O. Box 4163, Berkeley, CA 94704-0163.

PJM peer review and production are managed by EditFLOw ${ }^{\text {TM }}$ from Mathematical Sciences Publishers.
PUBLISHED BY PACIFIC JOURNAL OF MATHEMATICS
at the University of California, Berkeley 94720-3840
A NON-PROFIT CORPORATION
Typeset in $\mathrm{IAT}_{\mathrm{E}} \mathrm{X}$
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[^0]:    MSC2000: primary 22E35, 22E50; secondary 11F70.
    Keywords: metaplectic groups, unitary dual, theta correspondence.

