## Pacific

## Journal of

 Mathematics
# REFINED OPEN NONCOMMUTATIVE DONALDSON-THOMAS INVARIANTS FOR SMALL CREPANT RESOLUTIONS 

Kentaro Nagao

## REFINED OPEN NONCOMMUTATIVE DONALDSON-THOMAS INVARIANTS FOR SMALL CREPANT RESOLUTIONS

Kentaro Nagao


#### Abstract

We study analogs of noncommutative Donaldson-Thomas invariants corresponding to the refined topological vertex for small crepant resolutions of toric Calabi-Yau 3 -folds. We give three definitions of the invariants which are equivalent to each others and provide "wall-crossing" formulas for the invariants. In particular, we get normalized generating functions which are unchanged under wall-crossing.


## Introduction

Donaldson-Thomas theory [Thomas 2000] is intersection theory on the moduli spaces of ideal sheaves on a smooth variety, which is conjecturally equivalent to Gromov-Witten theory [Maulik et al. 2006]. For a Calabi-Yau 3-fold, the virtual dimension of the moduli space is zero and hence Donaldson-Thomas invariants are said to be counting invariants of ideal sheaves. It is known that they coincide with the weighted Euler characteristics of the moduli spaces weighted by the Behrend functions [2009]. Recently, the Donaldson-Thomas invariants of Calabi-Yau 3folds have been studied using categorical methods; see, for example, [Joyce 2008; 2007; Toda 2009; 2010; Kontsevich and Soibelman 2008; Joyce and Song 2010].

On the other hand, a smooth variety $Y$ sometimes has a noncommutative associative algebra $A$ such that the derived category of coherent sheaves on $Y$ is equivalent to the derived category of $A$-modules. Derived McKay correspondence [Kapranov and Vasserot 2000; Bridgeland et al. 2001] and Van den Bergh's noncommutative crepant resolutions [2004] are typical examples. In such cases, B. Szendrői proposed to study counting invariants of $A$-modules (noncommutative DonaldsonThomas invariants) and relations with the original Donaldson-Thomas invariants on $Y$ [Szendrői 2008]. In [Nagao and Nakajima 2011; Nagao 2011a], we provided wall-crossing formulas which relate generating functions of the DonaldsonThomas and noncommutative Donaldson-Thomas invariants for small crepant resolutions of toric Calabi-Yau 3-folds. (We say a resolution of a 3-fold is small if the dimension of each fiber is less than or equal to 1.)

[^0]The aim of this paper is to propose new invariants generalizing noncommutative Donaldson-Thomas invariants and to provide "wall-crossing formulas" for small crepant resolutions of toric Calabi-Yau 3-folds. We have two directions of generalizations:

- "open" version: ${ }^{1}$ corresponding to counting invariants of sheaves on $Y$ with noncompact supports, ${ }^{2}$
- refined version: corresponding to refined topological vertex [Iqbal et al. 2009]. ${ }^{3}$

Let $Y \rightarrow X$ be a projective small crepant resolution of an affine toric Calabi-Yau 3-fold. Recall that giving an affine toric Calabi-Yau 3-fold is equivalent to giving a convex lattice polygon. Existence of a small crepant resolution is equivalent to absence of interior lattice points in the polygon. It is easy to classify such polygons and $X$ is one of the following:

- $X=X_{L^{+}, L^{-}}:=\left\{\mathrm{xy}=\mathrm{z}^{L^{+}} \mathrm{w}^{L^{-}}\right\} \subset \mathbb{C}^{4}$ for $L^{+}>0$ and $L^{-} \geq 0$, or
- $X=X_{(\mathbb{Z} / 2 \mathbb{Z})^{2}}:=\mathbb{C}^{3} /(\mathbb{Z} / 2 \mathbb{Z})^{2}$ where $(\mathbb{Z} / 2 \mathbb{Z})^{2}$ acts on $\mathbb{C}^{3}$ with weights $(1,0)$, $(0,1)$ and $(1,1)$.


Figure 1. Polygons for $X_{L^{+}, L^{-}}$and $X_{(\mathbb{Z} / 2 \mathbb{Z})^{2}}$.
In this paper, we study the first case. We put $L:=L^{+}+L^{-}$. Note that $X_{1,1}$ is called the conifold and $X_{L, 0}$ is isomorphic to $\mathbb{C} \times \mathbb{C}^{2} /(\mathbb{Z} / L \mathbb{Z})$.

Given a pair of Young diagrams $v=\left(v_{+}, v_{-}\right)$and an $L$-tuple of Young diagrams

$$
\lambda=\left(\lambda^{(1 / 2)}, \ldots, \lambda^{(L-1 / 2)}\right),
$$

the generating function of refined open noncommutative Donaldson-Thomas invariants (roncDT, in short)

$$
\mathscr{L}_{\lambda, v}^{Y}(\vec{q})=\mathscr{L}_{\lambda, v}^{Y}\left(q_{+}, q_{-}, q_{1} \ldots, q_{L-1}\right),
$$

which is denoted by $\mathscr{F}_{\sigma, \lambda, v} \mathrm{RTV}$ in the body of this paper, is defined by counting the number of the following data:

- an $(L-1)$-tuple of Young diagrams $\vec{v}=\left(v^{(1)}, \ldots, v^{(L-1)}\right)$, and

[^1]- an $L$-tuple of 3-dimensional Young diagrams $\vec{\Lambda}=\left(\Lambda^{(1 / 2)}, \ldots, \Lambda^{(L-1 / 2)}\right)$ such that $\Lambda^{(j)}$ is of type $\left(\lambda^{(j)}, \nu^{(j+1 / 2)},{ }^{\mathrm{t}} \nu^{(j-1 / 2)}\right)$ or $\left(\lambda^{(j)},{ }^{\mathrm{t}} v^{(j-1 / 2)}, \nu^{(j+1 / 2)}\right)$ (see Section 5.3 for details).

Such data parametrize torus fixed ideal sheaves on the small crepant resolution $Y$. In particular,

$$
\left.\mathscr{E}_{\varnothing, \varnothing}^{Y}(\vec{q})\right|_{q_{+}=q_{-}}
$$

coincides with the generating function of Euler characteristic versions of the Don-aldson-Thomas invariants of $Y$. ${ }^{4}$

Let $A$ be a noncommutative crepant resolution of $X$. Let $\mathbb{Z}_{\mathrm{h}}$ denote the set of half integers and let $\theta: \mathbb{Z}_{\mathrm{h}} \rightarrow \mathbb{Z}_{\mathrm{h}}$ be a bijection such that $\theta(h+L)=\theta(h)+L$ and such that

$$
\theta(1 / 2)+\cdots+\theta(L-1 / 2)=1 / 2+\cdots+(L-1 / 2) .
$$

We will define generating functions $\mathscr{E}_{\lambda, v, \theta}^{A}(\vec{q})$, which are denoted by $\mathscr{I}_{\sigma, \lambda, v, \theta}(\vec{q})$ in the body of this paper (see Section 3.4), satisfying these properties:

- $\left.\mathscr{E}_{\varnothing, \varnothing, \mathrm{id}}^{A}(\vec{q})\right|_{q_{+}=q_{-}=q_{0}^{1 / 2}}$ coincides with the generating function $\mathscr{I}_{\mathrm{NCDT}, \mathrm{eu}}^{A}$ of Euler characteristic versions ${ }^{5}$ of noncommutative Donaldson-Thomas invariants for the noncommutative crepant resolution $A$; see [Mozgovoy and Reineke 2010] and the remark on page 184.
- " $\lim _{\theta \rightarrow \infty} \mathscr{E}_{\lambda, v, \theta}^{A}(\vec{q})=\mathscr{I}_{\lambda, v}^{Y}(\vec{q})$; see Theorem 5.4.8. (The limit in this equation is, in fact, equivalent to a limit in the space of stability conditions for the category of finite-dimensional $A$-modules. $)^{6}$

Moreover, for $i \in I:=\mathbb{Z} / L \mathbb{Z}$ we can define the new bijection $\mu_{i}(\theta): \mathbb{Z}_{\mathrm{h}} \rightarrow \mathbb{Z}_{\mathrm{h}}$ (see §1.2.1) and

- $\mathscr{L}_{\lambda, v, \mu_{i}(\theta)}^{A}(\vec{q}) / \mathscr{E}_{\lambda, v, \theta}^{A}(\vec{q})$ is given explicitly (Theorem 4.2.2 and 4.4.2).

In [Nagao and Nakajima 2011; Nagao 2011a], we realized the $\left.\mathscr{E}_{\varnothing, \varnothing, \theta}^{A}(\vec{q})\right|_{q_{+}=q_{-}}$as generating functions of virtual counting of certain moduli spaces and these moduli spaces are constructed using geometric invariant theory. In this story, $\theta$ determines a chamber in the space of stability parameters and the chamber corresponding to $\theta$ is adjacent to the chamber corresponding to $\mu_{i}(\theta)$ by a single wall. This is the reason we call Theorem 4.2.2 and 4.4.2 as wall-crossing formulas, even though our

[^2]definition of the invariants and the proof of the formula are given in combinatorial ways. In fact, in the subsequent paper [Nagao 2011b] we provide an alternative geometric definition, in which $\theta$ determines a chamber in the space of Bridgeland's stability conditions for the category of finite-dimensional $A$-modules.

As consequences of the wall-crossing formula, we get

- Corollaries 4.5.2 and 5.5.2: $\mathscr{Z}_{\lambda, v, \theta}^{A} / \mathscr{L}_{\lambda, \varnothing, \theta}^{A}=\mathscr{L}_{\lambda, \nu}^{Y} / \mathscr{L}_{\lambda, \varnothing}^{Y}$ for any $\theta, \lambda$ and $\nu$.
- Corollaries 4.5.4 and 5.5.4: $\left.\quad\left(\mathscr{L}_{\lambda, v, \theta}^{A} / \mathscr{L}_{\varnothing, \varnothing, \theta}^{A}\right)\right|_{q_{+}=q_{-}}=\left.\left(\mathscr{L}_{\lambda, v}^{Y} / \mathscr{L}_{\varnothing, \varnothing}^{Y}\right)\right|_{q_{+}=q_{-}}$ for any $\theta, \lambda$ and $v$ such that $c_{\lambda}[j]=0$ for any $j$ (see $\S 1.3 .1$ for notation).

By the results in [Nagao and Nakajima 2011; Nagao 2011a], these formulas should be interpreted as stability of the normalized generating functions under wall crossing. We can find such stability of normalized generating functions in other contexts such as flop invariance and DT-PT correspondence. Categorical interpretations of such normalized generating functions and their stability are expected.

Now, we summarize the prior study on noncommutative Donaldson-Thomas invariants:

- Szendrői's formula on the generating function of noncommutative DonaldsonThomas invariants of the conifold was shown by B. Young [2009] in a purely combinatorial way. The main tool is an operation called dimer shuffling.
- J. Brian and Young [2010] generalized the Szendrői-Young formula for $X_{L, 0}$ and $X_{(\mathbb{Z} / 2 \mathbb{Z})^{2}}$. The method is different from the one used in [Young 2009]: they use vertex operator method.
- In [Nagao and Nakajima 2011], we gave an interpretation of Szendrői-Young formula as a consequence of the wall-crossing formula. From our point of view, the argument there can be translated into combinatorial language by localization, yielding the argument in [Young 2009]. In particular, dimer shuffling is nothing but "mutation" in the categorical language.
- In [Nagao 2011a], we generalized the results in [Nagao and Nakajima 2011] for arbitrary small crepant resolutions of toric Calabi-Yau 3-folds.
- In [Joyce and Song 2010], the authors study noncommutative DonaldsonThomas invariants of small crepant resolutions of toric Calabi-Yau 3-folds as examples of their theory of generalized Donaldson-Thomas invariants.
- T. Dimofte and S. Gukov [2010] provided a refined version of SzendrőiYoung formula for the conifold.
- See [Jafferis and Moore 2008; Chuang and Jafferis 2009; Aganagic et al. 2011; Chuang and Pan 2010; Aganagic and Yamazaki 2010; Dimofte et al. 2011] for developments in physics.

In this paper, we define the roncDT invariants using a dimer model (Section 2), which is purely combinatorial.

In Section 3, we give an interpretation of the dimer model as a crystal melting model. ${ }^{7}$ We construct an $A$-module $M_{\sigma, \lambda, v, \theta}^{\max }$ such that giving a dimer configuration is equivalent to giving a finite-dimensional torus invariant quotient module of $M_{\sigma, \lambda, v, \theta}^{\max }$. Hence the roncDT invariant coincides with the Euler characteristic of the moduli space of finite-dimensional quotient modules of $M_{\sigma, \lambda, v, \theta}^{\max }$; see [Nagao 2011b]. ${ }^{8}$

In Section 4, we introduce the notion of dimer shuffling to prove the first main result of this paper: the wall-crossing formula (Theorems 4.2.2 and 4.4.2).

Finally we study the limit behavior of the dimer model in Section 5. The second main result is that the generating function given by the refined topological vertex for $Y$ appears as the limit (Theorem 5.4.8).

While preparing the papers, the author was informed from J. Bryan that he and his collaborators C. Cadman and B. Young provided an explicit formula of $\mathscr{I}_{\lambda, v, \text { id }}^{A} \mid q_{+=q_{-}}$for $X_{L, 0}$ and $X_{(\mathbb{Z} / 2 \mathbb{Z})^{2}}$ using vertex operator methods [Bryan et al. 2012; $\geq 2011$ ]. In a subsequent paper [Nagao 2011b], we provide an explicit formula of $\mathscr{E}_{\lambda, v, \theta}^{A}$ for $X_{L_{+}, L_{-}}$using vertex operator methods.

A physicist may refer to [Nagao and Yamazaki 2010], in which we explain the result of this paper in a physical context.

We conclude this introduction by definition some notation.
Indices. Let $\mathbb{Z}_{\mathrm{h}}$ denote the set of half integers and $L$ be a positive integer. We set $I:=\mathbb{Z} / L \mathbb{Z}$ and $I_{\mathrm{h}}:=\mathbb{Z}_{\mathrm{h}} / L \mathbb{Z}$. The two natural projections $\mathbb{Z} \rightarrow I$ and $\mathbb{Z}_{\mathrm{h}} \rightarrow I_{\mathrm{h}}$ are denoted by the same symbol $\pi$. We sometimes identify $I$ and $I_{\mathrm{h}}$ with $\{0, \ldots, L-1\}$ and $\{1 / 2, \ldots, L-1 / 2\}$ respectively.

The symbols $n, h, i$ and $j$ are used for elements in $\mathbb{Z}, \mathbb{Z}_{\mathrm{h}}, I$ and $I_{\mathrm{h}}$ respectively.
For $n \in \mathbb{Z}$ and $h \in \mathbb{Z}_{\mathrm{h}}$, we define $c(n), c(h) \in \mathbb{Z}$ by

$$
n=c(n) \cdot L+\pi(n), \quad h=c(h) \cdot L+\pi(h) .
$$

Young diagrams. A Young diagram $v$ is a map $v: \mathbb{Z} \rightarrow \mathbb{Z}$ such that $v(n)=|n|$ if $|n| \gg 0$ and $v(n)-v(n-1)= \pm 1$ for any $n \in \mathbb{Z}$. The map $\mathbb{Z}_{\mathrm{h}} \rightarrow\{ \pm 1\}$ given by $j \mapsto \nu(j+1 / 2)-v(j-1 / 2)$ is also denoted by $\nu$.

By an abuse of notation, we sometimes identify + and - with 1 and -1 .

[^3]A Young diagram can be represented by a nonincreasing sequence of positive integers. We fix the notation as in Figure 2.


Figure 2. $v=(1,1),{ }^{t} v=(2)$.

Formal variables. Let $q_{+}, q_{-}$and $q_{0}, \ldots, q_{L-1}$ be formal variables. We use $q_{+}, q_{-}$ and $q_{1}, \ldots, q_{L-1}$ for generating functions of refined invariants. Substituting $q_{+}=$ $q_{-}=\left(q_{0}\right)^{1 / 2}$, we get generating functions of nonrefined invariants.

Let $P:=\mathbb{Z} \cdot I$ be the lattice with the basis $\left\{\alpha_{i} \mid i \in I\right\}$. For an element $\alpha=$ $\sum \alpha^{i} \cdot \alpha_{i} \in P \quad\left(\alpha^{i} \in \mathbb{Z}\right)$, we put $q^{\alpha}:=\prod\left(q_{i}\right)^{\alpha^{i}}$.

For $\alpha, \alpha^{\prime} \in P$, we say $\alpha<\alpha^{\prime}$ or $q^{\alpha}<q^{\alpha^{\prime}}$ if $\alpha^{\prime}-\alpha \in P^{+}:=\mathbb{Z}_{\geq 0} \cdot I$.

## 1. Preliminaries

### 1.1. Affine root system.

1.1.1. For $h, h^{\prime} \in \mathbb{Z}_{\mathrm{h}}$, we define $\alpha_{\left[h, h^{\prime}\right]} \in P$ by

$$
\alpha_{\left[h, h^{\prime}\right]}:=\sum_{n=h+1 / 2}^{h^{\prime}-1 / 2} \alpha_{\pi(n)}
$$

if $h<h^{\prime}, \alpha_{\left[h, h^{\prime}\right]}=1$ if $h=h^{\prime}$ and $\alpha_{\left[h, h^{\prime}\right]}=-\alpha_{\left[h^{\prime}, h\right]}$ if $h>h^{\prime}$. We set

$$
\begin{aligned}
\Lambda & :=\left\{\alpha_{\left[h, h^{\prime}\right]} \in P \mid h \neq h^{\prime}\right\}, \\
\Lambda^{\mathrm{re},+} & :=\left\{\alpha_{\left[h, h^{\prime}\right]} \in \Lambda \mid h<h^{\prime}, h \not \equiv h^{\prime}(\bmod L)\right\} .
\end{aligned}
$$

An element in $\Lambda\left(\right.$ resp. $\left.\Lambda^{\mathrm{re},+}\right)$ is called a root (resp. positive real root) of the affine root system of type $A_{L-1}$.
1.1.2. The element $\delta:=\alpha_{0}+\cdots \alpha_{L-1} \in P$ is called the minimal imaginary root. We set

$$
\Lambda^{\mathrm{fin},+}:=\left\{\alpha_{\left[j, j^{\prime}\right]} \in \Lambda \mid 1 / 2 \leq j<j^{\prime} \leq L-1 / 2\right\}
$$

and

$$
\begin{equation*}
\Lambda_{+}^{\mathrm{re},+}:=\left\{\alpha_{\left[j, j^{\prime}\right]}+N \delta \mid \alpha_{\left[j, j^{\prime}\right]} \in \Lambda^{\mathrm{fin},+}, N \geq 0\right\} . \tag{1-1}
\end{equation*}
$$

Example 1.1.3. In the case of $L=4$, we have

$$
\Lambda^{\mathrm{fin},+}:=\left\{\alpha_{1}, \alpha_{2}, \alpha_{3}, \alpha_{2}+\alpha_{3}, \alpha_{1}+\alpha_{2}, \alpha_{1}+\alpha_{2}+\alpha_{3}\right\}
$$

1.1.4. For a root $\alpha \in \Lambda$, we take $h$ and $h^{\prime}$ such that $\alpha=\alpha_{\left[h, h^{\prime}\right]}$ and set

$$
j_{-}(\alpha):=\pi(h), \text { and } j_{+}(\alpha):=\pi\left(h^{\prime}\right) .
$$

We also put

$$
B^{\alpha}:=\left\{\left(h, h^{\prime}\right) \in\left(\mathbb{Z}_{\mathrm{h}}\right)^{2} \mid \alpha_{\left[h, h^{\prime}\right]}=\alpha\right\} .
$$

1.1.5. Let $\Theta$ denote the set of bijections $\theta: \mathbb{Z}_{\mathrm{h}} \rightarrow \mathbb{Z}_{\mathrm{h}}$ such that

- $\theta(h+L)=\theta(h)+L$ for any $h \in \mathbb{Z}_{\mathbf{h}}$, and
- $\sum_{h=1 / 2}^{L-1 / 2} \theta(h)=\sum_{h=1 / 2}^{L-1 / 2} h$.

Example 1.1.6. In the case of $L=4$, the correspondence

$$
\frac{1}{2} \mapsto-\frac{1}{2}, \quad \frac{3}{2} \mapsto \frac{3}{2}, \quad \frac{5}{2} \mapsto \frac{5}{2}, \quad \frac{7}{2} \mapsto \frac{9}{2}
$$

gives an elements in $\Theta$. Let $\mu_{0}(\mathrm{id})$ denote this map (see $\S 1.2 .1$ for notation).
1.1.7. For $\theta \in \Theta$ and $i \in I$, we define $\alpha(\theta, i):=\alpha_{[\theta(n-1 / 2), \theta(n+1 / 2)]}\left(n \in \pi^{-1}(i)\right)$.

## Example 1.1.8.

$$
\begin{array}{ll}
\alpha(\mathrm{id}, 0)=\alpha_{0}, & \alpha\left(\mu_{0}(\mathrm{id}), 0\right)=-\alpha_{0}, \\
\alpha(\mathrm{id}, 1)=\alpha_{1}, & \alpha\left(\mu_{0}(\mathrm{id}), 1\right)=\alpha_{0}+\alpha_{1}, \\
\alpha(\mathrm{id}, 2)=\alpha_{2}, & \alpha\left(\mu_{0}(\mathrm{id}), 2\right)=\alpha_{2}, \\
\alpha(\mathrm{id}, 3)=\alpha_{3}, & \alpha\left(\mu_{0}(\mathrm{id}), 3\right)=\alpha_{0}+\alpha_{3} .
\end{array}
$$

1.1.9. If $\alpha=\alpha_{\left[h, h^{\prime}\right]}$ is a positive real root, we write $\theta(\alpha)>0$ if $\theta^{-1}(h)>\theta^{-1}\left(h^{\prime}\right)$, and we write $\theta(\alpha)<0$ if $\theta^{-1}(h)<\theta^{-1}\left(h^{\prime}\right)$. We set

$$
\begin{equation*}
\Lambda_{\theta}^{\mathrm{re},+}:=\left\{\alpha \in \Lambda^{\mathrm{re},+} \mid \theta(\alpha)>0\right\} . \tag{1-2}
\end{equation*}
$$

Example 1.1.10. We have $\Lambda_{\mathrm{id}}^{\mathrm{re},+}=\varnothing$ and $\Lambda_{\mu_{0}(\mathrm{id})}^{\mathrm{re},+}=\left\{\alpha_{0}\right\}$.
Remark. As we mentioned in the introduction, we studied moduli spaces of representations of a noncommutative crepant resolution of $X_{L^{+}, L^{-}}$in [Nagao 2011a]. In this theory, the space of stability conditions can be canonically identified with $P^{*} \otimes \mathbb{R}$ and the walls are classified as follows:

- the walls $W_{\alpha}:=(\mathbb{R} \cdot \alpha)^{\perp} \subset P^{*} \otimes \mathbb{R} \quad\left(\alpha \in \Lambda^{\mathrm{re},+}\right)$, and
- the wall $W_{\delta}:=(\mathbb{R} \cdot \delta)^{\perp}$, which separates the Donaldson-Thomas and Pandhari-pande-Thomas chambers.
The maps $\theta: \mathbb{Z}_{\mathrm{h}} \rightarrow \mathbb{Z}_{\mathrm{h}}$ as above parametrize the chambers on one side of the wall $W_{\delta}$. The notation $\theta(\alpha) \gtrless 0$ respects this parametrization.


### 1.2. Wall-crossing.

1.2.1. For $i \in I$, let $\mu_{i}: \mathbb{Z}_{\mathrm{h}} \rightarrow \mathbb{Z}_{\mathrm{h}}$ be the map given by

$$
\mu_{i}(h)= \begin{cases}h-1 & \text { if } \pi(h-1 / 2)=i, \\ h+1 & \text { if } \pi(h+1 / 2)=i, \\ h & \text { otherwise }\end{cases}
$$

For $\theta \in \Theta$, we put $\mu_{i}(\theta):=\theta \circ \mu_{i}$.
Remark. The chambers corresponding to $\theta$ and $\mu_{i}(\theta)$ are separated by the wall $W_{\alpha(\theta, i)}$, which is the reason for the title of this subsection. From the viewpoint of the affine root system, wall crossing corresponds to simple reflection; from the viewpoint of noncommutative crepant resolutions, it corresponds to mutation; and from the viewpoint of dimer models, to dimer shuffling.
1.2.2. Let $\boldsymbol{i}=\left(i_{1}, i_{2}, \ldots\right) \in I^{\mathbb{Z}}>0$ be a sequence of elements in $I$. For $b>0$, we define

$$
\theta_{i, b}:=\mu_{i_{b-1}}\left(\cdots\left(\mu_{i_{1}}(\mathrm{id})\right) \cdots\right) \in \Theta, \quad \alpha_{i, b}:=\alpha\left(\theta_{i, b}, i_{b}\right) .
$$

We say $\boldsymbol{i} \in I^{\mathbb{Z}}>0$ is a minimal expression if $\theta_{i, b}\left(\alpha_{i, b}\right)<0$ for any $b>0$. For a minimal expression $\boldsymbol{i}$, we have

$$
\Lambda_{\theta_{i, b}}^{\mathrm{re},+}=\left\{\alpha_{i, 1}, \ldots, \alpha_{i, b-1}\right\} .
$$

### 1.3. Core and quotient of a Young diagram.

1.3.1. Let $\sigma: I_{\mathrm{h}} \rightarrow\{ \pm\}$ and $\lambda: \mathbb{Z}_{\mathrm{h}} \rightarrow\{ \pm\}$ be maps such that $\lambda(h)= \pm \sigma(\pi(h))$ if $\pm h \gg 0$. We define integers $c_{\lambda}[j]$ and Young diagrams $\lambda^{[j]}$ for $j \in I_{\mathrm{h}}$ by

$$
\lambda(h)=\lambda^{[\pi(h)]}\left(\sigma(j(h)) \cdot\left(c(h)-c_{\lambda}[\pi(h)]+1 / 2\right)\right) .
$$

Remark. In the case $\sigma \equiv+$ and $\sum c_{\lambda}[j]=0$, the sequence $\left(c_{\lambda}[j]\right)$ of integers and the sequence $\left(\lambda^{[j]}\right)$ of Young diagrams are called the $L$-core and the $L$-quotient of the Young diagram $\lambda$.
1.3.2. We put

$$
\begin{equation*}
B_{\sigma, \lambda}^{\alpha, \pm}:=\left\{\left(h, h^{\prime}\right) \in B^{\alpha} \mid-\lambda(h) \sigma(h)=\lambda\left(h^{\prime}\right) \sigma\left(h^{\prime}\right)= \pm\right\} . \tag{1-3}
\end{equation*}
$$

Lemma 1.3.3.

$$
\left|B_{\sigma, \lambda}^{\alpha,+}\right|-\left|B_{\sigma, \lambda}^{\alpha,-}\right|=\alpha^{0}+c_{\lambda}\left[j_{-}(\alpha)\right]-c_{\lambda}\left[j_{+}(\alpha)\right] .
$$

Proof. We write simply $j_{ \pm}$for $j_{ \pm}(\alpha)$. Note that we have

$$
B^{\alpha}=\left\{\left(c L+j_{-},\left(c+\alpha^{0}\right) L+j_{+}\right) \mid c \in \mathbb{Z}\right\} .
$$

For an integer $N$, we set

$$
B_{N}^{\alpha}:=\left\{\left(c L+j_{-},\left(c+\alpha^{0}\right) L+j_{+}\right) \mid c \in[-N, N-1]\right\} .
$$

Take a sufficiently large integer $N$. Then

$$
B_{\sigma, \lambda}^{\alpha,+}, B_{\sigma, \lambda}^{\alpha,-} \subset B_{N}^{\alpha}
$$

and so

$$
\begin{aligned}
\left|B_{\sigma, \lambda}^{\alpha,+}\right|- & \left|B_{\sigma, \lambda}^{\alpha,-}\right| \\
= & -\sharp\left\{\left(h, h^{\prime}\right) \in B_{N}^{\alpha} \mid \lambda(h) \sigma(h)=+\right\}+\sharp\left\{\left(h, h^{\prime}\right) \in B_{N}^{\alpha} \mid \lambda\left(h^{\prime}\right) \sigma\left(h^{\prime}\right)=+\right\} \\
= & -\sharp\left\{c \in[-N, N-1] \mid \lambda^{\left[j_{-}\right]}\left(\sigma\left(j_{-}\right) \cdot\left(c-c_{\lambda}\left[j_{-}\right]+1 / 2\right)\right)=\sigma\left(j_{-}\right)\right\} \\
& \quad+\sharp\left\{c \in[-N, N-1] \mid \lambda^{\left[j_{+}\right]}\left(\sigma\left(j_{+}\right) \cdot\left(c+\alpha^{0}-c_{\lambda}\left[j_{+}\right]+1 / 2\right)\right)=\sigma\left(j_{+}\right)\right\} \\
= & -\left(N-c_{\lambda}\left[j_{-}\right]-1 / 2\right)+\left(N+\alpha^{0}-c_{\lambda}\left[j_{+}\right]-1 / 2\right) \\
= & \alpha^{0}+c_{\lambda}\left[j_{-}\right]-c_{\lambda}\left[j_{+}\right] .
\end{aligned}
$$

For $\sigma, \lambda, \theta$ and $i$, we put

$$
\begin{equation*}
B_{\sigma, \lambda, \theta}^{i, \pm}:=\left\{n \in \pi^{-1}(i) \mid(\theta(n-1 / 2), \theta(n+1 / 2)) \in B_{\sigma, \lambda}^{\alpha(\theta, i), \pm}\right\} \tag{1-4}
\end{equation*}
$$

## 2. Dimer model

### 2.1. Dimer configurations.

2.1.1. We fix the following data:

- a map $\sigma: I_{\mathrm{h}} \rightarrow\{ \pm\}$,
- a map $\lambda: \mathbb{Z}_{\mathrm{h}} \rightarrow\{ \pm\}$ such that $\lambda(h)= \pm \sigma(\pi(h))$ for $\pm h \gg 0$,
- a pair of Young diagrams $v=\left(v_{+}, v_{-}\right)$,
- a bijection $\theta: \mathbb{Z}_{\mathrm{h}} \rightarrow \mathbb{Z}_{\mathrm{h}}$ in $\Theta$.

We put $\tilde{\sigma}:=\sigma \circ \pi \circ \theta, \tilde{\lambda}:=\lambda \circ \theta$ and $L_{ \pm}:=\left|\sigma^{-1}( \pm)\right|$.
2.1.2. We consider the following graph in the $(x, y)$-plane. First, we set
(2-1) $H(\sigma, \theta):=\{n \in \mathbb{Z} \mid \tilde{\sigma}(n-1 / 2)=\tilde{\sigma}(n+1 / 2)\}, \quad I_{H}(\sigma, \theta):=\pi(H(\sigma, \theta))$,
$(2-2) \quad S(\sigma, \theta):=\{n \in \mathbb{Z} \mid \tilde{\sigma}(n-1 / 2) \neq \tilde{\sigma}(n+1 / 2)\}, \quad I_{S}(\sigma, \theta):=\pi(S(\sigma, \theta))$ and for $n \in H(\sigma, \theta)$ we put $\tilde{\sigma}(n):=\tilde{\sigma}(n \pm 1 / 2)$.

The set of the vertices is given by

$$
\begin{aligned}
\mathscr{V}:= & \{(n, m) \mid n \in S(\sigma, \theta), n-m: \text { odd }\} \\
& \sqcup\{(n-1 / 2, m) \mid n \in H(\sigma, \theta), n-m: \text { odd }\} \\
& \sqcup\{(n+1 / 2, m) \mid n \in H(\sigma, \theta), n-m: \text { odd }\},
\end{aligned}
$$

which are denoted by $v(n, m), v_{1}(n-1 / 2, m)$ and $v_{\mathrm{r}}(n+1 / 2, m)$ respectively.

The set of the edges is given by

$$
\mathscr{E}:=\left\{e_{\mathrm{h}}(n, m) \mid n \in H(\sigma, \theta), n-m: \text { odd }\right\} \sqcup\left\{e_{\mathrm{s}}(h, k) \mid h, k \in \mathbb{Z}_{\mathrm{h}}\right\},
$$

where

- $e_{\mathrm{h}}(n, m)$ connects $v_{1}(n-1 / 2, m)$ and $v_{\mathrm{r}}(n+1 / 2, m)$,
- $e_{\mathrm{s}}(h, k)$ connects $v(h-1 / 2, k+1 / 2)$ or $v_{\mathrm{r}}(h, k+1 / 2)$ and $v(h+1 / 2, k-1 / 2)$ or $v_{1}(h, k-1 / 2)$ if $h-k$ is even, and
- $e_{\mathrm{s}}(h, k)$ connects $v(h-1 / 2, k-1 / 2)$ or $v_{\mathrm{r}}(h, k-1 / 2)$ and $v(h+1 / 2, k+1 / 2)$ or $v_{1}(h, k+1 / 2)$ if $h-k$ is odd.

We put

$$
\begin{equation*}
\mathscr{F}:=\left\{(n, m) \in \mathbb{Z}^{2} \mid n+m: \text { even }\right\}, \quad \mathscr{F}_{i}:=\left\{(n, m) \in \mathscr{F} \mid n \in \pi^{-1}(i)\right\} \tag{2-3}
\end{equation*}
$$

for $i \in I$. Note that $\mathscr{E}$ divides the plain into disjoint hexagons and quadrilaterals. The hexagons are parametrized by the set

$$
\mathscr{F}_{H}:=\{(n, m) \in \mathscr{F} \mid n \in H(\sigma, \theta)\}
$$

and the quadrilaterals are parametrized by the set

$$
\mathscr{F}_{\mathrm{S}}:=\{(n, m) \in \mathscr{F} \mid n \in S(\sigma, \theta)\} .
$$

For $(n, m) \in \mathscr{F}$, let $f(n, m)$ denote the corresponding hexagon or quadrilateral.
Example 2.1.3. In Figure 3, we show the graph associated with $L=3, \sigma$ given by

$$
\sigma(1 / 2)=+, \quad \sigma(3 / 2)=-, \quad \sigma(5 / 2)=-,
$$

and $\theta=\mathrm{id}\left(L_{+}=1, L_{-}=2\right)$.


Figure 3. Graph and $\mathscr{V}_{+}$for Example 2.1.3.

### 2.1.4. We set

$$
\begin{aligned}
\mathscr{V}_{ \pm}:= & \{v(n, m) \mid \tilde{\sigma}(n+1 / 2)= \pm\} \\
& \sqcup\left\{v_{1}(n-1 / 2, m) \mid \tilde{\sigma}(n)=\mp\right\} \sqcup\left\{v_{\mathrm{r}}(n+1 / 2, m) \mid \tilde{\sigma}(n)= \pm\right\}
\end{aligned}
$$

Note that $\mathscr{V}=\mathscr{V}_{+} \sqcup \mathscr{V}_{-}$and each element in $\mathscr{E}$ connects an element in $\mathscr{V}_{+}$and an element in $\mathscr{V}_{-}$(see Figure 3 for example).

A perfect matching is a subset of $\mathscr{E}$ giving a bijection between $\mathscr{V}_{+}$and $\mathscr{V}_{-}$.
2.1.5. We define the map $F_{\sigma, \lambda, \theta}: \mathbb{Z} \rightarrow \mathbb{Z}$ by $F_{\sigma, \lambda, \theta}(0)=0$ and

$$
\begin{equation*}
F_{\sigma, \lambda, \theta}(n)=F_{\sigma, \lambda, \theta}(n-1)-\tilde{\lambda}(n-1 / 2) \tag{2-4}
\end{equation*}
$$

For $k \in \mathbb{Z}_{\mathrm{h}}$, we set

$$
\begin{aligned}
\mathscr{P}_{\sigma, \lambda, \theta}^{k, \pm}:= & \left\{e_{\mathrm{h}}\left(n, F_{\sigma, \lambda, \theta}(n)+2 k\right) \mid n \in \mathbb{Z}, \tilde{\sigma}(n)=\mp\right\} \\
& \sqcup\left\{\left.e_{\mathrm{s}}\left(h, \frac{1}{2}\left(F_{\sigma, \lambda, \theta}(h-1 / 2)+F_{\sigma, \lambda, \theta}(h+1 / 2)\right)+2 k\right) \right\rvert\, h \in \mathbb{Z}_{\mathrm{h}}, \tilde{\sigma}(h)= \pm\right\} .
\end{aligned}
$$

For a Young diagram $\eta$, define the perfect matching

$$
\mathscr{P}_{\sigma, \lambda, \theta}^{\eta}:=\bigsqcup_{k \in \mathbb{Z}_{\mathrm{h}}} \mathscr{P}_{\sigma, \lambda, \theta}^{k, \eta(k)}
$$

Example 2.1.6. In Figure 4, we show the perfect matching associated with $\sigma$ as in Example 2.1.3, $\theta=\mathrm{id}, \eta=\varnothing$, and $\lambda$ given by

$$
\lambda(h)= \begin{cases}+ & \text { if } h=-5 / 2 \\ - & \text { if } h=1 / 2 \\ \operatorname{sgn}(h) \sigma(h) & \text { otherwise }\end{cases}
$$



Figure 4. Example 2.1.6: $\left\{f\left(n, F_{\sigma, \lambda, \mathrm{id}}(n)\right) \mid n \in \mathbb{Z}\right\}$ and $\mathscr{P}_{\sigma, \lambda, \mathrm{id}}^{\varnothing}$.
2.1.7. Define the perfect matching

$$
\mathscr{P}_{\sigma, \lambda, \theta}^{ \pm}:=\left\{e_{\mathrm{h}}(n, m) \mid \tilde{\sigma}(n)=\mp\right\} \sqcup\left\{e_{\mathrm{s}}(h, k) \mid \tilde{\sigma}(h)= \pm, h \cdot \tilde{\lambda}(h)-k: \text { even }\right\} .
$$

Definition 2.1.8. A perfect matching $D$ is said to be a dimer configuration of type $(\sigma, \lambda, \nu, \theta)$ if $D$ coincides with $\mathscr{P}_{\sigma, \lambda, \theta}^{\nu_{ \pm}}$in the area $\{ \pm x>m\}$ and $\mathscr{P}_{\sigma, \lambda, \theta}^{ \pm}$in the area $\{ \pm y>m\}$ for $m \gg 0$.
Remark. A dimer configuration of type ( $\sigma, \vec{\varnothing}, \vec{\varnothing}, \mathrm{id}$ ) is "a perfect matching congruent to the canonical perfect matching" in the terminology of [Mozgovoy and Reineke 2010].
2.1.9. For $f \in \mathscr{F}$, let $\partial f \subset \mathscr{E}$ denote the set of edges surrounding the face $f$. By moving $f$ around clockwise, we can determine an orientation for each element in $\partial f$. Let $\partial^{ \pm} f \subset \partial f$ denote the subset of edges starting from elements in $\mathscr{V}_{ \pm}$.

For an edge $e \in \mathscr{E}$, let $f^{ \pm}(e)$ denote the unique face such that $e \in \partial^{ \pm} f^{ \pm}(e)$.

### 2.2. Weights.

2.2.1. For $h \in \mathbb{Z}_{\mathrm{h}}$, we define the monomials $w_{\sigma, \lambda}(h)$ by the conditions

$$
w_{\sigma, \lambda}(h)= \begin{cases}\left(Q_{\sigma(h)}\right)^{c(h)-c_{\lambda}[j(h)]} q_{\sigma(h)}^{(j(h))} & \text { if } h \gg 0, \\ \left(Q_{-\sigma(h)}\right)^{c(h)-c_{\lambda}[j(h)]} q_{-\sigma(h)}^{(j(h))} & \text { if } h \ll 0,\end{cases}
$$

and

$$
w_{\sigma, \lambda}(h) / w_{\sigma, \lambda}(h-L)=q_{\lambda(h)} \cdot q_{\lambda(h-L)} \cdot q_{1} \cdots \cdots q_{L-1},
$$

where

$$
Q_{ \pm}:=\left(q_{ \pm}\right)^{2} \cdot q_{1} \cdots \cdots q_{L-1}, q_{ \pm}^{(j)}:=q_{ \pm} \cdot q_{1} \cdots \cdots q_{j-1 / 2}
$$

Note that for $h \neq h^{\prime}$ we have

$$
\begin{equation*}
w_{\lambda}\left(h^{\prime}\right) /\left.w_{\lambda}(h)\right|_{q_{+}=q_{-}=\left(q_{0}\right)^{1 / 2}}=q^{\alpha_{\left[l, h^{\prime}\right]}} . \tag{2-5}
\end{equation*}
$$

Example 2.2.2. Figure 5 shows the weight $w_{\sigma, \lambda}$ for $\sigma$ and $\lambda$ as in Example 2.1.6.


Figure 5. The weight $w_{\sigma, \lambda}$.
2.2.3. To an edge $e \in \mathscr{E}$ we associate the weight $w_{\sigma, \lambda, \theta}(e)$ by

$$
\begin{align*}
& w_{\sigma, \lambda, \theta}\left(e_{\mathrm{s}}(h, k)\right):=\left\{\begin{array}{cl}
w_{\sigma, \lambda}(\theta(h))^{\tilde{\sigma}(h) \cdot \tilde{\lambda}(h)} & \text { if } h \cdot \tilde{\lambda}(h)-k \text { is odd } \\
1 & \text { if } h \cdot \tilde{\lambda}(h)-k \text { is even }
\end{array}\right.  \tag{2-6}\\
& w_{\lambda, \sigma, \theta}\left(e_{\mathrm{h}}(n, m)\right):=1 \tag{2-7}
\end{align*}
$$

2.2.4. Fix $\sigma$ and $\lambda$. Then the set $\bigsqcup_{\alpha \in \Lambda^{\mathrm{re},+}} B_{\sigma, \lambda}^{\alpha,-}$ is finite. We define

$$
\begin{equation*}
F_{\sigma, \lambda}^{\alpha}:=\prod_{\substack{\left(h, h^{\prime}\right) \in B_{\sigma, \lambda}^{\alpha,-}}} \frac{w_{\sigma, \lambda}\left(h^{\prime}\right)}{w_{\sigma, \lambda}(h)}, \quad F_{\sigma, \lambda}^{\theta}:=\prod_{\substack{\alpha \in \Lambda^{\mathrm{re},+} ; \theta(\alpha)<0, \sigma\left(j^{-}(\alpha)\right) \neq \sigma\left(j^{+}(\alpha)\right)}} F_{\sigma, \lambda}^{\alpha} \tag{2-8}
\end{equation*}
$$

2.2.5. Note that for a dimer configuration $D$ of type $(\sigma, \lambda, \nu, \theta)$ we have only a finite number of $e \in D$ such that $w_{\sigma, \lambda, \theta}(e) \neq 1$.
Definition 2.2.6. For a dimer configuration $D$ of type ( $\sigma, \lambda, v, \theta$ ), we define the weight $w_{\sigma, \lambda, \theta}(D)$ by

$$
\begin{equation*}
w_{\sigma, \lambda, \theta}(D):=F_{\sigma, \lambda}^{\theta} \cdot \prod_{e \in D} w_{\sigma, \lambda, \theta}(e) \tag{2-9}
\end{equation*}
$$

(See (2-6)-(2-8) for notation.)
Remark. We will define the generating function $\mathscr{E}_{\sigma, \lambda, v, \theta}$ by the sum of weighs of all dimer configurations of type $(\sigma, \lambda, v, \theta) .{ }^{9}$
2.2.7. For a finite subset $\mathscr{E}^{\prime} \subset \mathscr{E}$, we put

$$
w_{\sigma, \lambda, \theta}\left(\mathscr{E}^{\prime}\right):=\prod_{e \in \mathscr{E}^{\prime}} w_{\sigma, \lambda, \theta}(e)
$$

and for a face $f \in \mathscr{F}$ we put

$$
\begin{equation*}
w_{\sigma, \lambda, \theta}(f):=\frac{w_{\sigma, \lambda, \theta}\left(\partial^{-} f\right)}{w_{\sigma, \lambda, \theta}\left(\partial^{+} f\right)} \tag{2-10}
\end{equation*}
$$

For an integer $n$ we set

$$
w_{\sigma, \lambda, \theta}(n):=\frac{w_{\sigma, \lambda}(\theta(n+1 / 2))}{w_{\sigma, \lambda}(\theta(n-1 / 2))}
$$

then

$$
w_{\sigma, \lambda, \theta}(f(n, m))=w_{\sigma, \lambda, \theta}(n)
$$

for any $(n, m) \in \mathscr{F}$. By (2-5), we have

$$
\left.w_{\sigma, \lambda, \theta}(n)\right|_{q_{+}=q_{-}=\left(q_{0}\right)^{1 / 2}}=q^{\alpha(\theta, i)}
$$

[^4]
## 3. The viewpoint of noncommutative crepant resolutions

3.1. Noncommutative crepant resolutions. Let $\Gamma$ be a lattice in the $(x, y)$-plane generated by $(L, 0)$ and $(0,2)$. The graph given in $\S 2.1 .2$ is invariant under the action of $\Gamma$ and so gives a graph on the torus $\mathbb{R}^{2} / \Gamma$. This gives a quiver with a potential $A=\left(Q_{\sigma, \theta}, w_{\sigma, \theta}\right)$ as in [Nagao 2011a]. The vertices of $Q_{\sigma, \theta}$ are parametrized by $I$ and the arrows are given by

$$
\left(\bigsqcup_{j \in I_{\mathrm{h}}} h_{j}^{+}\right) \sqcup\left(\bigsqcup_{j \in I_{\mathrm{h}}} h_{j}^{-}\right) \sqcup\left(\bigsqcup_{i \in I_{H}(\sigma, \theta)} r_{i}\right)
$$

(see (2-1) for notation). Here $h_{j}^{+}$(resp. $h_{j}^{-}$) is an edge from $j-1 / 2$ to $j+1 / 2$ (resp. from $j+1 / 2$ to $j-1 / 2$ ) and $r_{i}$ is an edge from $i$ to itself. See [Nagao 2011a, §1.2] for the definition of the potential $w_{\sigma, \theta}$.
Example 3.1.1. Here is the quiver $Q_{\sigma \text {,id }}$ for $\sigma$ as in Example 2.1.6:


Remarks. - The center of $A$ is isomorphic to $R:=\mathbb{C}[\mathrm{x}, \mathrm{y}, \mathrm{z}, \mathrm{w}] /\left(\mathrm{xy}=\mathrm{z}^{L_{+}} \mathrm{w}^{L_{-}}\right)$. In [Nagao 2011a, Theorem 1.14 and 1.19], we showed that $A$ is a noncommutative crepant resolution of $X=\operatorname{Spec} R$.

- The affine 3-fold $X$ is toric. In fact,

$$
T=\operatorname{Spec} \tilde{R}:=\operatorname{Spec} \mathbb{C}\left[\mathrm{x}^{ \pm}, \mathrm{y}^{ \pm}, \mathrm{z}^{ \pm}, \mathrm{w}^{ \pm}\right] /\left(\mathrm{xy}=\mathrm{z}^{L_{+}} \mathrm{w}^{L_{-}}\right) \subset X
$$

is a 3-dimensional torus.

### 3.2. Dimer model and noncommutative crepant resolution.

3.2.1. We will construct an $A$-module $M(D)$ for a dimer configuration $D$. Let $V_{i}=V_{i}(D)(i \in I)$ be vector space with the basis

$$
\left\{b[D ; x, y, z] \mid(x, y) \in \mathscr{F}_{i}, z \in \mathbb{Z}_{\geq 0}\right\}
$$

(see (2-3) for notation). We define the map $h_{j}^{ \pm}: V_{j \mp 1 / 2} \rightarrow V_{j \pm 1 / 2}$ by setting
$h_{j}^{ \pm}(b[D ; x, y, z])= \begin{cases}b[D ; x \pm 1, y-\tilde{\sigma}(j), z] & \text { if } e_{\mathrm{s}}\left(x \pm \frac{1}{2}, y-\frac{1}{2} \tilde{\sigma}(j)\right) \notin D, \\ b[D ; x \pm 1, y-\tilde{\sigma}(j), z+1] & \text { if } e_{\mathrm{s}}\left(x \pm \frac{1}{2}, y-\frac{1}{2} \tilde{\sigma}(j)\right) \in D,\end{cases}$


Figure 6. An example of $M(D)$.
and $r_{i}: V_{i} \rightarrow V_{i}$ by

$$
r_{i}(b[D ; x, y, z])= \begin{cases}b[D ; x, y+\tilde{\sigma}(j), z] & \text { if } e_{\mathrm{h}}(x, y+\tilde{\sigma}(j) / 2) \notin D \\ b[D ; x, y+\tilde{\sigma}(j), z+1] & \text { if } e_{\mathrm{h}}(x, y+\tilde{\sigma}(j) / 2) \in D\end{cases}
$$

3.2.2. Let $\mathscr{C} \subset \mathscr{E}$ be a subset which gives a closed zigzag curve without selfintersection. By moving along the zigzag curve clockwisely, we can determine an orientation for each element in $\mathscr{C}$. Let $\mathscr{C}^{ \pm} \subset \mathscr{C}$ denote the subset of edges starting from elements in $\mathscr{V}_{ \pm}$.

Let $D$ be a dimer configuration of type ( $\sigma, \lambda, \nu, \theta$ ). A subset $\mathscr{C}$ as above is said to be a positive cycle with respect to $D$ if $\mathscr{C} \cap D=\mathscr{C}^{+}$, and it is said to be a negative cycle with respect to $D$ if $\mathscr{C}^{-}$.
3.2.3. Given a dimer configuration $D$ and a positive cycle $\mathscr{C}$ with respect to $D$, let $D_{\mathscr{C}}$ be the dimer configuration given by

$$
D_{\mathscr{C}}=\left(D \backslash \mathscr{C}^{+}\right) \cup \mathscr{C}^{-} .
$$

Then we can check the following lemma:
Lemma 3.2.4. The surjection $M(D) \rightarrow M\left(D_{6}\right)$ given by

$$
b[D ; x, y, z] \mapsto \begin{cases}0 & \text { if }(x, y) \in \mathscr{C}^{\circ} \text { and } z=0, \\ b\left[D_{\mathscr{C}} ; x, y, z-1\right] & \text { if }(x, y) \in \mathscr{C}^{\circ} \text { and } z \geq 1, \\ b\left[D_{\mathscr{C}} ; x, y, z\right] & \text { if }(x, y) \notin \mathscr{C}^{\circ},\end{cases}
$$

is a homomorphism of A-modules, where $\mathscr{C}^{\circ}$ is the interior of the closed zigzag curve. Moreover,

$$
w_{\sigma, \lambda, \theta}\left(D_{\mathscr{C}}\right)=w_{\sigma, \lambda, \theta}(D) \cdot \prod_{f \in \mathscr{母}^{\circ}} w_{\sigma, \lambda, \theta}(f) .
$$

3.3. Crystal melting interpretation. In this subsection, we show that a dimer configuration of type ( $\sigma, \lambda, \nu, \theta$ ) corresponds to a (torus invariant) quotient $A$-module of the $A$-module $M^{\max }=M_{\sigma, \lambda, v, \theta}^{\max }$. In the physicists' terminology, studying such quotient modules is called the crystal melting model (see [Ooguri and Yamazaki 2009]) and $M^{\text {max }}$ is called the grand state of the model.
3.3.1. We define a Young diagram $G_{\sigma, \lambda, \theta}: \mathbb{Z} \rightarrow \mathbb{Z}$ by the following conditions:

- $G_{\sigma, \lambda, \theta}(n)=|n|$ if $|n| \gg 0$, and
- $G_{\sigma, \lambda, \theta}(n)=G_{\sigma, \lambda, \theta}(n-1)+\tilde{\sigma}(n-1 / 2) \tilde{\lambda}(n-1 / 2)$ for any $n$.

We define a map $G_{\sigma, \lambda, \theta}: \mathscr{F} \rightarrow \mathbb{Z}$ by

$$
\begin{equation*}
G_{\sigma, \lambda, \theta}(n, m):=G(n)_{\sigma, \lambda, \theta}+2 \cdot\left|m-F_{\sigma, \lambda, \theta}(n)\right|, \tag{3-1}
\end{equation*}
$$

where $F_{\sigma, \lambda, \theta}(n)$ is given in (2-4).
Example 3.3.2. In the case of Example 2.1.6, we have

$$
\left(G_{\sigma, \lambda, \mathrm{id}}(n)\right)_{n \in \mathbb{Z}}=(\ldots, 6,5,4,3,4,3,2,1,2,3,4,5,6, \ldots)
$$

and $G_{\sigma, \lambda, \mathrm{id}}(n, m)$ is given in Figure 7.
3.3.3. We define two maps $F_{\sigma, \lambda, \theta}^{ \pm}: \mathbb{Z} \rightarrow \mathbb{Z}$ by the following conditions:

- $F_{\sigma, \lambda, \theta}^{ \pm}(n)=F_{\sigma, \lambda, \theta}(n)$ if $\pm n \gg 0$.
- $F_{\sigma, \lambda, \theta}^{ \pm}(n)=F_{\sigma, \lambda, \theta}^{ \pm}(n-1) \mp \tilde{\sigma}(n-1 / 2)$ for any $n$.

Then we define two maps $G_{\sigma, \lambda, \theta}^{v_{ \pm}, \pm}: \mathscr{F} \rightarrow \mathbb{Z}$ by

$$
\begin{equation*}
G_{\sigma, \lambda, \theta}^{v_{ \pm}, \pm}(n, m):=v_{ \pm}\left(m-F_{\sigma, \lambda, \theta}^{ \pm}(n)\right) \pm n . \tag{3-2}
\end{equation*}
$$



Figure 7. $G_{\sigma, \lambda, \mathrm{id}}(n, m)$.

Example 3.3.4. Figure 8 shows $G_{\sigma, \lambda, \text { id }}^{\varnothing,+}$ and $G_{\sigma, \lambda, \mathrm{id}}^{\square,-}$ for $\sigma$ and $\lambda$ as in Example 2.1.6.


Figure 8. $G_{\sigma, \lambda, \theta}^{\varnothing,+}$ (top) and $G_{\sigma, \lambda, \theta}^{\square,-}$ (bottom).
3.3.5. We define a map $G_{\sigma, \lambda, \theta}^{v}: \mathscr{F} \rightarrow \mathbb{Z}$ by

$$
G_{\sigma, \lambda, \theta}^{v}(n, m):=\max \left(G_{\sigma, \lambda, \theta}(n, m), G_{\sigma, \lambda, \theta}^{v_{+}++}(n, m), G_{\sigma, \lambda, \theta}^{\mu_{-,-}-}(n, m)\right) .
$$

We can verify that

$$
G_{\sigma, \lambda, \theta}^{v}\left(f^{+}(e)\right)=G_{\sigma, \lambda, \theta}^{v}\left(f^{-}(e)\right)+1 \text { or } G_{\sigma, \lambda, \theta}^{v}\left(f^{-}(e)\right)-3 .
$$

for an edge $e \in \mathscr{E}$ (see $\S 2.1 .9$ for notation). We define a perfect matching $D^{\max }=$ $D_{\sigma, \lambda, v, \theta}^{\max }$ by

$$
e \in D^{\max } \Longleftrightarrow G_{\sigma, \lambda, \theta}^{v}\left(f^{+}(e)\right)=G_{\sigma, \lambda, \theta}^{v}\left(f^{-}(e)\right)-3 .
$$

Let $M^{\max }=M_{\sigma, \lambda, v, \theta}^{\max }:=M\left(D_{\sigma, \lambda, v, \theta}^{\max }\right)$ denote the corresponding $A$-module.
Example 3.3.6. In Figure 9, we show $G_{\sigma, \lambda, \text { id }}^{\vec{\theta}}$ and $D_{\sigma, \lambda, \vec{\varnothing}, \text { id }}^{\max }$ for $\sigma$ and $\lambda$ as in Example 2.1.6.


Figure 9. $G_{\sigma, \lambda, \mathrm{id}}^{\vec{\varnothing}}$ and $D_{\sigma, \lambda, \vec{\varnothing}, \mathrm{id}}^{\max }$.

Remark. The graph of the map $m \mapsto G_{\sigma, \lambda, \theta}^{v}(n, m)$ determines a Young diagram. This is what we denote by $\mathscr{V}_{\min }(n)$ in [Nagao 2011b, §3.1].

Lemma 3.3.7. There is no positive cycle with respect to $D^{\max }$.
Proof. Assume that we have a positive cycle $\mathscr{C}$. For an edge $e \in \partial \mathscr{C}$, let $f_{\text {in }}(e)$ (resp. $\left.f_{\text {out }}(e)\right)$ be the unique face such that $e \in \partial f_{\text {in }}(e)$ and $f_{\text {in }}(e) \in \mathscr{C}^{\circ}$ (resp. $e \in \partial f_{\text {out }}(e)$ and $\left.f_{\text {out }}(e) \notin \mathscr{C}^{\circ}\right)$. Then we have

$$
\begin{equation*}
G_{\sigma, \lambda, \theta}^{v}\left(f_{\text {in }}(e)\right)>G_{\sigma, \lambda, \theta}^{v}\left(f_{\text {out }}(e)\right) . \tag{3-3}
\end{equation*}
$$

Take a face $(n, m) \in \mathscr{C}^{\circ}$. If $G_{\sigma, \lambda, \theta}^{v}(n, m)=G_{\sigma, \lambda, \theta}^{\nu, \pm}(n, m)$, then

$$
\left(n \pm n^{\prime}, F_{\sigma, \lambda, \theta}^{ \pm}\left(n \pm n^{\prime}\right)-F_{\sigma, \lambda, \theta}^{ \pm}(n)+m\right) \in \mathscr{C}^{\circ}
$$

for any $n^{\prime} \geq 0$ by (3-2) and (3-3), and this is a contradiction. On the other hands, if $G_{\sigma, \lambda, \theta}^{\mu}(n, m)=G_{\sigma, \lambda, \theta}(n, m)$ and $\pm m \mp F_{\sigma, \lambda, \theta}(n) \geq 0$, then $\left(n, m \pm m^{\prime}\right) \in \mathscr{C}^{\circ}$ for any $m^{\prime} \geq 0$ by (3-1) and (3-3), and this is also a contradiction. Hence the claim follows.
3.3.8. For a map $H: \mathscr{F} \rightarrow \mathbb{Z}_{\geq 0}$, let $V_{i}^{H} \subset V_{i}\left(D^{\max }\right)(i \in I)$ be the subspace spanned by the elements

$$
\left\{b\left[D^{\max } ; x, y, z\right] \mid(x, y) \in \mathscr{F}_{i}, z \geq H(x, y)\right\} .
$$

The following proposition gives a one-to-one correspondence between dimer configurations of type ( $\sigma, \lambda, \nu, \theta$ ) and finite-dimensional quotient modules of $M_{\sigma, \lambda, v, \theta}^{\max }$.

Proposition 3.3.9. Given a monomial $\boldsymbol{q}$, we have a natural bijection between

- the set of dimer configurations of type $(\sigma, \lambda, \nu, \theta)$ with weight $\boldsymbol{q}$, and
- the set of maps $H: \mathscr{F} \rightarrow \mathbb{Z}_{\geq 0}$ satisfying the following conditions:
- $H(f)=0$ except for only a finite number of $f \in \mathscr{F}$,
- $\left(V_{i}^{H}\right)_{i \in I}$ is stable under the action of $A$, and
$-w_{\sigma, \lambda, \theta}\left(D^{\text {max }}\right) \cdot \prod_{f} w_{\sigma, \lambda, \theta}(f)^{H(f)}=\boldsymbol{q}$.
Proof. Let $D$ be a dimer configuration of type ( $\sigma, \lambda, v, \theta$ ). By Lemma 3.3.7, $\left(D \cup D^{\max }\right) \backslash\left(D \cap D^{\max }\right)$ is a disjoint union $\sqcup \mathscr{C}_{\gamma}$ of a finite number of positive cycles. We define a map $H_{D}: \mathscr{F} \rightarrow \mathbb{Z}_{\geq 0}$ by

$$
H_{D}(f):=\sharp\left\{\mathscr{C}_{\gamma} \mid f \in \mathscr{C}_{\gamma}^{\circ}\right\} .
$$

Then we can verify the claim using Lemma 3.2.4.
Remark. The graph of the map $m \mapsto G_{\sigma, \lambda, \theta}^{v}(n, m)+2 H(n, m)$ determines a Young diagram. This is what we denote by ${ }^{\mathscr{V}}(n)$ in [Nagao 2011b, §3.1].
3.4. Generating function. From the description given by Proposition 3.3.9, we can verify that, fixing a monomial $\boldsymbol{q}$, we have only a finite number of dimer configurations of type ( $\sigma, \lambda, \nu, \theta$ ) with weight $\boldsymbol{q}$.

Definition 3.4.1. We define the generating function by

$$
\mathscr{L}_{\sigma, \lambda, v, \theta}=\mathscr{L}_{\sigma, \lambda, v, \theta}(\vec{q}):=\sum_{D} w_{\sigma, \lambda, \theta}(D),
$$

where the sum is taken over all dimer configurations of type $(\sigma, \lambda, \nu, \theta)$. In particular, we put

$$
\mathscr{E}_{\sigma, \lambda, v}^{\mathrm{NCDT}}:=\mathscr{Z}_{\sigma, \lambda, v, \text {,id } \mathbb{I d}_{\mathrm{h}}} .
$$

Remark. Note that $\mathscr{L}_{\sigma, \lambda, v}^{\mathrm{NCDT}} \cdot w_{\sigma, \lambda, \theta}\left(D_{\sigma, \lambda, v . \mathrm{id}}^{\max }\right)^{-1}$ is a formal power series in $q_{+}, q_{-}$ and $q_{1}, \ldots, q_{L-1}$.

## 4. Dimer shuffling and wall-crossing formula

4.1. Dimer shuffling at a hexagon. In this and next subsections, we study the relation between dimer configurations of type ( $\sigma, \lambda, \nu, \theta$ ) and of type ( $\sigma, \lambda, \nu, \mu_{i}(\theta)$ ) for $i \in I_{H}(\sigma, \theta)$.
4.1.1. For $(n, m) \in \mathscr{F}$ and $M \in \mathbb{Z}_{>0} \sqcup\{\infty\}$. we put

$$
f(n, m ; \pm, M):=\bigcup_{m^{\prime}=0}^{M-1} f\left(n, m \pm m^{\prime}\right)
$$

We define $\partial f(n, m ; \pm, M)$ and $\partial^{ \pm} f(n, m ; \pm, M)$ in the same way as in $\S 2.1 .9$ and §3.2.2.
4.1.2. For a dimer configuration $D$ and $n \in B_{\sigma, \lambda, \theta}^{i, \pm}$, let $m(D, n)$ denote the unique integer such that

$$
\partial f(n, m(D, n) ; \sigma(i), \infty) \cap D=\partial^{ \pm} f(n, m(D, n) ; \sigma(i), \infty)
$$

4.1.3. For a dimer configuration $D$ and $i \in I$, we consider the following conditions:
(4-1) $\partial f \cap D \neq \partial^{-} f$ for any $f \in \mathscr{F}_{i}$,
(4-2) $\quad \partial f \cap D \neq \partial^{+} f$ if $f \in \mathscr{F}_{i} \backslash\left\{f(n, m(D, n)) \mid n \in B_{\sigma, \lambda, \theta}^{i, \pm}\right\}$,
(4-3) $\quad \partial f(n, m(D, n)-2 \sigma(i)) \cap D \neq \partial^{-} f(n, m(D, n)-2 \sigma(i))$ for $n \in B_{\sigma, \lambda, \theta}^{i, \pm}$.
4.1.4. For a dimer configuration $D^{\circ}$ of type ( $\sigma, \lambda, \nu, \theta$ ) satisfying the condition (4-1), we set

$$
E_{i}\left(D^{\circ}\right):=\left\{(n, m) \in \mathscr{F}_{i} \mid \partial f(n, m) \cap D^{\circ}=\partial^{+} f(n, m)\right\},
$$

and define the map $M_{D^{\circ}}^{i}: E_{i}\left(D^{\circ}\right) \rightarrow \mathbb{Z}_{>0} \sqcup\{\infty\}$ by

$$
M_{D^{\circ}}^{i}(n, m):=\max \left\{M \mid \partial f(n, m ; \sigma(i), M) \cap D^{\circ}=\partial^{+} f(n, m ; \sigma(i), M)\right\} .
$$

Note that

$$
\left(M_{D^{\circ}}^{i}\right)^{-1}(\infty)=\left\{\left(n, m_{n}\right) \mid n \in B_{\sigma, \lambda, \theta}^{i,+}\right\} .
$$

We put $E_{i}^{\mathrm{fin}}\left(D^{\circ}\right):=E_{i}\left(D^{\circ}\right) \backslash\left(M_{D^{\circ}}^{i}\right)^{-1}(\infty)$.
Definition 4.1.5. For a dimer configuration $D^{\circ}$ of type ( $\sigma, \lambda, \nu, \theta$ ) satisfying the condition (4-1), let $\mu_{i}\left(D^{\circ}\right)$ be the a dimer configuration of type ( $\sigma, \lambda, \nu, \mu_{i}(\theta)$ ) given by

$$
\begin{aligned}
&\left(D^{\circ} \backslash\left(\bigcup_{(n, m) \in E_{i}\left(D^{\circ}\right)} \partial^{+} f\left(n, m ; \sigma(i), M_{D^{\circ}}^{i}(n, m)\right) \cup \bigcup_{n \in B_{\sigma, \lambda, \theta}^{i,-}} \partial^{-} f(n, m ; \sigma(i), \infty)\right)\right) \\
& \sqcup\left(\bigcup_{(n, m) \in E_{i}\left(D^{\circ}\right)} \partial^{-} f\left(n, m ; \sigma(i), M_{D^{\circ}}^{i}(n, m)\right) \cup \bigcup_{n \in B_{\sigma, \lambda, \theta}^{i,-}} \partial^{+} f(n, m ; \sigma(i), \infty)\right) .
\end{aligned}
$$

Note that $\mu_{i}\left(D^{\circ}\right)$ satisfies the condition (4-2) and (4-3).
Example 4.1.6. Here are some examples of dimer shuffling at hexagons.








Lemma 4.1.7.

$$
w_{\sigma, \lambda, \mu_{i}(\theta)}\left(\mu_{i}\left(D^{\circ}\right)\right)=w_{\sigma, \lambda, \theta}\left(D^{\circ}\right)
$$

Proof. For $n \in \pi^{-1}(i)$ and $m \in \mathbb{Z}$ such that $n+m$ is odd, we put

$$
D^{\circ}(n, m):=\left\{e_{\mathrm{s}}\left(n+\varepsilon_{1}, m+\varepsilon_{2}\right)\left(\varepsilon_{1}, \varepsilon_{2}= \pm 1 / 2\right)\right\} \cap D^{\circ}
$$

Assume that

$$
\begin{equation*}
(n, m-1),(n, m+1) \notin \bigcup_{(n, m) \in E_{i}\left(D^{\circ}\right)} f\left(n, m ; \sigma(i), M_{D^{\circ}}^{i}(n, m)\right) . \tag{4-4}
\end{equation*}
$$

Then $D^{\circ}(n, m)$ is one of the following:

$$
\varnothing, \quad\left\{e_{\mathrm{s}}(n \pm 1 / 2, m \pm 1 / 2)\right\}, \quad\left\{e_{\mathrm{s}}(n \pm 1 / 2, m \mp 1 / 2)\right\} .
$$

In particular, we have

$$
w_{\sigma, \lambda, \theta}\left(D^{\circ}(n, m)\right)=w_{\sigma, \lambda, \mu_{i}(\theta)}\left(D^{\circ}(n, m)\right)
$$

Hence

$$
w_{\sigma, \lambda, \theta}\left(D^{\circ} \cap \mu_{i}\left(D^{\circ}\right)\right)=w_{\sigma, \lambda, \mu_{i}(\theta)}\left(D^{\circ} \cap \mu_{i}\left(D^{\circ}\right)\right)
$$

The claim follows from this and the fact that

$$
w_{\sigma, \lambda, \theta}\left(\partial^{ \pm} f(n, m, M)\right)=w_{\sigma, \lambda, \mu_{i}(\theta)}\left(\partial^{\mp} f(n, m, M)\right)
$$

for $n \in \pi^{-1}(i)$.

### 4.2. Wall-crossing formula at a hexagon.

## Lemma 4.2.1.

$\mathscr{L}_{\sigma, \lambda, v, \theta}=\sum_{D^{\circ}} w_{\sigma, \lambda, \theta}\left(D^{\circ}\right) \prod_{n \in B_{\sigma, \lambda, \theta}^{i,+}} \frac{1}{1+w_{\sigma, \lambda, \theta}(n)} \prod_{(n, m) \in E_{i}^{\text {fin }}\left(D^{\circ}\right)} \frac{1+w_{\sigma, \lambda, \theta}(n)^{M_{D^{\circ}(n, m)+1}^{i}}}{1+w_{\sigma, \lambda, \theta}(n)}$,
where the sum is taken over all dimer configurations $D^{\circ}$ of type $(\sigma, \lambda, \nu, \theta)$ satisfying the condition (4-1).

Proof. For a map $s: E_{i}\left(D^{\circ}\right) \rightarrow \mathbb{Z}_{\geq 0}$ such that $s(n, m) \leq M_{D^{\circ}}^{i}(n, m)$, we define the dimer configuration

$$
\begin{aligned}
& D_{s}^{\circ}:=\left(D^{\circ} \backslash \bigcup_{(n, m) \in E_{i}\left(D^{\circ}\right)} \partial^{+} f(n, m ; \sigma(i), s(n, m))\right) \\
& \sqcup \bigcup_{(n, m) \in E_{i}\left(D^{\circ}\right)} \partial^{-} f(n, m ; \sigma(i), s(n, m))
\end{aligned}
$$

Then

$$
w_{\sigma, \lambda, \theta}\left(D_{s}^{\circ}\right)=w_{\sigma, \lambda, \theta}\left(D^{\circ}\right) \prod_{(n, m) \in E_{i}\left(D^{\circ}\right)} w_{\sigma, \lambda, \theta}(n)^{s(n, m)}
$$

Note that any dimer configuration $D$ is uniquely realized as $D^{\circ}(s)$ by some $D^{\circ}$ and $s$. Hence we have

$$
\begin{aligned}
& \mathscr{Z}_{\sigma, \lambda, v, \theta}= \sum_{D^{\circ}} w_{\sigma, \lambda, \theta}\left(D^{\circ}\right) \cdot\left(\sum_{s} \prod_{(n, m) \in E_{i}\left(D^{\circ}\right)} w_{\sigma, \lambda, \theta}(n)^{s(n, m)}\right) \\
&=\sum_{D^{\circ}} w_{\sigma, \lambda, v, \theta}\left(D^{\circ}\right) \prod_{n \in B_{\sigma, \lambda, \theta}^{i,+}} \frac{1}{1-w_{\sigma, \lambda, \theta}(n)} \\
& \quad \times \prod_{(n, m) \in E_{i}^{\text {fin }}\left(D^{\circ}\right)} \frac{1+w_{\sigma, \lambda, \theta}(n)^{M_{D^{\circ}}^{i}(n, m)+1}}{1+w_{\sigma, \lambda, \theta}(n)} .
\end{aligned}
$$

## Theorem 4.2.2.

$$
\mathscr{L}_{\sigma, \lambda, v, \mu_{i}(\theta)}=\mathscr{L}_{\sigma, \lambda, v, \theta} \prod_{n \in B_{\sigma, \lambda, \theta}^{i,+}}\left(1-w_{\sigma, \lambda, \theta}(n)\right) \prod_{n \in B_{\sigma, \lambda, \theta}^{i,-}} \frac{1}{1-w_{\sigma, \lambda, \theta}(n)} .
$$

Proof. As Lemma 4.2.1, we get

$$
\begin{aligned}
\mathscr{L}_{\sigma, \lambda, v, \mu_{i}(\theta)}=\sum_{D^{\bullet}} w_{\sigma, \lambda, \mu_{i}(\theta)}\left(D^{\bullet}\right) \prod_{n \in B_{\sigma, \lambda, \mu_{i}(\theta)}^{i,+}} & \frac{1}{1-w_{\sigma, \lambda, \mu_{i}(\theta)}(n)^{-1}} \\
& \times \prod_{(n, m) \in \check{E}_{i}\left(D^{\bullet}\right)} \frac{1+w_{\sigma, \lambda, \mu_{i}(\theta)}(n)^{-\check{M}_{D^{\bullet}}^{i}(n, m)-1}}{1+w_{\sigma, \lambda, \mu_{i}(\theta)}(n)^{-1}},
\end{aligned}
$$

where the sum is taken over all dimer configurations $D^{\bullet}$ of type $\left(\sigma, \lambda, \nu, \mu_{i}(\theta)\right)$ satisfying (4-2), (4-3), and

$$
\begin{aligned}
\check{E}_{i}\left(D^{\bullet}\right) & :=\left\{(n, m) \in \mathscr{F}_{i} \mid \partial f(n, m) \cap D^{\bullet}=\partial^{-} f(n, m)\right\}, \\
\check{M}_{D}^{i}(n, m) & :=\max \left\{M \mid \partial f(n, m ; \sigma(i), M) \cap D^{\bullet}=\partial^{-} f(n, m ;-\sigma(i), M)\right\} .
\end{aligned}
$$

Note that $\mu_{i}$ gives a one-to-one correspondence between dimer configurations of type ( $\sigma, \lambda, \nu, \theta$ ) satisfying (4-1) and those of type ( $\sigma, \lambda, \nu, \mu_{i}(\theta)$ ) satisfying (4-2) and (4-3). Hence the claim follows from

- $B_{\sigma, \lambda, \mu_{i}(\theta)}^{i, \pm}=B_{\sigma, \lambda, \theta}^{i, \mp}$,
- $w_{\sigma, \lambda, \mu_{i}(\theta)}(n)=w_{\sigma, \lambda, \theta}(n)^{-1}$ for $n \in \pi^{-1}(i)$,
- $(n, m) \mapsto\left(n, m+\sigma(i) \cdot\left(M_{D^{\circ}}^{i}(n, m)-1\right)\right)$ gives a bijection between $E_{i}^{\text {fin }}\left(D^{\circ}\right)$ and $\check{E}_{i}\left(\mu_{i}\left(D^{\circ}\right)\right)$ which respects $M_{D^{\circ}}^{i}$ and $\check{M}_{\mu_{i}\left(D^{\circ}\right)}^{i}$,
and Lemma 4.2.1.
4.3. Dimer shuffing at a quadrilateral. In this subsection, we study the relation between dimer configurations of type ( $\sigma, \lambda, \nu, \theta$ ) and of type $\left(\sigma, \lambda, \nu, \mu_{i}(\theta)\right.$ ) for $i \in I_{S}(\sigma, \theta)$.
4.3.1. For a dimer configuration $D^{\circ}$ of type ( $\sigma, \lambda, \nu, \theta$ ) satisfying the condition (4-1) and $n \in \pi^{-1}(i)$, we define

$$
\begin{aligned}
& E_{n}^{1}\left(D^{\circ}\right):=\left\{(n, m) \in \mathscr{F} \mid \partial f(n, m) \cap D^{\circ}=\partial^{+} f(n, m)\right\}, \\
& E_{n}^{2}\left(D^{\circ}\right):=\left\{(n, m) \in \mathscr{F} \mid \partial f(n, m) \cap D^{\circ}=\varnothing\right\} .
\end{aligned}
$$

Lemma 4.3.2. $\quad\left|E_{n}^{1}\left(D^{\circ}\right)\right|-\left|E_{n}^{2}\left(D^{\circ}\right)\right|= \begin{cases}\mp 1 & \text { if } n \in B_{\sigma, \lambda, \theta}^{i, \pm}, \\ 0 & \text { otherwise } .\end{cases}$ (See (1-4) for notation.)

Proof. For $n, m \in \mathbb{Z}$ such that $n+m$ is odd, we define $\varepsilon_{D^{\circ}}(n, m)$ by

$$
\varepsilon_{D^{\circ}}(n, m):= \begin{cases}+ & \text { if } e_{\mathrm{s}}(n+1 / 2, m+1 / 2), e_{\mathrm{s}}(n-1 / 2, m+1 / 2) \notin D, \\ - & \text { if } e_{\mathrm{s}}(n+1 / 2, m-1 / 2), e_{\mathrm{s}}(n-1 / 2, m-1 / 2) \notin D .\end{cases}
$$

Then for $(n, m) \in \mathscr{F}$, we have

$$
\begin{aligned}
& (n, m) \in E_{n}^{1}\left(D^{\circ}\right) \Longleftrightarrow \varepsilon_{D^{\circ}}(n, m \pm 1)= \pm, \\
& (n, m) \in E_{n}^{2}\left(D^{\circ}\right) \Longleftrightarrow \varepsilon_{D^{\circ}}(n, m \pm 1)=\mp,
\end{aligned}
$$

and $\varepsilon_{D^{\circ}}(n, m)=\mp \tilde{\lambda}(n \pm 1 / 2)$ if $\tilde{\sigma}(n \pm 1 / 2) \cdot m \gg 0$. The claim follows.
4.3.3. For a dimer configuration $D^{\circ}$ of type ( $\sigma, \lambda, \nu, \theta$ ) satisfying the condition (4-1), we define a dimer configuration $\mu_{i}\left(D^{\circ}\right)$ of type $\left(\sigma, \lambda, v, \mu_{i}(\theta)\right)$ as follows:

- If $\pi(h) \neq i \pm 1 / 2$, we have

$$
e_{\mathrm{s}}(h, k) \in D^{\circ} \Longleftrightarrow e_{\mathrm{s}}(h, k) \in \mu_{i}\left(D^{\circ}\right),
$$

- If $n \in I_{H}(\sigma, \theta)$ and $\pi(n) \neq i \pm 1$, we have

$$
e_{\mathrm{h}}(n, m) \in D^{\circ} \Longleftrightarrow e_{\mathrm{h}}(n, m) \in \mu_{i}\left(D^{\circ}\right)
$$

- For $(n, m) \in \mathscr{F}_{i}$ we have

$$
\begin{aligned}
D^{\circ}(f(n, m))=\varnothing & \Longleftrightarrow \mu_{i}\left(D^{\circ}\right)(f(n, m))=\partial_{\sigma, \mu_{i}(\theta)}^{-}(f(n, m)), \\
D^{\circ}(f(n, m))=\partial_{\sigma, \theta}^{+}(f(n, m)) & \Longleftrightarrow \mu_{i}\left(D^{\circ}\right)(f(n, m))=\varnothing,
\end{aligned}
$$

(Here we use notation such as $\partial_{\sigma, \theta}^{ \pm}(f(n, m))$ in order to emphasize that the notions like $\partial^{ \pm}(f(n, m))$ given in $\S 2.1 .9$ depend on $\sigma$ and $\theta$.)

- If $D^{\circ}(f(n, m)) \neq \varnothing, \partial_{\sigma, \theta}^{+}(f(n, m))$ for $(n, m) \in \mathscr{F}_{i}$, we have

$$
e_{\mathrm{s}}\left(n+\varepsilon_{1}, m+\varepsilon_{2}\right) \in D^{\circ} \Longleftrightarrow e_{\mathrm{s}}\left(n-\varepsilon_{1}, m-\varepsilon_{2}\right) \in \mu_{i}\left(D^{\circ}\right) \quad\left(\varepsilon_{1}, \varepsilon_{2}= \pm 1 / 2\right)
$$

- If $\sigma(i \pm 3 / 2) \neq \sigma(i \pm 1 / 2)$, we have

$$
e_{\mathrm{s}}(n \pm 1 / 2, m-1), e_{\mathrm{s}}(n \pm 1 / 2, m+1) \notin D^{\circ} \Longleftrightarrow e_{\mathrm{h}}(n \pm 1, m) \in \mu_{i}\left(D^{\circ}\right)
$$

Note that $\mu_{i}\left(D^{\circ}\right)$ satisfies the condition

$$
\begin{equation*}
D(f) \neq \partial^{+} f \text { for any } f \in \mathscr{F}_{i} . \tag{4-5}
\end{equation*}
$$

Example 4.3.4. Here are some examples of dimer shuffling at squares.






Lemma 4.3.5.

$$
w_{\sigma, \lambda, \mu_{i}(\theta)}\left(\mu_{i}\left(D^{\circ}\right)\right)=w_{\sigma, \lambda, \theta}\left(D^{\circ}\right) .
$$

Proof. We have $w_{\sigma, \lambda, \theta}\left(\partial_{\sigma, \theta}^{+} f\right)=w_{\sigma, \lambda, \mu_{i}(\theta)}\left(\partial_{\sigma, \mu_{i}(\theta)}^{-} f\right)$ for $f \in \mathscr{F}_{i}$, and

$$
w_{\sigma, \lambda, \theta}\left(\partial_{\sigma, \theta}^{+} f\right)= \begin{cases}1 & \text { if } n \in B_{\sigma}^{i,+,}, \theta \\ w_{\sigma, \lambda, \theta}(n)^{-1} & \text { if } n \in B_{\sigma, \lambda, \theta}^{i--}\end{cases}
$$

Thus, the claim follows from Lemma 4.3.2 and (2-9).

### 4.4. Wall-crossing formula at a quadrilateral.

Lemma 4.4.1. $\quad \mathscr{L}_{\sigma, \lambda, v, \theta}=\sum_{D^{\circ}} w_{\sigma, \lambda, \theta}\left(D^{\circ}\right) \cdot \prod_{n \in \pi^{-1}(i)}\left(1+w_{\sigma, \lambda, \theta}(n)\right)^{\left|E_{n}^{1}\left(D^{\circ}\right)\right|}$.
Proof. We set

$$
E_{i}^{1}\left(D^{\circ}\right):=\bigcup_{n \in \pi^{-1}(i)} E_{n}^{1}\left(D^{\circ}\right), \quad E_{i}^{2}\left(D^{\circ}\right):=\bigcup_{n \in \pi^{-1}(i)} E_{n}^{2}\left(D^{\circ}\right) .
$$

Given a subset $S \subset E_{i}^{1}\left(D^{\circ}\right)$, we get a dimer configuration $D_{S}^{\circ}$ of type ( $\sigma, \lambda, \nu, \theta$ ) such that

$$
D_{S}^{\circ}:=\left(D \backslash \bigcup \partial^{+} f\right) \cup \bigcup \partial^{+} f,
$$

and we have

$$
w_{\sigma, \lambda, \theta}\left(D_{S}^{\circ}\right)=w_{\sigma, \lambda, \theta}\left(D^{\circ}\right) \prod_{(n, m) \in S} w_{\sigma, \lambda, \theta}(n) .
$$

Note that any dimer configuration $D$ is uniquely realized as $D_{S}^{\circ}$ by some $D^{\circ}$ and $S$. Hence we have

$$
\begin{aligned}
\mathscr{Z}_{\sigma, \lambda, v, \theta} & =\sum_{D^{\circ}} w_{\sigma, \lambda, \theta}\left(D^{\circ}\right)\left(\sum_{S} \prod_{(n, m) \in S} w_{\sigma, \lambda, \theta}(n)\right) \\
& =\sum_{D^{\circ}} w_{\sigma, \lambda, \theta}\left(D^{\circ}\right) \prod_{(n, m) \in E_{i}^{1}\left(D^{\circ}\right)}\left(1+w_{\sigma, \lambda, \theta}(n)\right) \\
& =\sum_{D^{\circ}} w_{\sigma, \lambda, \theta}\left(D^{\circ}\right) \prod_{n \in \pi^{-1}(i)}\left(1+w_{\sigma, \lambda, \theta}(n)\right)^{\left|E_{n}^{1}\left(D^{\circ}\right)\right|} .
\end{aligned}
$$

## Theorem 4.4.2.

$$
\mathscr{L}_{\sigma, \lambda, v, \mu_{i}(\theta)}=\mathscr{L}_{\sigma, \lambda, v, \theta} \prod_{n \in B_{\sigma, \lambda, \theta}^{i,+}}\left(1+w_{\sigma, \lambda, \theta}(n)\right)^{-1} \prod_{n \in B_{\sigma, \lambda, \theta}^{i,-}}\left(1+w_{\sigma, \lambda, \theta}(n)\right) .
$$

Proof. Let $D^{\bullet}$ be a dimer configuration of type ( $\sigma, \lambda, \nu, \mu_{i}(\theta)$ ) satisfying (4-5). We put

$$
\tilde{E}_{n}^{1}\left(D^{\bullet}\right):=\left\{(n, m) \in \mathscr{F} \mid \partial_{\sigma, \mu_{i}(\theta)} f(n, m) \cap D^{\bullet}=\partial_{\sigma, \mu_{i}(\theta)}^{-} f(n, m)\right\} .
$$

Then, as Lemma 4.4.1, we get

$$
\mathscr{L}_{\sigma, \lambda, \nu, \mu_{i}(\theta)}=\sum_{D^{\bullet}} w_{\sigma, \lambda, \mu_{i}(\theta)}\left(D^{\bullet}\right) \prod_{n \in \pi^{-1}(i)}\left(1+w_{\sigma, \lambda, \mu_{i}(\theta)}(n)^{-1}\right)^{\left|\tilde{E}_{n}^{1}\left(D^{\bullet}\right)\right|},
$$

where the sum is taken over all dimer configurations $D^{\bullet}$ of type $\left(\sigma, \lambda, \nu, \mu_{i}(\theta)\right)$ satisfying the condition (4-5). Note that $\mu_{i}$ gives a one-to-one correspondence of dimer configurations of type ( $\sigma, \lambda, \nu, \theta$ ) satisfying the condition (4-1) and ones of type ( $\sigma, \lambda, \nu, \mu_{i}(\theta)$ ) satisfying the condition (4-5). Hence the claim follows from the equalities $\tilde{E}_{n}^{1}\left(\mu_{i}\left(D^{\circ}\right)\right)=E_{n}^{2}\left(D^{\circ}\right)$ and $w_{\sigma, \lambda, \mu_{i}(\theta)}(n)=w_{\sigma, \lambda, \theta}(n)^{-1}$, both valid for $n \in \pi^{-1}(i)$, together with Lemma 4.4.1.
4.5. Conclusion. For $\sigma$ and $\alpha \in \Lambda^{\mathrm{re},+}$, we put

$$
\begin{equation*}
\sigma(\alpha):=\sigma\left(j^{-}(\alpha)\right) \cdot \sigma\left(j^{+}(\alpha)\right) . \tag{4-6}
\end{equation*}
$$

Combining Theorem 4.2.2 and 4.4.2, we get:

Theorem 4.5.1. $\mathscr{L}_{\sigma, \lambda, v, \theta}$ has the value
$\mathscr{Z}_{\sigma, \lambda, \nu}^{\mathrm{NCDT}} \prod_{\alpha \in \Lambda_{\theta}^{\mathrm{r},+}}\left(\prod_{\left(h, h^{\prime}\right) \in B_{\sigma, \lambda}^{\alpha,+}}\left(1-\sigma(\alpha) \frac{w_{\lambda}\left(h^{\prime}\right)}{w_{\lambda}(h)}\right)^{\sigma(\alpha)} \prod_{\left(h, h^{\prime}\right) \in B_{\sigma, \lambda}^{\alpha,-}}\left(1-\sigma(\alpha) \frac{w_{\lambda}\left(h^{\prime}\right)}{w_{\lambda}(h)}\right)^{-\sigma(\alpha)}\right)$.
(See (1-2) and (1-3) for notation.)
Since the second term in this expression does not depend on $v$, we have:

Corollary 4.5.2.

$$
\frac{\mathscr{L}_{\sigma, \lambda, v, \theta}}{\mathscr{Z}_{\sigma, \lambda, \vec{\varnothing}, \theta}}=\frac{\mathscr{L}_{\sigma, \lambda, \nu}^{\mathrm{NCDT}}}{\mathscr{Z}_{\sigma, \lambda, \vec{\varnothing}}^{\mathrm{NCDT}}} .
$$

Lemma 1.3.3 and Theorem 4.5.1 yield:
Theorem 4.5.3. (See (1-2) for notation.)

$$
\begin{aligned}
&\left.\mathscr{L}_{\sigma, \lambda, v, \theta}\right|_{q_{+}=}=q_{-}=\left(q_{0}\right)^{1 / 2} \\
&=\left.\mathscr{L}_{\sigma, \lambda, v}^{\mathrm{NCDT}}\right|_{q_{+}=q_{-}=\left(q_{0}\right)^{1 / 2}} \prod_{\alpha \in \Lambda_{\theta}^{\mathrm{re},+}}\left(1-\sigma(\alpha) \cdot q^{\alpha}\right)^{\sigma(\alpha)\left[\alpha^{0}+c_{\lambda}\left(j_{-}(\alpha)\right)-c_{\lambda}\left(j_{+}(\alpha)\right)\right]} .
\end{aligned}
$$

Since the second term on the right depends only on the $c_{\lambda}[j]$ and not on $\lambda$ and $v$, we have:

Corollary 4.5.4. If $c_{\lambda}[j]=0$ for any $j$, we have

$$
\frac{\mathscr{Z}_{\sigma, \lambda, v, \theta}}{\left.\mathscr{L}_{\sigma, \vec{\varnothing}, \vec{\varnothing}, \theta}\right|_{q_{+}=q_{-}}}=\frac{\mathscr{Z}_{\sigma, \lambda, v}^{\mathrm{NCDT}}}{\left.\mathscr{E}_{\sigma, \vec{\varnothing}, \vec{\varnothing}}\right|_{q_{+}=q_{-}}} .
$$

## 5. Refined topological vertex via dimer model

### 5.1. Refined topological vertex for $\mathbb{C}^{3}$.

5.1.1. A Young diagram can be regarded as a subset of $\left(\mathbb{Z}_{\geq 0}\right)^{2}$. For a Young diagram $\lambda$, let

$$
\begin{aligned}
\Lambda^{\mathrm{x}}(\lambda) & =\left\{(x, y, z) \in\left(\mathbb{Z}_{\geq 0}\right)^{3} \mid(y, z) \in \lambda\right\} \\
\Lambda^{\mathrm{y}}(\lambda) & =\left\{(x, y, z) \in\left(\mathbb{Z}_{\geq 0}\right)^{3} \mid(z, x) \in \lambda\right\} \\
\Lambda^{\mathrm{z}}(\lambda) & =\left\{(x, y, z) \in\left(\mathbb{Z}_{\geq 0}\right)^{3} \mid(x, y) \in \lambda\right\} .
\end{aligned}
$$

5.1.2. Given a triple $\left(\lambda_{x}, \lambda_{y}, \lambda_{z}\right)$ of Young diagrams, define

$$
\Lambda^{\min }:=\Lambda^{\mathrm{x}}\left(\lambda_{x}\right) \cup \Lambda^{\mathrm{y}}\left(\lambda_{y}\right) \cup \Lambda^{\mathrm{z}}\left(\lambda_{z}\right) \subset\left(\mathbb{Z}_{\geq 0}\right)^{3}
$$

5.1.3. A subset $\Lambda$ of $\left(\mathbb{Z}_{\geq 0}\right)^{3}$ is said to be a 3-dimensional Young diagram of type ( $\lambda_{x}, \lambda_{y}, \lambda_{z}$ ) if the following conditions are satisfied:

- If $(x, y, z) \notin \Lambda$, then $(x+1, y, z),(x, y+1, z),(x, y, z+1) \notin \Lambda$.
- $\Lambda \supset \Lambda^{\min }$.
- $\left|\Lambda \backslash \Lambda^{\min }\right|<\infty$.
5.1.4. For a Young diagram $\lambda$, we define a monomial $w_{\lambda}(m)$ for each $m \in \mathbb{Z}$ by

$$
\begin{equation*}
w_{\lambda}(m)=q_{\lambda(m-1 / 2)} \cdot q_{\lambda(m+1 / 2)} \cdot q_{1} \cdots \cdots q_{L-1} . \tag{5-1}
\end{equation*}
$$

For a finite subset $S$ of $\left(\mathbb{Z}_{\geq 0}\right)^{3}$ we define the weight $w(S)$ by

$$
w(S):=\prod_{(x, y, z) \in S} w_{\lambda_{x}}(y-z)
$$

For a positive integer $N$, we set $C_{N}:=[0, N]^{3}$. Given a 3 -dimensional Young diagram $\Lambda$ of type ( $\lambda_{x}, \lambda_{y}, \lambda_{z}$ ), we take a sufficiently large $N$ such that $\Lambda \backslash \Lambda^{\min } \subset$ $C_{N}$ and define the weight $w(\Lambda)$ of $\Lambda$ by

$$
w(\Lambda):=\frac{w\left(\Lambda \cap C_{N}\right)}{w\left(\Lambda^{\mathrm{x}}\left(\lambda_{x}\right) \cap C_{N}\right) w\left(\Lambda^{\mathrm{y}}\left(\lambda_{y}\right) \cap C_{N}\right) w\left(\Lambda^{z}\left(\lambda_{z}\right) \cap C_{N}\right)} .
$$

Note that this is well-defined.
Remarks. - In the definition of $w(\Lambda)$, the three axes do not play the same role. The $x$-axis is called the preferred axis for the refined topological vertex.

- If we replace the definition (5-1) with

$$
\left(q_{\lambda(m-1 / 2)}\right)^{2} \cdot q_{1} \cdots \cdots q_{L-1}
$$

then the weight coincides with the one in [Iqbal et al. 2009]. Our weight coincides with the one in [Dimofte and Gukov 2010].
We define the generating function

$$
G_{\lambda_{x}, \lambda_{y}, \lambda_{z}}(\vec{q}):=\sum w(\Lambda),
$$

where the sum is taken over all 3 -dimensional Young diagrams of type $\left(\lambda_{x}, \lambda_{y}, \lambda_{z}\right)$.
5.2. Dimer model for $L=1$. In the case $L=1$, the graph in $\S 2.1 .2$ gives a hexagon lattice. As we have only two choices of $\sigma$, we put $\sigma(1 / 2)=+$. We take id as $\theta$. We omit $\sigma$ and id from the notation in this subsection. Note that $\lambda$ is a single 2-dimensional Young diagram.

It is well-known that giving a dimer configuration of type $(\lambda, \nu)$ is equivalent to giving a 3-dimensional Young diagram of type $\left(\lambda, v_{+},{ }^{t} v_{-}\right)$. Let $D(\Lambda)$ be the dimer configuration corresponding to a 3 -dimensional Young diagram $\Lambda$.

For a Young diagram $\eta=\left(\eta_{(1)}, \eta_{(2)}, \ldots\right)$ and a monomial $p$, we put

$$
w(\eta ; p, Q):=\prod\left(p Q^{i-1}\right)^{\eta_{(i)}}
$$

Then we can verify the following:

$$
\begin{equation*}
w_{\lambda}(D(\Lambda))=w\left(v_{-} ; q_{+}, Q_{+}\right) w\left(v_{+} ; q_{-}, Q_{-}\right) w(\Lambda) . \tag{5-2}
\end{equation*}
$$

Example 5.2.1. As we show in Figure 10, we have

$$
\begin{aligned}
& w_{\varnothing}\left(\Lambda_{\varnothing,(1), \varnothing}^{\min }\right)=w\left((1) ; q_{-}, Q_{-}\right)=q_{-}, \\
& w_{\varnothing}\left(\Lambda_{\varnothing,(2), \varnothing}^{\min }\right)=w\left((2) ; q_{-}, Q_{-}\right)=q_{-}^{2}, \\
& w_{\varnothing}\left(\Lambda_{\varnothing,(1,2), \varnothing}^{\min }\right)=w\left((2,1) ; q_{-}, Q_{-}\right)=q_{-}^{3} Q_{-} .
\end{aligned}
$$



Figure 10. $D\left(\Lambda_{\varnothing,(1), \varnothing}^{\min }\right), D\left(\Lambda_{\varnothing,(2), \varnothing}^{\min }\right)$ and $D\left(\Lambda_{\varnothing,(1,2), \varnothing}^{\min }\right)$.
In particular, we have

$$
\mathscr{L}_{\lambda, v}=w\left(v_{-} ; q_{+}, Q_{+}\right) \cdot w\left(v_{+} ; q_{-}, Q_{-}\right) \cdot G_{\lambda, v_{+},{ }^{v_{-}}},
$$

where $\mathscr{L}_{\lambda, \nu}$ is the generating function given in Definition 3.4.1.
5.3. Refined topological vertex for a small resolution. We will define generating functions $\mathscr{£}_{\sigma, \lambda, v}^{\mathrm{RTV}}(\vec{q})$. First, we consider the following data: let $\vec{v}=\left(v^{(1)}, \ldots, v^{(L-1)}\right)$ be an $(L-1)$-tuple of Young diagrams and $\vec{\Lambda}=\left(\Lambda^{(1 / 2)}, \ldots, \Lambda^{(L-1 / 2)}\right)$ be an $L$ tuple of 3-dimensional Young diagrams such that $\Lambda^{(j)}$ is

- of type $\left(\lambda^{(j)}, v^{(j+1 / 2)},{ }^{\mathrm{t}} v^{(j-1 / 2)}\right)$ if $\sigma(j)=+$,
- of type $\left(\lambda^{(j)},{ }^{\mathrm{t}} \nu^{(j-1 / 2)}, v^{(j+1 / 2)}\right)$ if $\sigma(j)=-$,
where we put $\nu^{(0)}:=v_{-}$and $\nu^{(L)}:=v_{+}$. We say that the data $(\vec{\Lambda}, \vec{v})$ is of type $(\sigma, \lambda, v)$. We define the weight $w(\vec{\Lambda}, \vec{v})$ of the data $(\vec{\Lambda}, \vec{v})$ by

$$
w_{\sigma}(\vec{\Lambda}, \vec{v}):=w\left(v_{+} ; q_{-}, Q_{-}\right) \cdot w\left(v_{-} ; q_{+}, Q_{+}\right)\left(\prod_{j=1 / 2}^{L-1 / 2} w\left(\Lambda^{(j)}\right)\right)\left(\prod_{i=1}^{L-1} w_{\sigma}^{i}\left(\mu^{(i)}\right)\right)
$$

where $w_{\sigma}^{i}\left(\mu^{(i)}\right)$ is given by

$$
w_{\sigma}^{i}\left(\mu^{(i)}\right):=\prod_{(\alpha, \beta) \in \mu^{i}} \begin{cases}q_{i} \cdot Q^{2 \alpha+1} & \text { if } \sigma\left(i-\frac{1}{2}\right)=\sigma\left(i+\frac{1}{2}\right)=+ \\ q_{i} \cdot Q^{2 \beta+1} & \text { if } \sigma\left(i-\frac{1}{2}\right)=\sigma\left(i+\frac{1}{2}\right)=- \\ q_{i} \cdot Q \cdot Q_{+}^{\alpha} \cdot Q_{-}^{\beta} & \text { if } \sigma\left(i-\frac{1}{2}\right)=+, \quad \sigma\left(i+\frac{1}{2}\right)=- \\ q_{i} \cdot Q \cdot Q_{-}^{\alpha} \cdot Q_{+}^{\beta} & \text { if } \sigma\left(i-\frac{1}{2}\right)=-, \quad \sigma\left(i+\frac{1}{2}\right)=+\end{cases}
$$

We consider the generating function

$$
\mathscr{E}_{\sigma, \lambda, \mu}^{\mathrm{RTV}}(\vec{q}):=\sum w_{\sigma}(\vec{\Lambda}, \vec{v})
$$

where the sum is taken over all the data as above.
Remark. This is the generating function of the refined topological vertex associated to $Y_{\sigma}$, where $Y_{\sigma} \rightarrow X$ is the crepant resolution constructed from $\sigma$ (see [Nagao 2011a, §1.1] for the construction of $Y_{\sigma}$ ). Here is the polygon corresponding to $Y_{\sigma}$, for $\sigma$ given by

$$
(\sigma(1 / 2), \ldots, \sigma(11 / 2))=(+,-,+,+,-,+):
$$



### 5.4. Limit behavior of the dimer model.

5.4.1. Let $\boldsymbol{i} \in I^{\mathbb{Z}_{>0}}$ be a minimal expression such that for any $N \in \mathbb{Z}_{\geq 0}$ we have $b(N) \in \mathbb{Z}_{>0}$ such that $\alpha_{i, b}>N \delta$ for any $b>b(N)$.

Lemma 5.4.2. Given $\sigma, \lambda$ and a monomial $\boldsymbol{q}$, there exists an integer $B_{1}$ such that the following condition holds: for any $b \geq B_{1}$,

- any dimer configuration of type ( $\sigma, \lambda, v, \theta_{i, b}$ ) with weight $\boldsymbol{q}$ satisfies (4-1),
- any dimer configuration of type $\left(\sigma, \lambda, v, \theta_{i, b+1}\right)$ with weight $\boldsymbol{q}$ satisfies (4-2), and
- $\mu_{i_{b}}$ gives a one-to-one correspondence between dimer configurations of type $\left(\sigma, \lambda, v, \theta_{i, b}\right)$ with weight $\left(\sigma, \lambda, v, \theta_{i, b+1}\right)$ with weight $\boldsymbol{q}$.

Proof. Take $N_{2}$ such that

$$
q^{N_{2} \delta}>\boldsymbol{q} \cdot w_{\sigma, \lambda, \theta}\left(D_{\sigma, \lambda, v . \mathrm{id}}^{\max }\right)^{-1}
$$

By Theorem 4.5.3 and the remark just before Section 4,

$$
\left.\mathscr{L}_{\sigma, \lambda, v, \theta} \cdot w_{\sigma, \lambda, \theta}\left(D_{\sigma, \lambda, v . \mathrm{id}}^{\max }\right)^{-1}\right|_{q_{+}=q_{-}=\left(q_{0}\right)^{1 / 2}}
$$

is a polynomial in $q_{0}, \ldots, q_{L-1}$. Thus, there does not exist any dimer configuration with weight $\boldsymbol{q}-\alpha(\boldsymbol{i}, b)$ for any $b>b\left(N_{2}\right)=: B_{1}$, where $b\left(N_{2}\right)$ is taken as in §5.4.1.

Assume that we have a dimer configuration type ( $\sigma, \lambda, v, \theta_{i, b}$ ) with weight $\boldsymbol{q}$ and $f \in \mathscr{F}$ such that $D(f)=\partial^{-}(f)$. Then we get a dimer configuration $D \cup$ $\partial^{+}(f) \backslash \partial^{-}(f)$ with weight $\boldsymbol{q}-\alpha(\theta, i)$, which is a contradiction. We can check the second claim similarly and the third claim immediately follows from the first and second ones.
5.4.3. Given $\sigma, \lambda$, we can take an integer $N_{2}$ such that

- $\tilde{\sigma}(h)= \pm \tilde{\lambda}(h)$ for any $h \in \mathbb{Z}_{\mathrm{h}}$ such that $\pm h>N_{2} L$,
- $e^{\mathrm{s}}(h, k) \notin D_{\sigma, \lambda, \theta_{i, B_{1}}}^{\max }$ for any $h$ and $k$ such that $h<N_{2} L$ and $h \cdot \tilde{\sigma}(h)-k$ is even, and
- $e^{\mathrm{s}}(h, k) \notin D_{\sigma, \lambda, \theta_{i, B_{1}}}^{\max }$ for any $h$ and $k$ such that $h>N_{2} L$ and $h \cdot \tilde{\sigma}(h)-k$ is odd.

Take a monomial $\boldsymbol{q}$. Since we have only a finite number of dimer configuration of type ( $\sigma, \lambda, v, \theta_{\boldsymbol{i}, B_{1}}$ ) with weight $\boldsymbol{q}$ and each dimer configuration has only finite difference with $D_{\sigma, \lambda, \nu, \theta_{i, B_{1}}}^{\max }$, we can take an integer $N_{4}$ such that

- $\tilde{\sigma}(h)= \pm \tilde{\lambda}(h)$ for any $h \in \mathbb{Z}_{\mathrm{h}}$ such that $\pm h>L N_{4}$,
- $e^{\mathrm{s}}(h, k) \notin D$ for any $h$ and $k$ such that $h<L N_{4}$ and $h \cdot \tilde{\sigma}(h)-k$ is even, and
- $e^{\mathrm{s}}(h, k) \notin D$ for any $h$ and $k$ such that $h>L N_{4}$ and $h \cdot \tilde{\sigma}(h)-k$ is odd.

Lemma 5.4.4. Let $D$ be a dimer configuration of type $(\sigma, \lambda, v, \theta)$ satisfying the condition (4-1). Take $h \in \pi^{-1}(i+1 / 2)$ such that $\tilde{\sigma}(h)=\tilde{\lambda}(h)$ and assume that $e_{\mathrm{s}}(h, k) \notin D$ for any $k \in \mathbb{Z}_{\mathrm{h}}$ such that $h \tilde{\sigma}(h)-k$ is odd. Then $e_{\mathrm{S}}(h-1, k-\tilde{\sigma}(h))$ is not in $\mu_{i}(D)$.

Similarly, take $h \in \pi^{-1}(i+1 / 2)$ such that $\tilde{\sigma}(h)=-\tilde{\lambda}(h)$ and assume that $e_{\mathrm{s}}(h, k) \notin D$ for any $k \in \mathbb{Z}_{\mathrm{h}}$ such that $h \tilde{\sigma}(h)-k$ is even. Then $e_{\mathrm{s}}(h+1, k+\tilde{\sigma}(h))$ is not in $\mu_{i}(D)$.
Proof. In the case $i \in I_{S}$, for any $h, k \in \mathbb{Z}_{\mathrm{h}}$ such that $\tilde{\sigma}(h)=\tilde{\lambda}(h)$ and $h \tilde{\sigma}(h)-k$ is odd, we can verify

$$
e_{\mathrm{s}}(h, k) \notin D \Longrightarrow e_{\mathrm{s}}(h-1, k-\tilde{\sigma}(h)) \notin \mu_{i}(D)
$$

from the definition of $\mu_{i}(D)$ in $\S 4.3 .3$.

In the case $i \in I_{S}$, assume we have $k \in \mathbb{Z}_{\mathrm{h}}$ such that $h \tilde{\sigma}(h)-k$ is odd and $e_{\mathrm{S}}(h-1, k-\tilde{\sigma}(h)) \in \mu_{i}(D)$. From Definition 4.1.5, we have $e_{\mathrm{S}}(h-1, k-\tilde{\sigma}(h)) \in D$. Since $e_{\mathrm{s}}(h, k-2 \tilde{\sigma}(h)) \notin D$, we have $e_{\mathrm{s}}(h, k-\tilde{\sigma}(h)) \in D$. Then, since $\tilde{\sigma}(h)=\tilde{\lambda}(h)$, there exists $m$ such that $\sigma(i)(m-k)>0$ and $\partial f(h-1 / 2, m) \cap D=\partial^{-} f(h-1 / 2, m)$, which is a contradiction.
5.4.5. Given $\sigma, \lambda$ and a monomial $\boldsymbol{q}$, take $B_{1}$ and $N_{4}$ as in Lemma 5.4.2 and §5.4.3. By the definition of $N_{4}$ and Lemma 5.4.4, we have the following lemma:

Lemma 5.4.6. For any $b \geq B_{1}$ and any dimer configuration of type ( $\sigma, \lambda, \nu, \theta_{i, b}$ ) with weight $\boldsymbol{q}$, we have

- $e_{\mathrm{s}}(h, k) \notin D$ for any $h$ and $k$ such that $h<\theta_{i, b}^{-1}(\pi(h))-2 L N_{4}$ and $h \cdot \tilde{\sigma}(h)-k$ is even, and
- $e_{\mathrm{s}}(h, k) \notin D$ for any $h$ and $k$ such that $h<\theta_{i, b}^{-1}(\pi(h))+2 L N_{4}$ and $h \cdot \tilde{\sigma}(h)-k$ is odd.
5.4.7. We assume that

$$
\theta_{i, b}^{-1}(1 / 2)<\theta_{i, b}^{-1}(3 / 2)<\cdots<\theta_{i, b}^{-1}(L-1 / 2)
$$

for any $b>0$.
Given $\sigma, \lambda$ and a monomial $\boldsymbol{q}$, take $B_{5}$ such that $B_{5}>b\left(2 N_{4}\right)$ and $B_{5}>B_{1}$. The following theorem is the main result of this section:

Theorem 5.4.8. For any $b>B_{5}$, we have a bijection between

- the set of dimer configurations of type ( $\sigma, \lambda, \nu, \theta_{i, b}$ ) with weights $\boldsymbol{q}$, and
- the set of data $(\vec{\Lambda}, \vec{v})$ as in Section 5.3 of type $(\sigma, \lambda, \nu)$ with weights $\boldsymbol{q}$.

Proof. First, we divide the ( $x, y$ )-plane into the following $2 L+1$ areas:

$$
\begin{aligned}
& C_{j}:=\left\{\theta^{-1}(j)-2 L N_{4}<x<\theta^{-1}(j)+2 L N_{4}\right\} \quad\left(j \in I_{\mathrm{h}}\right), \\
& C_{0}:=\left\{x<\theta^{-1}(1 / 2)-2 L N_{4}\right\}, \\
& C_{i}:=\left\{\theta^{-1}(i-1 / 2)+2 L N_{4}<x<\theta^{-1}(i+1 / 2)-2 L N_{4}\right\} \quad(1 \leq i \leq L-1), \\
& C_{L}:=\left\{\theta^{-1}(L-1 / 2)+2 L N_{4}<x\right\} .
\end{aligned}
$$

By Lemma 5.4.6, in the area $C_{j}$ we have

- $e^{\mathrm{s}}[h, k] \notin D$ for any $h$ and $k$ such that $\pi(h)>j$ and $h \cdot \tilde{\sigma}(h)-k$ is even;
- $e^{\mathrm{s}}[h, k] \notin D$ for any $h$ and $k$ such that $\pi(h)<j$ and $h \cdot \tilde{\sigma}(h)-k$ is odd.

Removing these edges, we get a new graph. A face of the new graph is a union of $L$-tuple of elements in $\mathscr{F}$. If we regard such a union as a hexagon, the dimer configuration $D$ gives a dimer configuration for the hexagon lattice - in other words, a
three-dimensional diagram. Let $\Lambda^{(j)}$ denote this three-dimensional diagram. (See Example 5.4.9.)

Similarly, in the area $C_{j}$ we have

- $e^{\mathrm{s}}[h, k] \notin D$ for any $h$ and $k$ such that $\pi(h)>i$ and $h \cdot \tilde{\sigma}(h)-k$ is even;
- $e^{\mathrm{s}}[h, k] \notin D$ for any $h$ and $k$ such that $\pi(h)<i$ and $h \cdot \tilde{\sigma}(h)-k$ is odd.

Removing these edges, we get a new graph, which is an infinite disjoint union of zigzag paths. For each zigzag path, we have two choices of perfect matching and so the dimer configuration $D$ gives a Young diagram $v^{(i)}$. We can verify that the datum $(\vec{\Lambda}, \vec{v})$ satisfies the conditions in Section 5.3. Note that the reverse construction also works.

We have to check the correspondence above respects the weights. Note that all edges of in the area $C_{i}$ have weights $=1$. By (5-2), the contribution of the part in the area $C_{j}$ is given by

$$
\begin{aligned}
& w\left(v^{(j-1 / 2)} ; q_{+}^{\left(s_{i}(j)\right)}, Q_{+}\right) w\left(v^{(j+1 / 2)} ;\left(q_{+}^{\left(s_{i}(j)\right)}\right)^{-1} Q, Q_{-}\right) w\left(\Lambda^{(j)}\right) \text { if } \sigma(j)=+ \\
& w\left({ }^{\mathrm{t}} v^{(j-1 / 2)} ; q_{+}^{\left(s_{i}(j)\right)}, Q_{+}\right) w\left({ }^{\mathrm{t}} v^{(j+1 / 2)} ;\left(q_{+}^{\left(s_{i}(j)\right)}\right)^{-1} Q, Q_{-}\right) w\left(\Lambda^{(j)}\right) \text { if } \sigma(j)=-
\end{aligned}
$$

Combining these contributions, we get the claim.
Example 5.4.9. We take $\sigma$ as in Example 2.1.3 and $\lambda=\varnothing$. Assume that $\theta(1 / 2)=$ $N+1 / 2$ and $\theta(5 / 2)=-N+5 / 2$ for $N \gg 0$. In Figure 11, we show the weight (after putting $q_{+}=q_{-}=q_{0}^{1 / 2}$ ) of edges in the area $C_{1 / 2}$. We can idenfity the graph in the area $C_{1 / 2}$ with a hexagon lattice as shown in Figure 12.


Figure 11. The graph in the area $C_{1 / 2}$.


Figure 12. Identification with a hexagon lattice.

Remark. In general, we have the permutation $s_{i} \in \mathfrak{S}_{I_{\mathrm{h}}}$ of the set $I_{\mathrm{h}}$ satisfying the following condition: for sufficiently large $b$ we have

$$
\theta_{i, b}^{-1}\left(s_{i}(1 / 2)\right)<\theta_{i, b}^{-1}\left(s_{i}(3 / 2)\right)<\cdots<\theta_{i, b}^{-1}\left(s_{i}(L-1 / 2)\right) .
$$

The permutation $s_{i}$ determines the direnction in which we take limit in the space of stability conditions. It is the refine topological vertex associated to $Y_{\sigma \circ s_{i}}$ what we get in the limit.
5.5. Conclusion. Note that

$$
\bigcup_{b=1}^{\infty} \Lambda_{\theta_{i, b}}^{\mathrm{re},+\mathrm{e}}=\Lambda_{+}^{\mathrm{re},++}
$$

Combining Theorem 4.5.1 and Theorem 5.4.8, we have:
Theorem 5.5.1. $£_{\sigma, \lambda, v}^{\mathrm{RTV}}$ has the value

$$
\mathscr{E}_{\sigma, \lambda, v}^{\mathrm{NCDT}} \prod_{\alpha \in \Lambda_{+}^{\mathrm{rec},+}}\left(\prod_{\left(h, h^{\prime}\right) \in B_{\sigma, \lambda}^{\alpha,+}}\left(1-\sigma(\alpha) \frac{w_{\lambda}\left(h^{\prime}\right)}{w_{\lambda}(h)}\right)^{\sigma(\alpha)} \prod_{\left(h, h^{\prime}\right) \in B_{\sigma, \lambda}^{\alpha,-}}\left(1-\sigma(\alpha) \frac{w_{\lambda}\left(h^{\prime}\right)}{w_{\lambda}(h)}\right)^{-\sigma(\alpha)}\right) .
$$

(See (1-1), (1-3) and (4-6) for notation.)
Since the second term in this expression does not depend on $v$, we have:
Corollary 5.5.2.

$$
\frac{\mathscr{L}_{\sigma, \lambda, v}^{\mathrm{RTV}}}{\mathscr{L}_{\sigma, \lambda, \lambda, \overline{\mathrm{R}}}^{\mathrm{RTV}}}=\frac{\mathscr{E}_{\sigma, \lambda, v}^{\mathrm{NCDT}}}{\mathscr{E}_{\sigma, \lambda, \bar{\varnothing}}^{\mathrm{NCDT}}} .
$$

Combining Theorem 4.5.3 and Theorem 5.4.8, we have:

## Theorem 5.5.3.

$$
\begin{aligned}
\left.\mathscr{L}_{\sigma, \lambda, \nu}^{\mathrm{RTV}}\right|_{q_{+}=q_{-}} & =\left(q_{0}\right)^{1 / 2} \\
& =\left.\mathscr{L}_{\sigma, \lambda, \nu}^{\mathrm{NCDT}}\right|_{q_{+}=q_{-}=\left(q_{0}\right)^{1 / 2}} \prod_{\alpha \in \Lambda_{+}^{\mathrm{re},+}}\left(1-\sigma(\alpha) \cdot q^{\alpha}\right)^{\sigma(\alpha) \cdot\left[\alpha^{0}+c_{\lambda}\left(j_{-}(\alpha)\right)-c_{\lambda}\left(j_{+}(\alpha)\right)\right]} .
\end{aligned}
$$

(See (1-1), (1-3) and (4-6) for notation.)
Since the second term in the right-hand side depend only on $c_{\lambda}[j]$ 's but not on $\lambda$ and $\nu$, we have the following:

Corollary 5.5.4. If $c_{\lambda}[j]=0$ for any $j$, we have

$$
\frac{\mathscr{L}_{\sigma, \lambda, v}^{\mathrm{RTV}}}{\left.\mathscr{L}_{\sigma, \vec{\varnothing}, \vec{\varnothing} \mathrm{RTV}}\right|_{q_{+}=q_{-}}}=\frac{\mathscr{L}_{\sigma, \lambda, \nu}^{\mathrm{NCDT}}}{\left.\mathscr{L}_{\sigma, \overparen{\varnothing}, \varnothing}^{\mathrm{NCDT}}\right|_{q_{+}=q_{-}}} .
$$

## Acknowledgement

The author is supported by JSPS Fellowships for Young Scientists (No. 19-2672). He thanks to Hiroaki Kanno and Masahito Yamazaki for useful comments.

## References

[Aganagic and Yamazaki 2010] M. Aganagic and M. Yamazaki, "Open BPS wall crossing and Mtheory", Nuclear Phys. B 834:1-2 (2010), 258-272. MR 2011e:81238 Zbl 1204.81132
[Aganagic et al. 2005] M. Aganagic, A. Klemm, M. Mariño, and C. Vafa, "The topological vertex", Comm. Math. Phys. 254:2 (2005), 425-478. MR 2006e:81263 Zbl 1114.81076
[Aganagic et al. 2011] M. Aganagic, H. Ooguri, C. Vafa, and M. Yamazaki, "Wall crossing and M-theory", Publ. Res. Inst. Math. Sci. 47:2 (2011), 569-584. Zbl 05931051 arXiv 0908.1194
[Behrend 2009] K. Behrend, "Donaldson-Thomas type invariants via microlocal geometry", Ann. Math. (2) 170:3 (2009), 1307-1338. MR 2011d:14098 Zbl 1191.14050
[Behrend et al. 2009] K. Behrend, J. Bryan, and B. Szendrői, "Motivic degree zero DonaldsonThomas invariants", preprint, 2009. arXiv 0909.5088
[van den Bergh 2004] M. van den Bergh, "Non-commutative crepant resolutions: the Abel bicentennial", pp. 749-770 in The legacy of Niels Henrik Abel (Oslo, 2002), edited by O. A. Laudal and R. Piene, Springer, Berlin, 2004. Updated in 2009 on arXiv. MR 2005e:14002 Zbl 1082.14005
[Bridgeland et al. 2001] T. Bridgeland, A. King, and M. Reid, "The McKay correspondence as an equivalence of derived categories", J. Amer. Math. Soc. 14:3 (2001), 535-554. MR 2002f:14023 Zbl 0966.14028
[Bryan and Young 2010] J. Bryan and B. Young, "Generating functions for colored 3D Young diagrams and the Donaldson-Thomas invariants of orbifolds", Duke Math. J. 152:1 (2010), 115-153. MR 2011b:14125 Zbl 05692596 arXiv 0802.3948
[Bryan et al. 2012] J. Bryan, C. Cadman, and B. Young, "The orbifold topological vertex", Adv. Math. 229:1 (2012), 531-595.
[Bryan et al. $\geq$ 2011] J. Bryan, C. Cadman, and B. Young, "The crepant resolution conjecture in Donaldson-Thomas theory". To appear.
[Chuang and Jafferis 2009] W.-Y. Chuang and D. L. Jafferis, "Wall crossing of BPS states on the conifold from Seiberg duality and pyramid partitions", Comm. Math. Phys. 292:1 (2009), 285-301. MR 2010e:81222 Zbl 1179.81079
[Chuang and Pan 2010] W.-Y. Chuang and G. Pan, "Bogomolny-Prasad-Sommerfeld state counting in local obstructed curves from quiver theory and Seiberg duality", J. Math. Phys. 51:5 (2010), 052305. MR 2011i:53146
[Dimofte and Gukov 2010] T. Dimofte and S. Gukov, "Refined, motivic, and quantum", Lett. Math. Phys. 91:1 (2010), 1-27. MR 2011m:14094 Zbl 1180.81112
[Dimofte et al. 2011] T. Dimofte, S. Gukov, and Y. Soibelman, "Quantum wall crossing in $\mathcal{N}=2$ gauge theories", Lett. Math. Phys. 95:1 (2011), 1-25. MR 2764330 Zbl 1205.81113
[Iqbal et al. 2009] A. Iqbal, C. Kozçaz, and C. Vafa, "The refined topological vertex", J. High Energy Phys. 10:069 (2009). MR 2011d:81266
[Ishii and Ueda 2008] A. Ishii and K. Ueda, "On moduli spaces of quiver representations associated with dimer models", pp. 127-141 in Higher dimensional algebraic varieties and vector bundles (Kyoto, 2007), edited by S. Mukai, RIMS Kôkyûroku Bessatsu B9, Res. Inst. Math. Sci. (RIMS), Kyoto, 2008. MR 2011e:16024 Zbl 1214.16012
[Jafferis and Moore 2008] D. L. Jafferis and G. W. Moore, "Wall crossing in local Calabi-Yau manifolds", preprint, 2008. arXiv 0810.4909
[Joyce 2007] D. Joyce, "Holomorphic generating functions for invariants counting coherent sheaves on Calabi-Yau 3-folds", Geom. Topol. 11 (2007), 667-725. MR 2008d:14062 Zbl 1141.14023
[Joyce 2008] D. Joyce, "Configurations in abelian categories, IV: Invariants and changing stability conditions", Adv. Math. 217:1 (2008), 125-204. MR 2009d:18015b Zbl 1134.14008
[Joyce and Song 2010] D. Joyce and Y. Song, "A theory of generalized Donaldson-Thomas invariants", Mem. Amer. Math. Soc. (2010). arXiv 0810.5645
[Kapranov and Vasserot 2000] M. Kapranov and E. Vasserot, "Kleinian singularities, derived categories and Hall algebras", Math. Ann. 316:3 (2000), 565-576. MR 2001h:14012 Zbl 0997.14001
[Kontsevich and Soibelman 2008] M. Kontsevich and Y. Soibelman, "Stability structures, motivic Donaldson-Thomas invariants and cluster transformations", preprint, 2008. arXiv 0811.2435
[Maulik et al. 2006] D. Maulik, N. Nekrasov, A. Okounkov, and R. Pandharipande, "GromovWitten theory and Donaldson-Thomas theory, I", Compos. Math. 142:5 (2006), 1263-1285. MR 2007i:14061 Zbl 1108.14046
[Mozgovoy and Reineke 2010] S. Mozgovoy and M. Reineke, "On the noncommutative DonaldsonThomas invariants arising from brane tilings", Adv. Math. 223:5 (2010), 1521-1544. MR 2011e: 16026 Zbl 1191.14008
[Nagao 2011a] K. Nagao, "Derived categories of small toric Calabi-Yau 3-folds and counting invariants", preprint, 2011. arXiv 0809.2994
[Nagao 2011b] K. Nagao, "Non-commutative Donaldson-Thomas theory and vertex operators", Geom. Topol. 15:3 (2011), 1509-1543. Zbl 1219.14066 arXiv 0910.5477
[Nagao and Nakajima 2011] K. Nagao and H. Nakajima, "Counting invariant of perverse coherent sheaves and its wall-crossing", Internat. Math. Res. Notices 17 (2011), 3885-3938. Zbl 05957514 arXiv 0809.2992
[Nagao and Yamazaki 2010] K. Nagao and M. Yamazaki, "The non-commutative topological vertex and wall crossing phenomena", Adv. Theor. Math. Phys. 14:4 (2010), 1147-1181. Zbl 05973794 arXiv 0910.5479
[Ooguri and Yamazaki 2009] H. Ooguri and M. Yamazaki, "Crystal melting and toric Calabi-Yau manifolds", Comm. Math. Phys. 292:1 (2009), 179-199. MR 2011d:81268 Zbl 1179.81139
[Szendrői 2008] B. Szendrői, "Non-commutative Donaldson-Thomas invariants and the conifold", Geom. Topol. 12:2 (2008), 1171-1202. MR 2009e:14100 Zbl 1143.14034
[Thomas 2000] R. P. Thomas, "A holomorphic Casson invariant for Calabi-Yau 3-folds, and bundles on K3 fibrations", J. Differential Geom. 54:2 (2000), 367-438. MR 2002b:14049 Zbl 1034.14015
[Toda 2009] Y. Toda, "Limit stable objects on Calabi-Yau 3-folds", Duke Math. J. 149:1 (2009), 157-208. MR 2011b:14043 Zbl 1172.14007
[Toda 2010] Y. Toda, "Generating functions of stable pair invariants via wall-crossings in derived categories", pp. 389-434 in New developments in algebraic geometry, integrable systems and mirror symmetry (Kyoto, 2008), edited by M.-H. Saito et al., Adv. Stud. Pure Math. 59, Math. Soc. Japan, Tokyo, 2010. MR 2683216 Zbl 1216.14009 arXiv 0806.0062
[Young 2009] B. Young, "Computing a pyramid partition generating function with dimer shuffling", J. Combin. Theory Ser. A 116:2 (2009), 334-350. MR 2009k:05016 Zbl 1191.05007

Received November 13, 2009. Revised February 2, 2010.
Kentaro Nagao
Graduate School of Mathematics
Nagoya University
Furocho, Chikusaku
NAGOYA, 464-8602
JAPAN
kentaron@math.nagoya-u.ac.jp

# PACIFIC JOURNAL OF MATHEMATICS 

http://pacificmath.org<br>Founded in 1951 by<br>E. F. Beckenbach (1906-1982) and F. Wolf (1904-1989)

## EDITORS

V. S. Varadarajan (Managing Editor)

Department of Mathematics University of California
Los Angeles, CA 90095-1555
pacific@math.ucla.edu

| Vyjayanthi Chari | Darren Long | Sorin Popa |
| :---: | :---: | :---: |
| Department of Mathematics | Department of Mathematics | Department of Mathematics |
| University of California | University of California | University of California |
| Riverside, CA 92521-0135 | Santa Barbara, CA 93106-3080 | Los Angeles, CA 90095-1555 |
| chari@math.ucr.edu | long@ math.ucsb.edu | popa@ math.ucla.edu |
| Robert Finn | Jiang-Hua Lu | Jie Qing |
| Department of Mathematics <br> Stanford University | Department of Mathematics | Department of Mathematics |
| Stanford, CA 94305-2125 | The University of Hong Kong | University of California |
| finn@math.stanford.edu | Pokfulam Rd., Hong Kong | Santa Cruz, CA 95064 |
| Kefeng Liu | jhlu@ maths.hku.hk | qing @cats.ucsc.edu |
| Department of Mathematics | Alexander Merkurjev | Jonathan Rogawski |
| University of California | Department of Mathematics | Department of Mathematics |
| Los Angeles, CA 90095-1555 | University of California | University of California |
| liu@math.ucla.edu | Los Angeles, CA 90095-1555 | Los Angeles, CA 90095-1555 |
|  | merkurev@ math.ucla.edu | jonr@ math.ucla.edu |

## PRODUCTION

pacific@math.berkeley.edu

Silvio Levy, Scientific Editor Matthew Cargo, Senior Production Editor

## SUPPORTING INSTITUTIONS

ACADEMIA SINICA, TAIPEI
CALIFORNIA INST. OF TECHNOLOGY
INST. DE MATEMÁTICA PURA E APLICADA KEIO UNIVERSITY

MATH. SCIENCES RESEARCH INSTITUTE NEW MEXICO STATE UNIV.
OREGON STATE UNIV

STANFORD UNIVERSITY
UNIV. OF BRITISH COLUMBIA
UNIV. OF CALIFORNIA, BERKELEY
UNIV. OF CALIFORNIA, DAVIS
UNIV. OF CALIFORNIA, LOS ANGELES
UNIV. OF CALIFORNIA, RIVERSIDE
UNIV. OF CALIFORNIA, SAN DIEGO
UNIV. OF CALIF., SANTA BARBARA

UNIV. OF CALIF., SANTA CRUZ
UNIV. OF MONTANA
UNIV. OF OREGON
UNIV. OF SOUTHERN CALIFORNIA
UNIV. OF UTAH
UNIV. OF WASHINGTON
WASHINGTON STATE UNIVERSITY

These supporting institutions contribute to the cost of publication of this Journal, but they are not owners or publishers and have no responsibility for its contents or policies.

See inside back cover or pacificmath.org for submission instructions.
The subscription price for 2011 is US $\$ 420 /$ year for the electronic version, and $\$ 485 /$ year for print and electronic.
Subscriptions, requests for back issues from the last three years and changes of subscribers address should be sent to Pacific Journal of Mathematics, P.O. Box 4163, Berkeley, CA 94704-0163, U.S.A. Prior back issues are obtainable from Periodicals Service Company, 11 Main Street, Germantown, NY 12526-5635. The Pacific Journal of Mathematics is indexed by Mathematical Reviews, Zentralblatt MATH, PASCAL CNRS Index, Referativnyi Zhurnal, Current Mathematical Publications and the Science Citation Index.

The Pacific Journal of Mathematics (ISSN 0030-8730) at the University of California, c/o Department of Mathematics, 969 Evans Hall, Berkeley, CA 94720-3840, is published monthly except July and August. Periodical rate postage paid at Berkeley, CA 94704, and additional mailing offices. POSTMASTER: send address changes to Pacific Journal of Mathematics, P.O. Box 4163, Berkeley, CA 94704-0163.

## PACIFIC JOURNAL OF MATHEMATICS

Volume 254 No. $1 \quad$ November 2011
A mean curvature estimate for cylindrically bounded submanifolds ..... 1
Luis J. Alías and Marcos Dajczer
Weyl group multiple Dirichlet series of type $C$ ..... 11
Jennifer Beineke, Benjamin Brubaker and Sharon Frechette
Milnor open books of links of some rational surface singularities ..... 47
Mohan Bhupal and Burak Ozbagci
Simple closed curves, word length, and nilpotent quotients of free groups ..... 67
Khalid Bou-Rabee and Asaf Hadari
Strong submodules of almost projective modules ..... 73
GÁbor Braun and Jan Trlifaj
Interlacing log-concavity of the Boros-Moll polynomials ..... 89
William Y. C. Chen, Larry X. W. Wang and Ernest X. W. Xia
Schwarzian norms and two-point distortion ..... 101
Martin Chuaqui, Peter Duren, William Ma, Diego Mejía, David Minda and Brad Osgood
The principle of stationary phase for the Fourier transform of $D$-modules ..... 117 Jianguue Fang
Monotonicity and uniqueness of a 3D transonic shock solution in a conic nozzle with ..... 129 variable end pressureJun Li, Zhouping Xin and Huicheng Yin
Refined open noncommutative Donaldson-Thomas invariants for small crepant ..... 173 resolutionsKentaro Nagao
The Dirichlet problem for harmonic functions on compact sets ..... 211
Tony L. Perkins
Extension of an analytic disc and domains in $\mathbb{C}^{2}$ with noncompact automorphism ..... 227group
Minju Song
Regularity of the first eigenvalue of the $p$-Laplacian and Yamabe invariant along ..... 239
geometric flowsEr-Min Wang and Yu Zheng


[^0]:    MSC2000: 14N10, 14N35.
    Keywords: Donaldson-Thomas theory, dimer model, topological vertex.

[^1]:    ${ }^{1}$ The word "open" stems from such terminologies as "open topological string theory". According to [Aganagic et al. 2005], open topological string partition function is given by summing up the generating functions of these invariants over Young diagrams.
    ${ }^{2}$ As far as we know, there is no definition of "open" invariants for general Calabi-Yau 3-folds.
    ${ }^{3}$ See [Behrend et al. 2009] for a geometric definition of refined invariants.

[^2]:    ${ }^{4}$ The Euler characteristic version of the Donaldson-Thomas invariant coincides with the Donald-son-Thomas invariant up to sign [Maulik et al. 2006].
    ${ }^{5}$ The Euler characteristic version of the noncommutative Donaldson-Thomas invariant coincides with the noncommutative Donaldson-Thomas invariant up to sign [Nagao 2011a; Mozgovoy and Reineke 2010].
    ${ }^{6}$ A moduli space of stable $A$-modules with the specific numerical data gives a crepant resolution of $X$ [Ishii and Ueda 2008]. The direction in which we take limit in the space of stability conditions determines a stability parameter in the construction of a crepant resolution.

[^3]:    ${ }^{7}$ From the geometric point of view, the crystal melting model is more natural. But in this paper we adapt the definition using the dimer model since it is more convenient when we prove some technical lemmas, which we also use in [Nagao 2011b].
    ${ }^{8}$ In the case when $v_{+}=v_{-}=\varnothing$, the moduli spaces have symmetric obstruction theory and the invariant in this paper coincides with the weighted Euler characteristic up to sign.

[^4]:    ${ }^{9}$ We will leave the definition of the generating function until Section 3.4 since we will use Proposition 3.3.9 to prove that the number of dimer configurations of type $(\sigma, \lambda, \nu, \theta)$ is finite.

